



33<sup>rd</sup> International Symposium of the Society of Core Analysts, Pau, France, Aug 26-29, 2019.

SOCIETY OF  
CORE ANALYSTS

# Short Course

## Interpretation of Pore Scale Experiments - Alternative Descriptions of Porous Media Flow by Topological Means

Ryan T. Armstrong<sup>1</sup> and Steffen Berg<sup>2,3</sup>



**UNSW**  
SYDNEY

<sup>1</sup>University of New South Wales, Australia.

<sup>2</sup>Shell Global Solutions International B.V.

<sup>3</sup>Imperial College London



**Imperial College  
London**

With key contributions from

**James McClure** (Virginia Tech), **Steffen Schlüter** (UFZ Potsdam), **Anna Herring** (ANU), **Christoph Arns** (UNSW)



# Outline

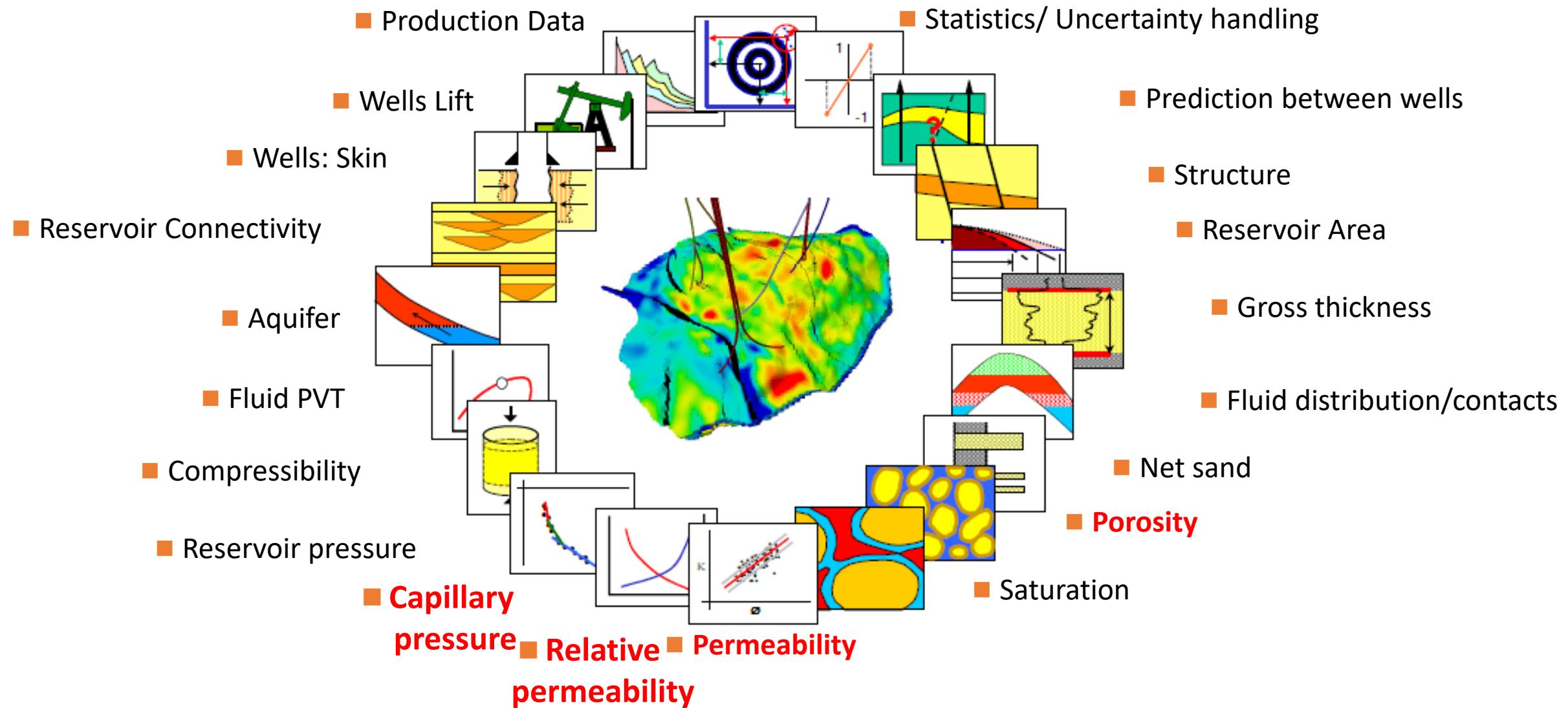
Why

How

What

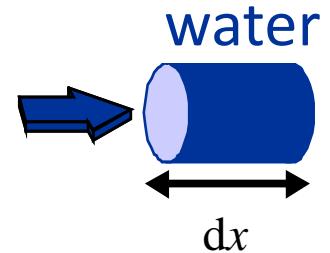
1. Motivation / necessity for and advantages of a geometric state function
  - 2-phase Darcy = phenomenological extension from 1-phase to 2-phase flow
  - example of having incomplete state variables: ideal gas equation of state
2. Introduction:
  - introducing Minkowski functionals
  - Hadwiger's theorem, Gauss-Bonnet Theorem, Steiner's formula
3. Description how we found it
  - general background: ganglion dynamics → topological changes
  - beamline data of 2013 experiment: → Hysteresis in Euler characteristic  $\chi$
  - proof: (1) Hadwiger's theorem, (2) 220000 LBM simulations
4. Software: Avizo, FIJI (BoneJ plugin), Matlab, Python, Dragonfly/DeepRocks
5. Applications,
  - hysteresis model for relative permeability
  - new route to pc and relative permeability
  - Digital Rock: validation of pore scale simulation techniques
  - wettability: description of contact angle as deficit curvature

# Integrated Subsurface Workflow



# Multiphase Flow in Porous Media at Darcy Scale

## Single-Phase



$$v_{Darcy} = -\frac{K}{\mu} \frac{dp}{dx}$$

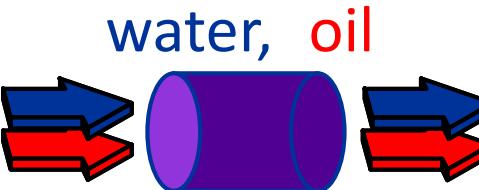
**Darcy's law**

$\mu$  viscosity

$P$  pressure

$K$  ( $K_{abs}$ ) absolute permeability

Viscous law (similar to pipe flow)  
can be derived from upscaling  
Stokes flow at pore scale by  
homogenization



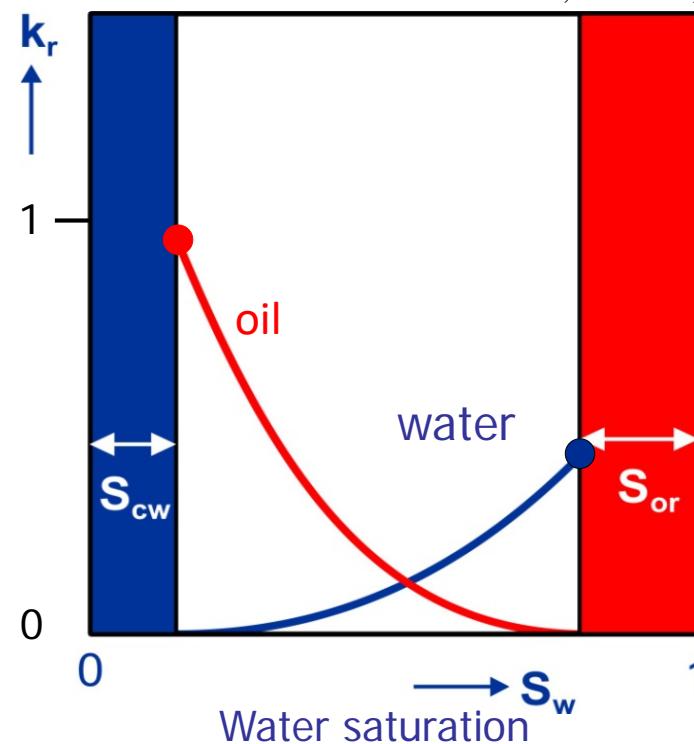
## Two-Phase

$i = w, o$

$$v_i = -k_{r,i} \frac{K}{\mu_i} \frac{dp_i}{dx}$$

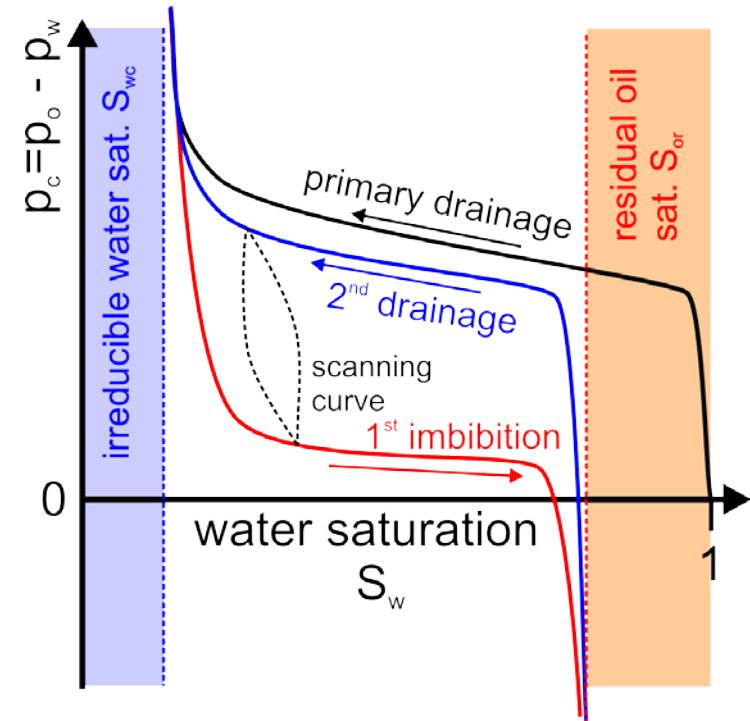
Phenomenological  
extension of Darcy's law

relative permeability  $k_{r,i} = k_{r,i}(S_w)$



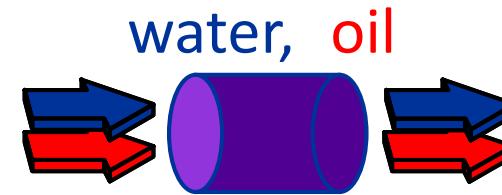
Capillary pressure

$$p_c = p_o - p_w$$



# Historical Overview – Classification of the Problem

- 1930s: **flow problem**: mass & momentum balance  
(Wykoff & Botset, Muskat & Meres, Leverett)

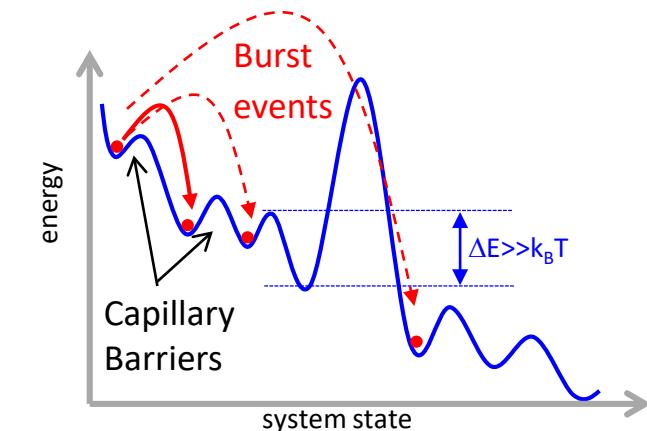


- 1970s: thermodynamics of pore scale displacements  
(Morrow, Swanson & Yuan)

$$\Delta F = -S\Delta T - \sum_{\alpha=1}^2 p_{\alpha} \Delta V_{\alpha} + \sigma_{1,2} \Delta A_{1,2}$$

Morrow 1970

- 1990s: **equilibrium thermodynamics problem**:  
mass, momentum & energy balance (Hassanizadeh & Gray)  
ganglion dynamics in 2D micromodels (Avraam & Payatakes)

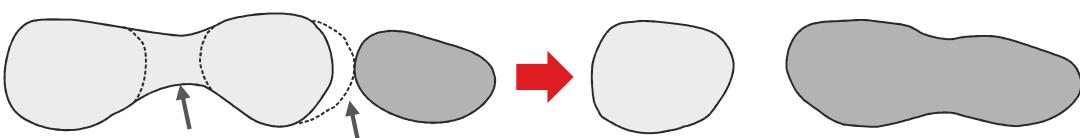


- 2000s: **non-equilibrium thermodynamic problem** – no global energy minimum  
TCAT (Gray & Miller)

- 2010s: **non-local dynamics**, ganglion dynamics in 3D

$S_w + A_{nw}$  does not close pc hysteresis

Armstrong, McClure et al. 2018



- 2018: **state variables**:  $S_w$ ,  $A_{nw}$ ,  $p_c$ ,  $\chi$

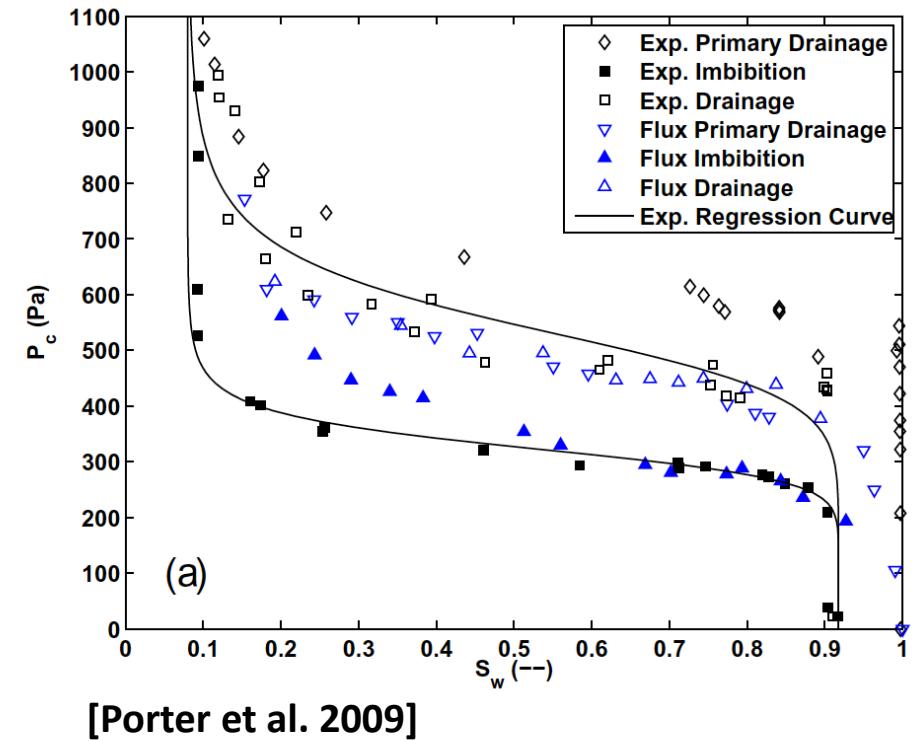
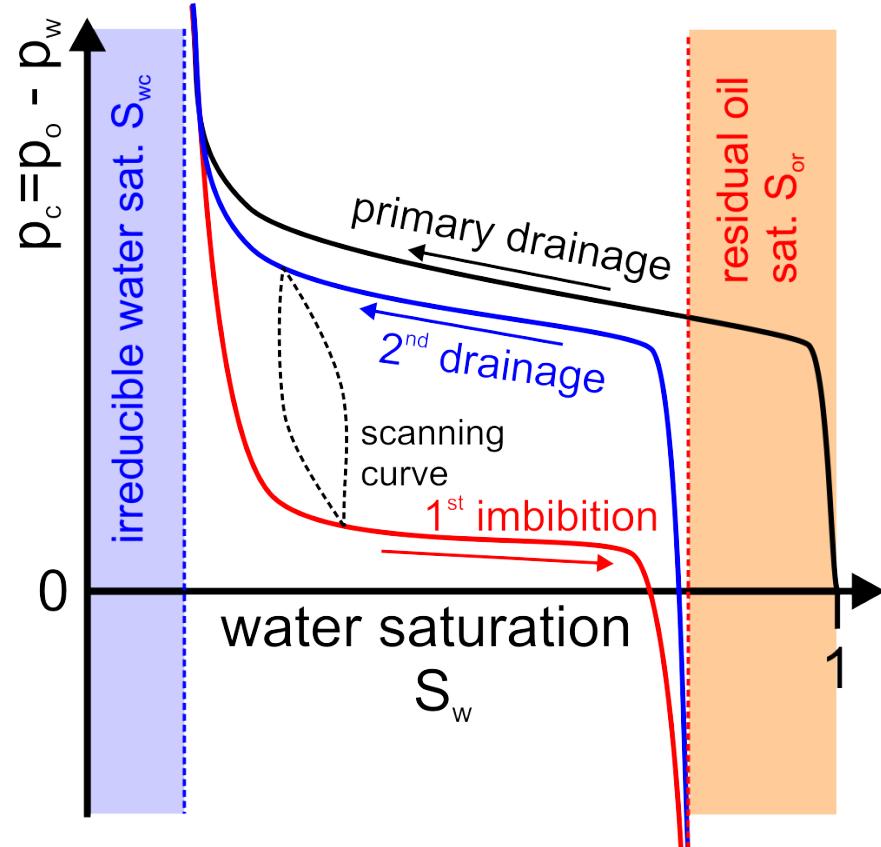
- 2019: **wettability** (upscaled contact angle as deficit Gaussian curvature)

# Unresolved Issue: Capillary Pressure - Hysteresis

Capillary pressure function of saturation only

$$P_{mw} - P_w = P_c = f(S_w)$$

But: hysteresis



Hysteresis challenges the validity of 2-phase Darcy. But is it really hysteresis ?  
What constitutive relationships properly represent multiphase flow?

# Consequence of Insufficient Number of State Variables

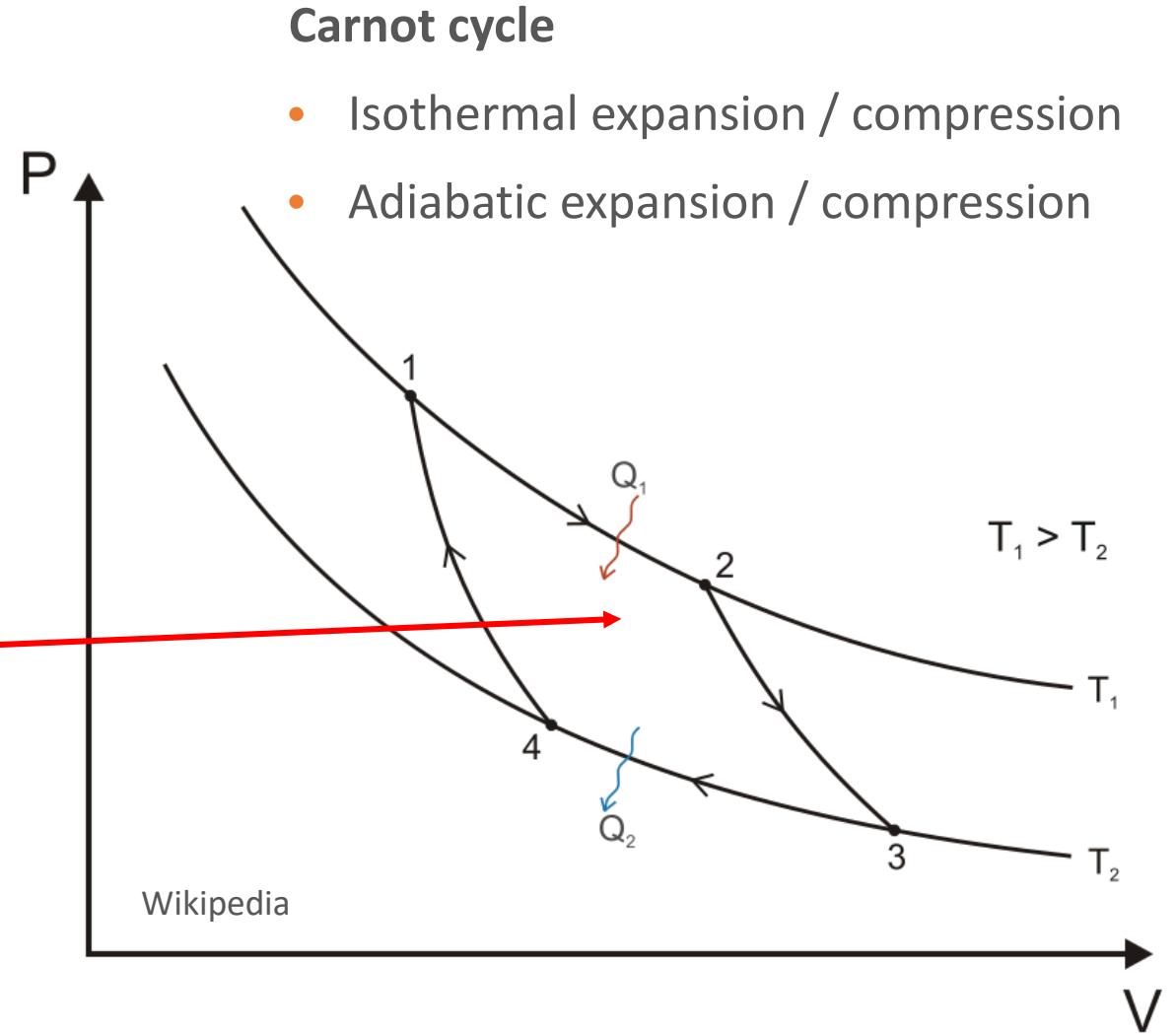
Example: ideal gas, equation of state

$$p \cdot V = nRT$$

Imagine we did not know about Temperature ...  
and only do  $P - V$  experiments,  
without measuring or controlling  $T$  ...

Apparent  
hysteresis

But with correct number of state variables,  $p, V, T$ ,  
there is no hysteresis (for an ideal gas).



# Saturation $S_w$ and Interfacial Area $A_{nw}$ are State Variables of $P_c$

$$\Delta F = -S\Delta T - \sum_{\alpha=1}^2 p_\alpha \Delta V_\alpha + \sigma_{1,2} \Delta A_{1,2}$$

Pressure-volume work      Interfacial Energy term

Thermodynamics: free energy

Morrow, 1970

But that did not close  
the capillary hysteresis  
(McClure, Gray, Miller ...)



This term was not included in the traditional theories

4 Minkowski Functionals

$m_0$  = volume (saturation)

$m_1$  = interfacial area

$m_2$  = mean curvature

$m_3$  = Gauss curvature =  $2\pi\chi$      $M_3^n = \int_{\Gamma_n} \frac{1}{R_1 R_2} dr .$

$$M_0^n = \lambda(\Omega_n) = \int_{\Omega_n} dr$$

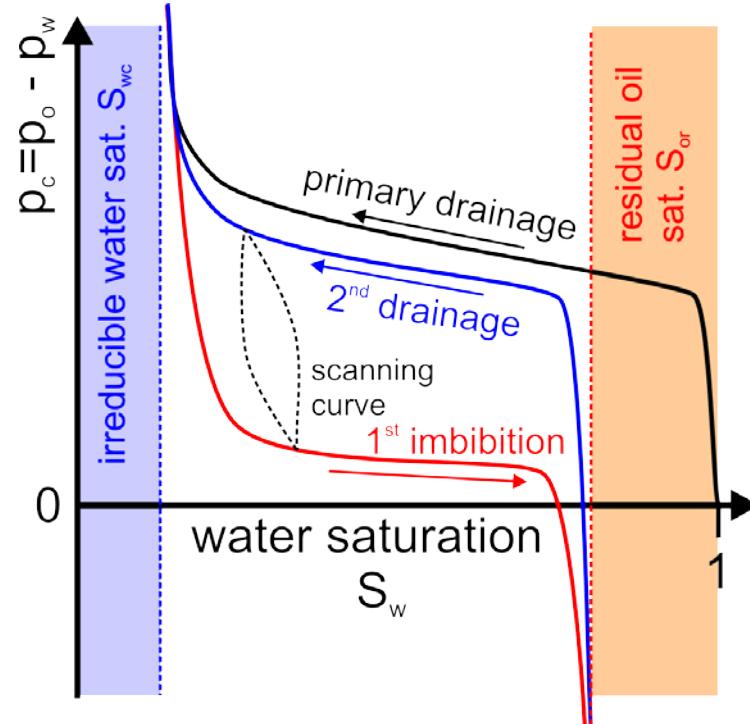
$$M_1^n = \lambda(\Gamma_n) = \int_{\Gamma_n} dr$$

$$M_2^n = \int_{\Gamma_n} \left( \frac{1}{R_1} + \frac{1}{R_2} \right) dr$$

# The Source of Capillary Pressure Hysteresis

Darcy Scale

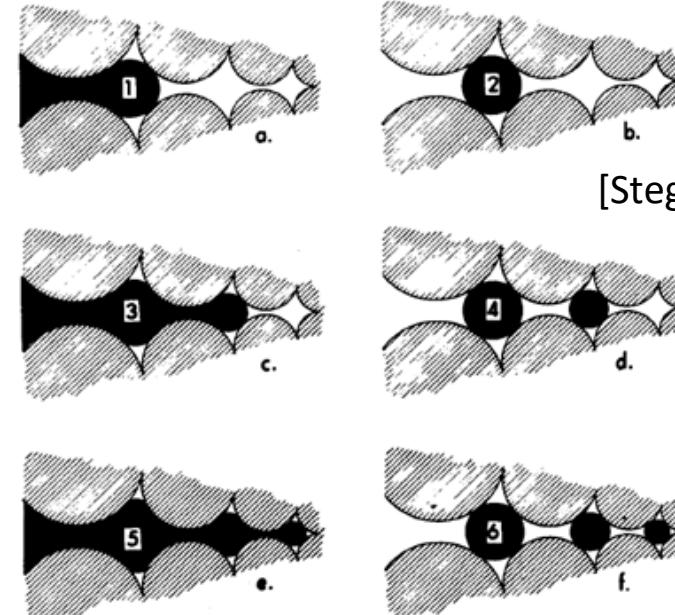
Macroscopic Darcy-scale  
“saturation functions”



=

Pore Scale

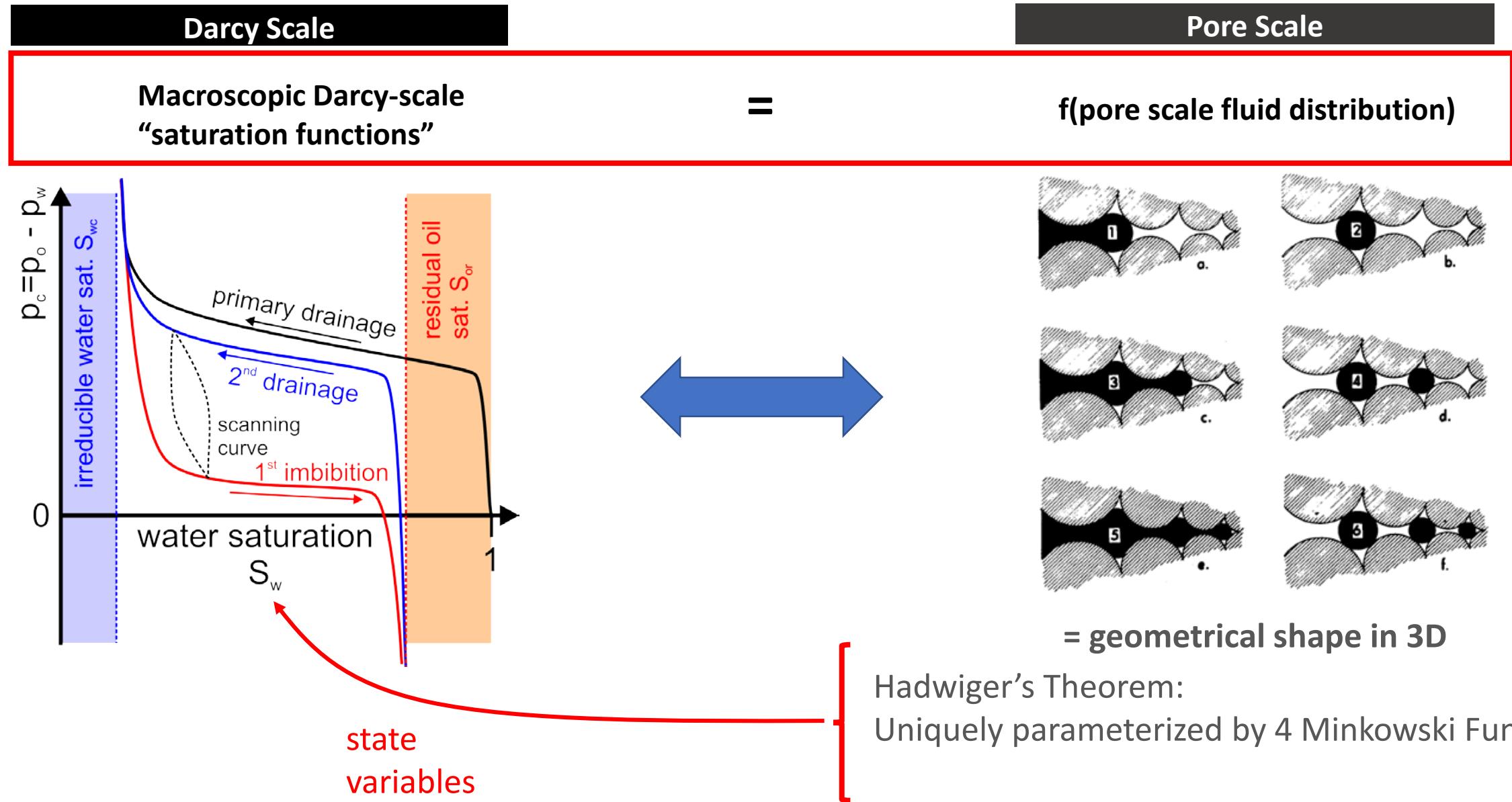
$f(\text{pore scale fluid distribution})$



[Stegemeier 1977]

= geometrical shape in 3D

# The Source of Capillary Pressure Hysteresis



# The Minkowski Functionals

→ Apply to Multiphase Flow

Named after **Hermann Minkowski**, Mathematician (1864-1909)

Herring et al. Advances in Water Resources 62, 47-58, 2013.

McClure et al. Phys. Rev. Fluids, 2018

**Hadwiger's theorem:** unique characterization of 3D objects by 4 Minkowski functionals

$m_0$  = volume (saturation)

$$M_0^n = \lambda(\Omega_n) = \int_{\Omega_n} dr$$

$m_1$  = interfacial area

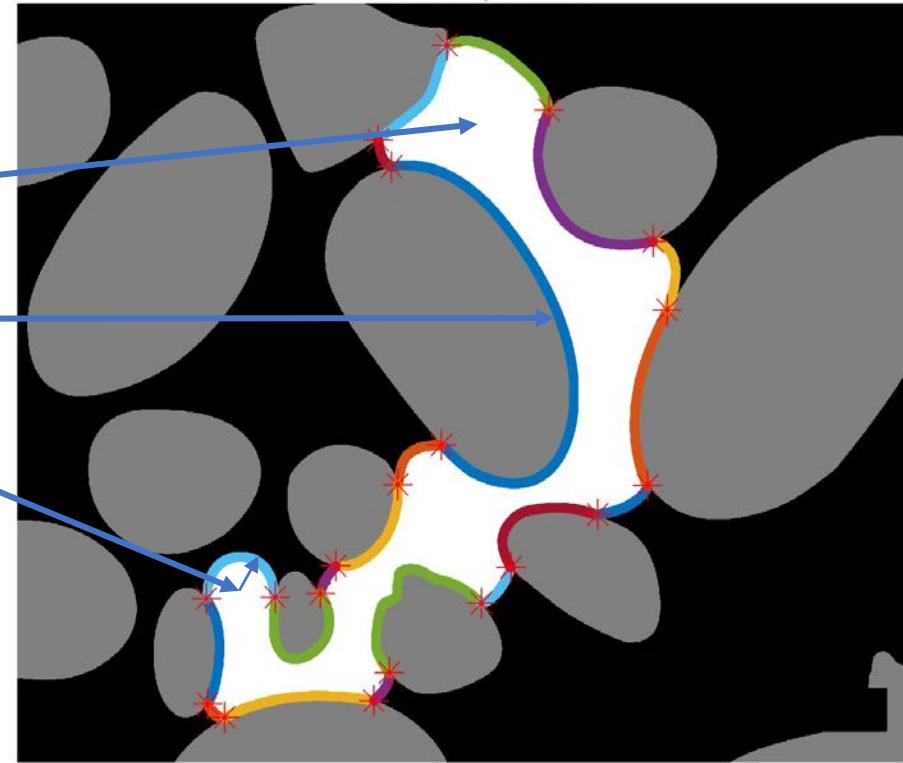
$$M_1^n = \lambda(\Gamma_n) = \int_{\Gamma_n} dr$$

$m_2$  = mean curvature (cap. pressure)

$$M_2^n = \int_{\Gamma_n} \left( \frac{1}{R_1} + \frac{1}{R_2} \right) dr$$

$m_3$  = integral curvature =  $2\pi\chi$

$$M_3^n = \int_{\Gamma_n} \frac{1}{R_1 R_2} dr .$$



Klaus R. Mecke, Dietrich Stoyan, Statistical Physics and Spatial Statistics. The Art of Analyzing and Modeling Spatial Structures and Pattern Formation, Lecture Notes in Physics, Springer, 2000.

C. H. Arns, M. A. Knackstedt, K. Mecke, 3D Structural Analysis: Sensitivity of Minkowski Functionals. Journal of Microscopy 240, 181-196, 2010.

H.J. Vogel, U. Weller, S. Schlüter, Quantification of Soil Structure Based on Minkowski Functions, Computers & Geosciences 36, 126-1251, 2010.

# The Euler Characteristic

Named after **Leonhard Euler**, German Mathematician (1707-1783)

$$M_3(X) = \int_{\delta X} [1/(r_1 r_2)] ds = 2\pi \chi(\delta X) = 4\pi \chi(X)$$

Gaussian curvature

$$\frac{1}{r_1} + \frac{1}{r_2}$$
 mean curvature

$$\frac{1}{r_1} \cdot \frac{1}{r_2}$$
 Gaussian curvature

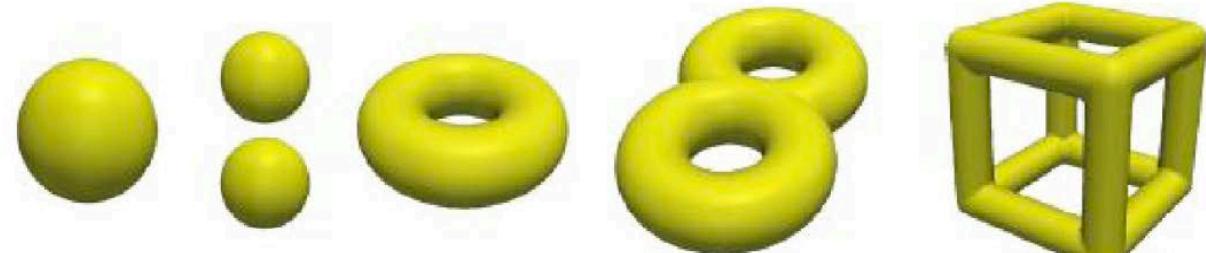
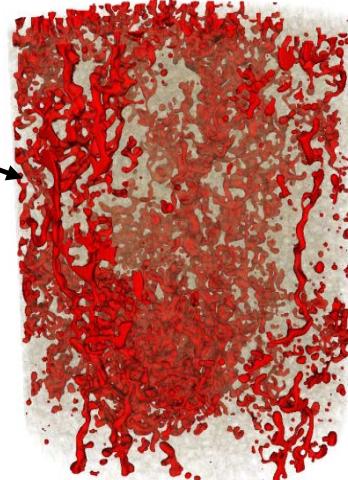
- Euler Characteristic measures the **bulk connectivity** of an object

$$X = Objects - Loops + Voids$$

[Herring et al. 2012]

Or a collection of objects:

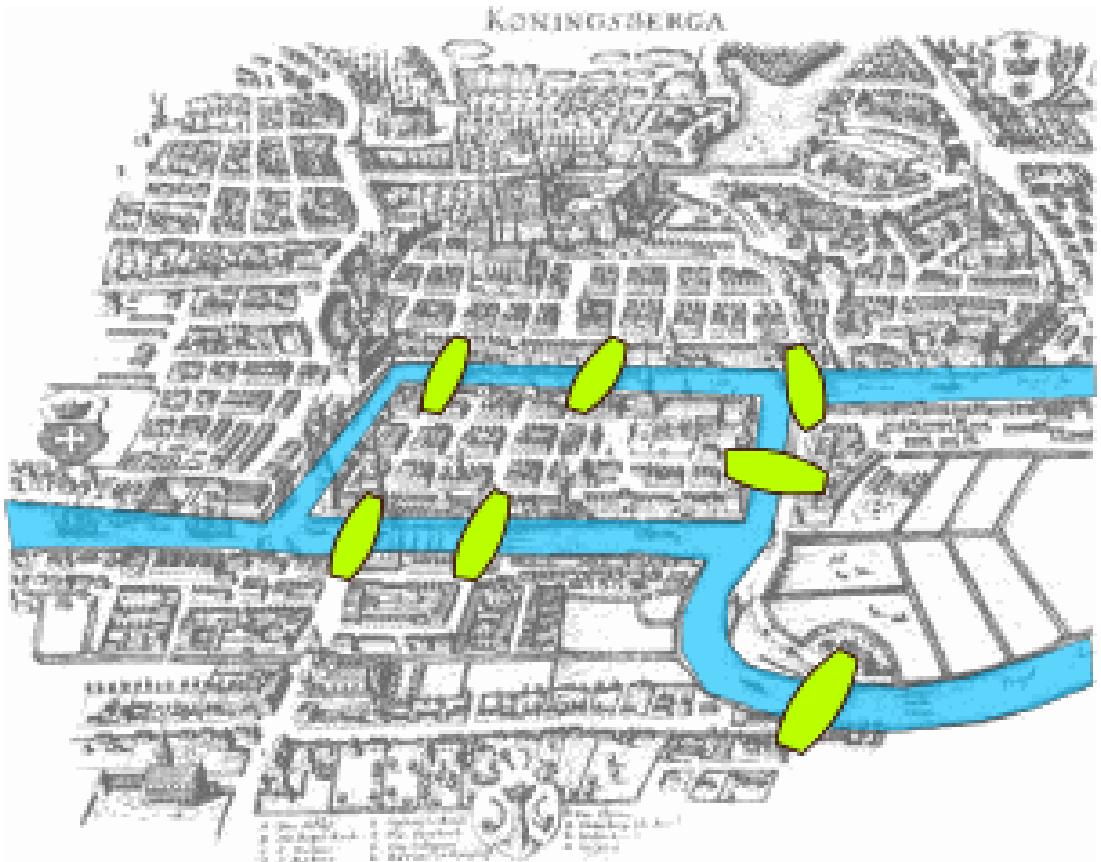
OIL



$$X(OIL) = -40$$

... a cube has 6 faces

# Leonhard Euler and the Seven Bridges of Königsberg



2 islands  
separated by a river  
7 bridges

Walk to cross each bridge only once ?

L. Euler: not possible + proof

The city of [Königsberg](#) in [Prussia](#) (now [Kaliningrad, Russia](#)) was set on both sides of the [Pregel River](#), and included two large islands - [Kneiphof](#) and [Lomse](#) - which were connected to each other, or to the two mainland portions of the city, by seven bridges. The problem was to devise a walk through the city that would cross each of those bridges once and only once (Source: Wikipedia)

[https://en.wikipedia.org/wiki/Seven\\_Bridges\\_of\\_K%C3%B6nigsberg](https://en.wikipedia.org/wiki/Seven_Bridges_of_K%C3%B6nigsberg)

# The Euler Characteristic

$$M_3(X) = \int_{\delta X} [1/(r_1 r_2)] ds = 2\pi\chi(\delta X) = 4\pi\chi(X)$$

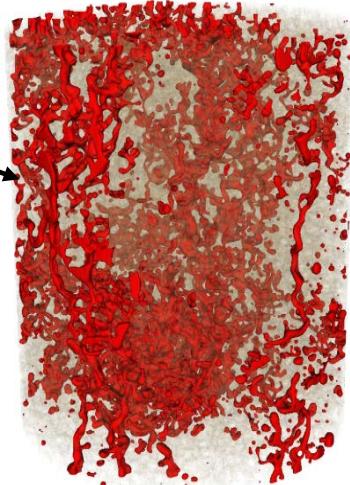
- Euler Characteristic measures the **bulk connectivity** of an object

$$\chi = \text{Objects} - \text{Loops} + \text{Voids}$$

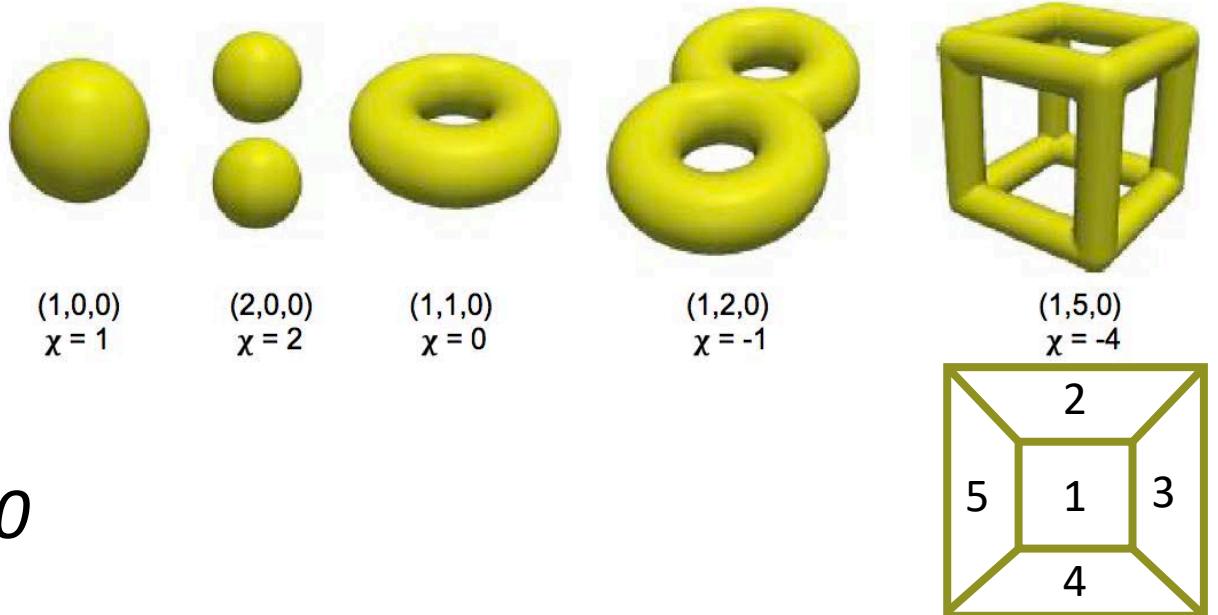
[Herring et al. 2012]

Or a collection of objects:

OIL



$$\chi(\text{OIL}) = -40$$



# (1) Hadwiger's Theorem

A remarkable theorem is the ‘completeness’ of the Minkowski functionals proven 1957 by H. Hadwiger [21]. This *characterization theorem* asserts that any additive, motion-invariant and conditionally continuous functional  $\mathcal{M}$  is a linear combination of the  $d + 1$  Minkowski functionals  $M_\nu$ ,

$$\mathcal{M}(A) = \sum_{\nu=0}^d c_\nu M_\nu(A) , \quad (7)$$

with real coefficients  $c_\nu$  independent of  $A$ . Motion-invariance of the functional means that the functional  $\mathcal{M}$  does not depend on the location and orientation of the grain  $A$ . Since quite often the assumption of a homogeneous and isotropic system is made in physics, motion-invariance is not a very restrictive constraint on the functional. Nevertheless, in the case where external fields are applied

## (2) Gauss-Bonnet Theorem

Explains the relationship between Gaussian curvature and topology.

Gaussian curvature

$$k_f = \left( \frac{1}{r_1} \cdot \frac{1}{r_2} \right)$$

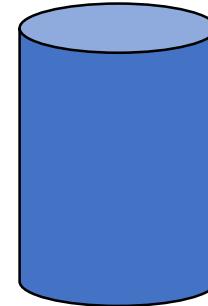
$$p_c(S_w) = \gamma \cdot \left( \frac{1}{r_1} + \frac{1}{r_2} \right)$$



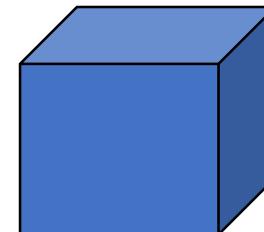
Mean curvature

$$\int_S [1/(r_1 r_2)] ds + \int_B k_g dl = 2\pi\chi(S)$$

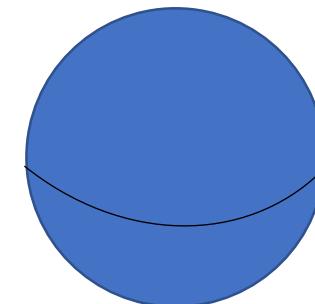
$$\chi(S) = 2$$



$$\chi(S) = 2$$



$$\chi(S) = 2$$



The curvature is either on the **surface (Gaussian Curvature)** or at the **edges (Geodesic Curvature)**

### (3) Steiner's Formula

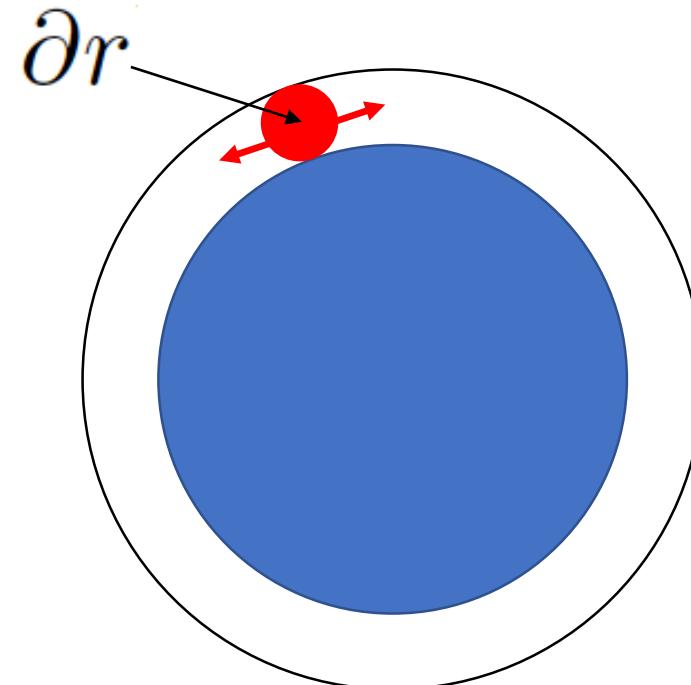
Provides a means to identify relationships between the Minkowski functionals

Volume



$$\lambda(X \oplus \partial r) - \lambda(X) = \sum_{i=1}^3 a_i M_i r^i$$

Explains how the volume (dependent) of an object changes depending on the objects morphology (independent)



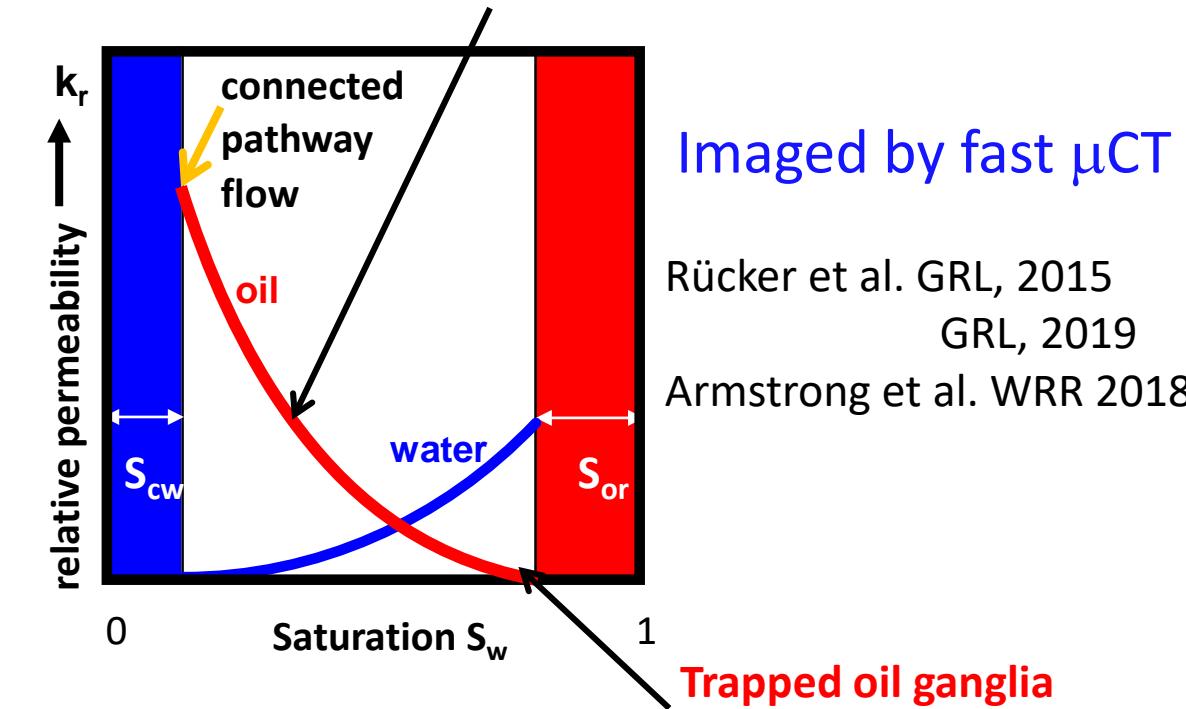
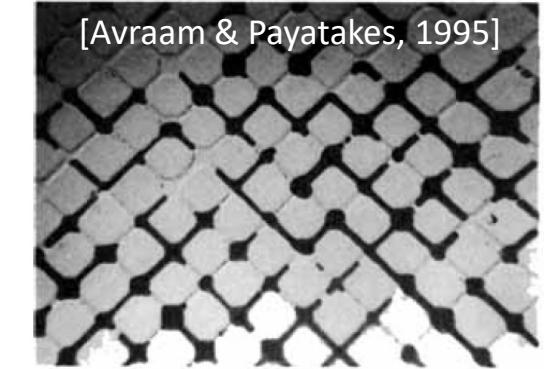
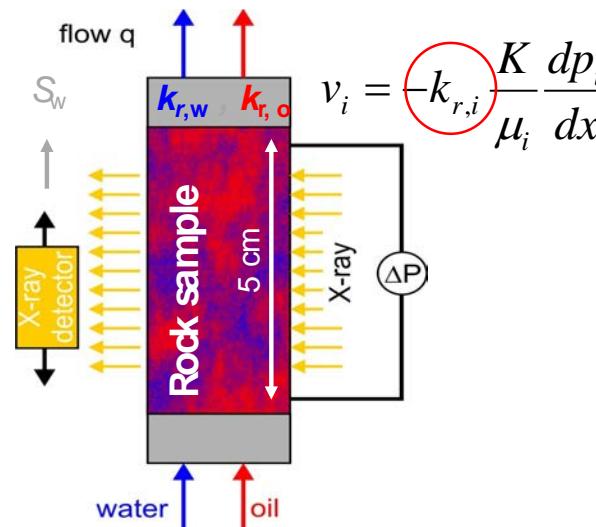
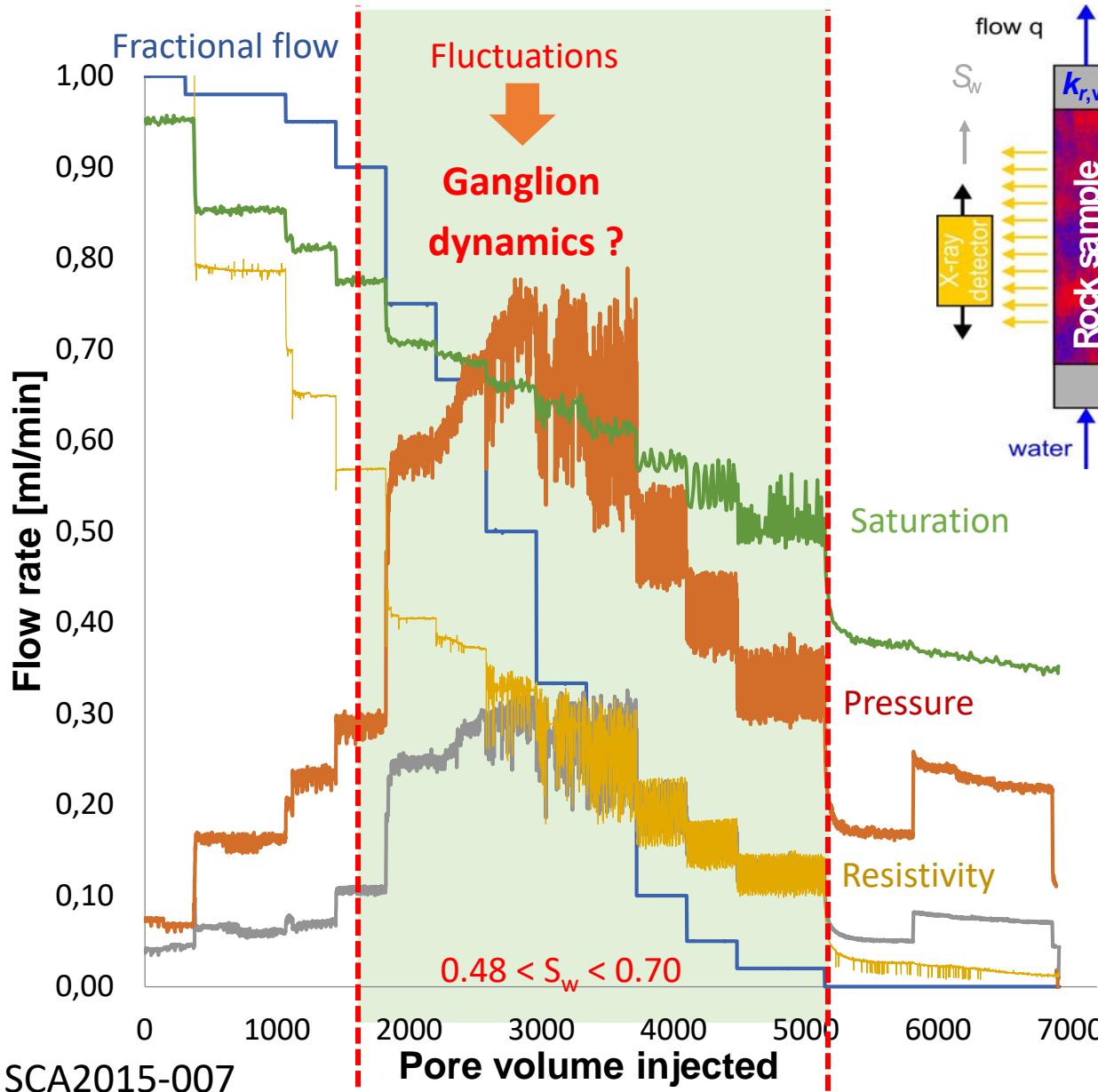
Example: sphere ...

$$V(\Omega_i \oplus \delta\zeta) - V(\Omega_i) = \frac{4}{3}\pi(r + \delta r)^2 - \frac{4}{3}\pi(r)^2 = A_i \delta r + H_i (\delta r)^2 + \frac{4}{3}\pi \chi_i (\delta r)^3$$

↑  
Mean width  $H_i = \int_{\Gamma_i} \frac{\kappa_1 + \kappa_2}{2} dS$

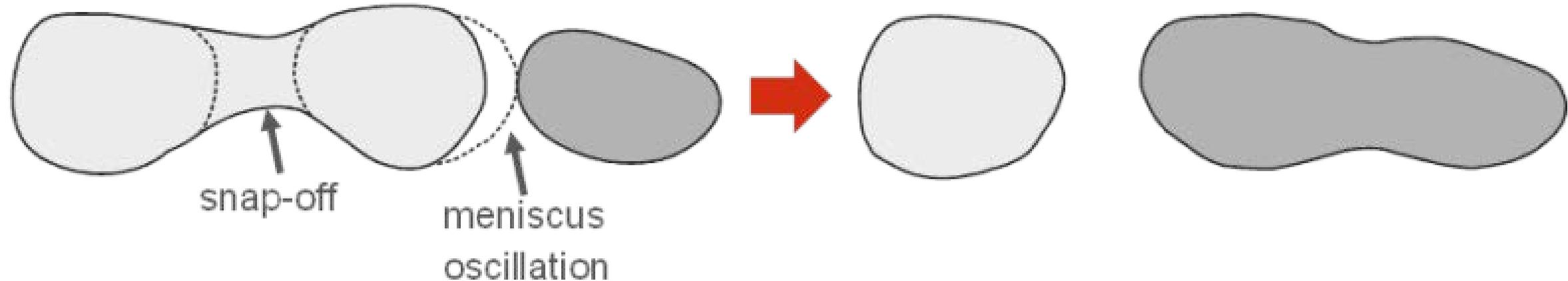
$A_i = 4\pi r^2$   
 $H_i = 4\pi r$   
 $\chi_i = 1$

# Background: Cluster Dynamics in SCAL Experiments



# Cluster Dynamics Introduces Topological Changes

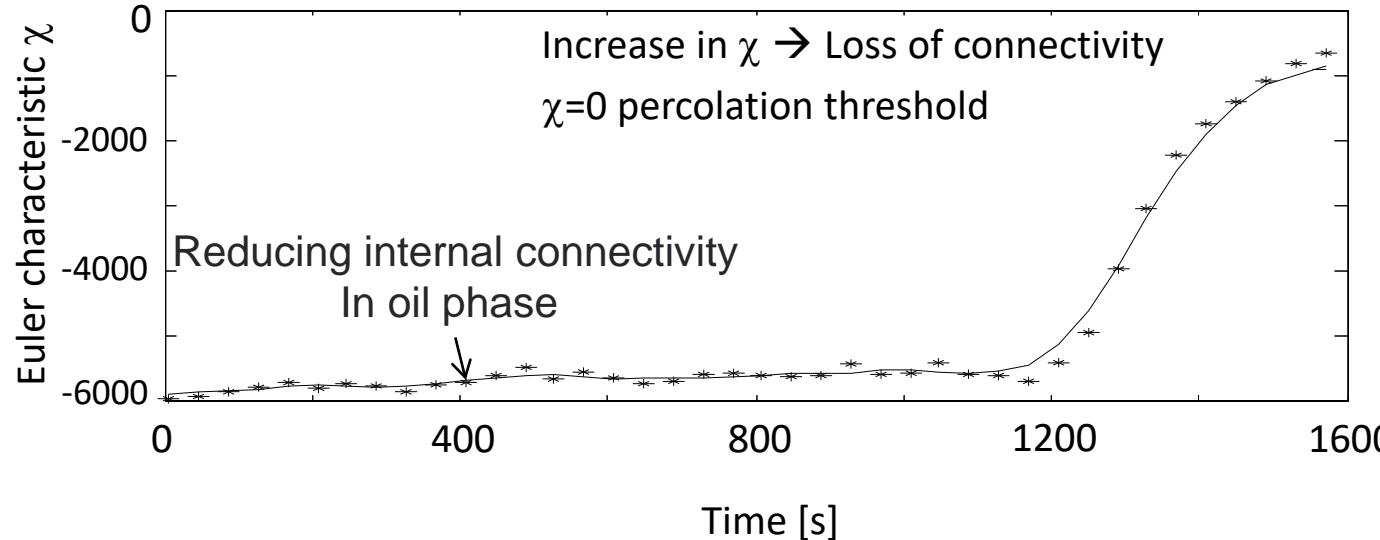
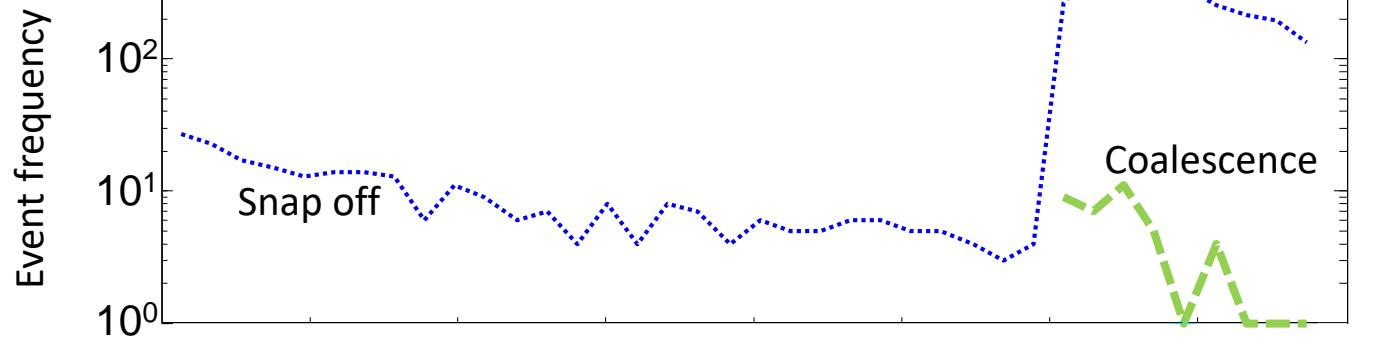
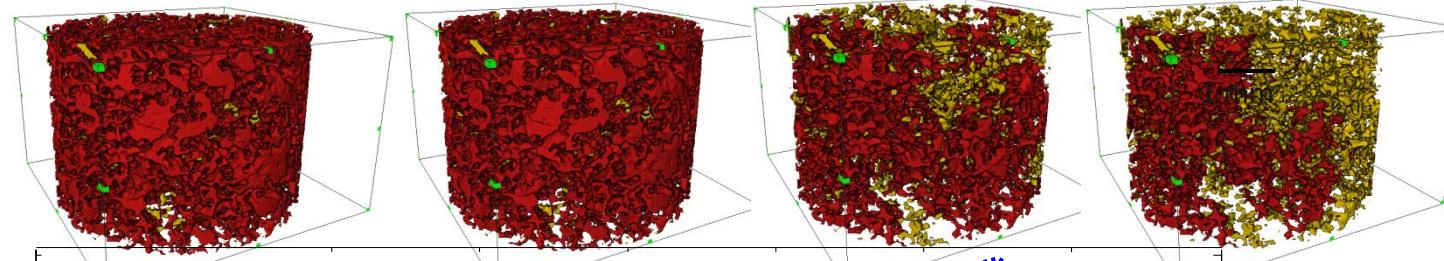
Strongly water-wet



Introduces Topological Changes

Rücker et al. GRL 2015

# Cluster Dynamics Introduces Topological Changes



Rücker et al. GRL, 2015

$$\chi = \text{Objects} - \text{Loops} + \text{Voids}$$

Snap-off

→ disconnection

→ #objects increases

→  $\chi$  increases

Coalescence

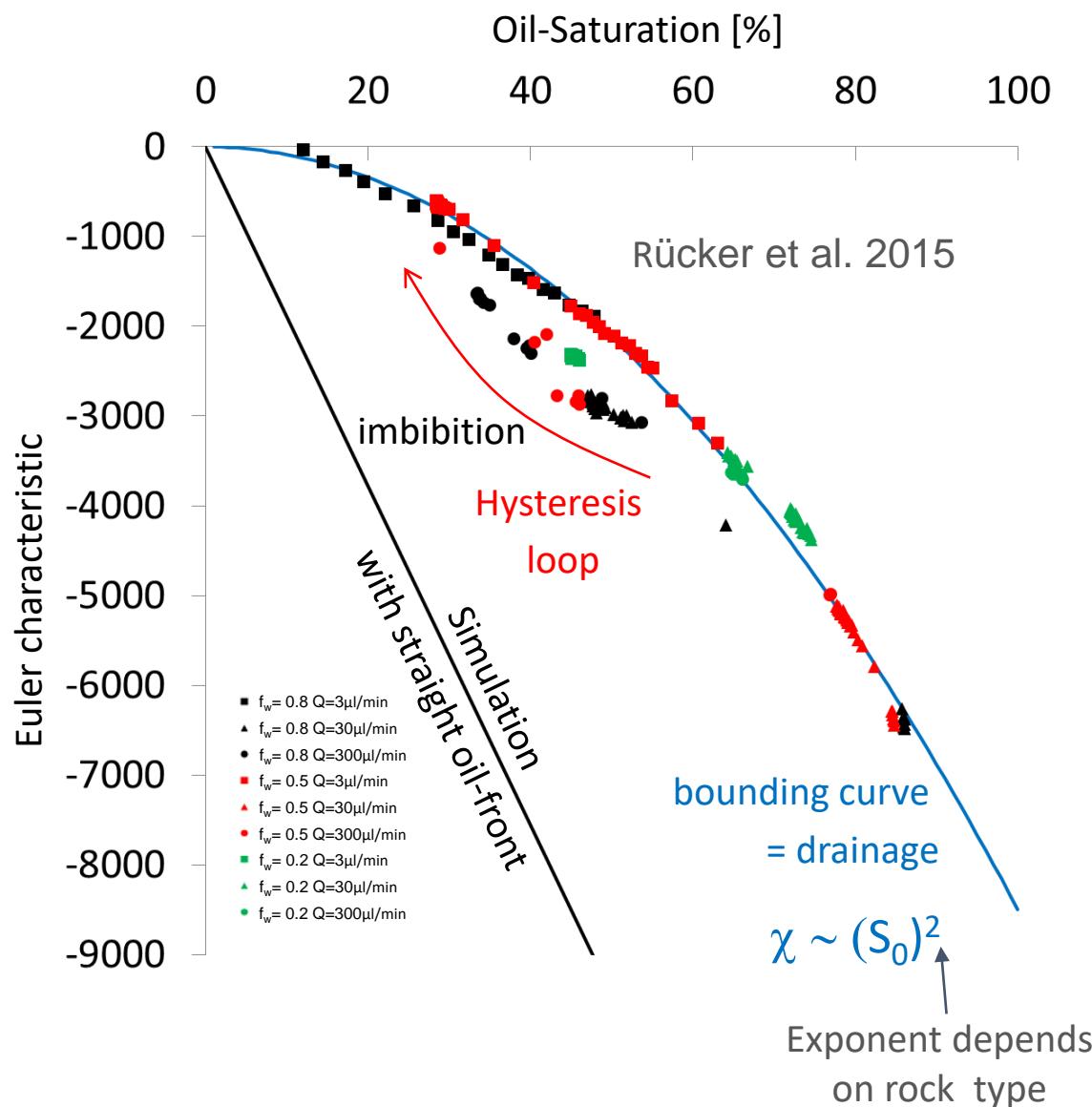
→ connection

→ #objects decreases

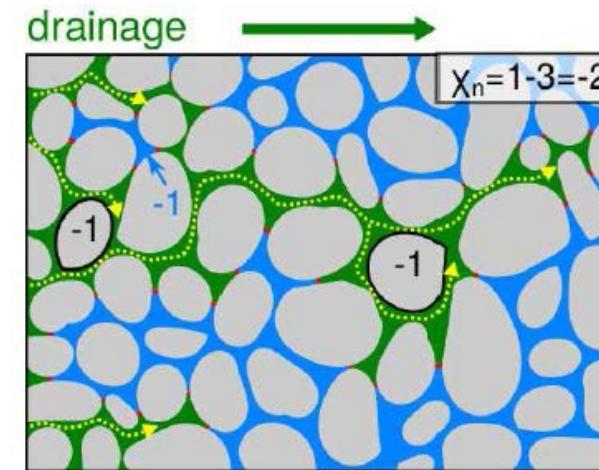
→  $\chi$  decreases

Here: #snap-off > #coalescence  
→ net  $\chi$  increase

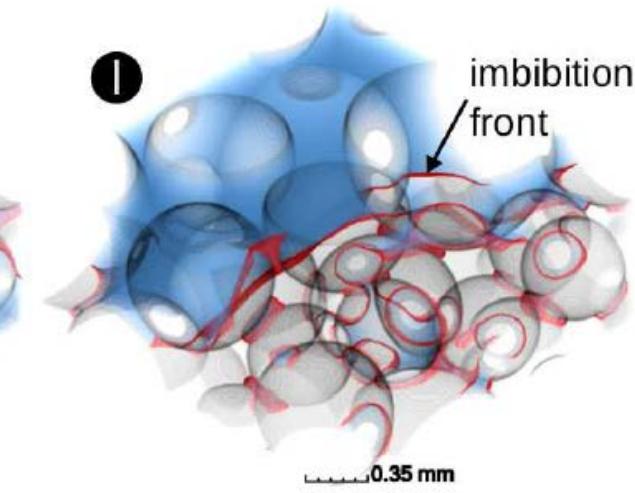
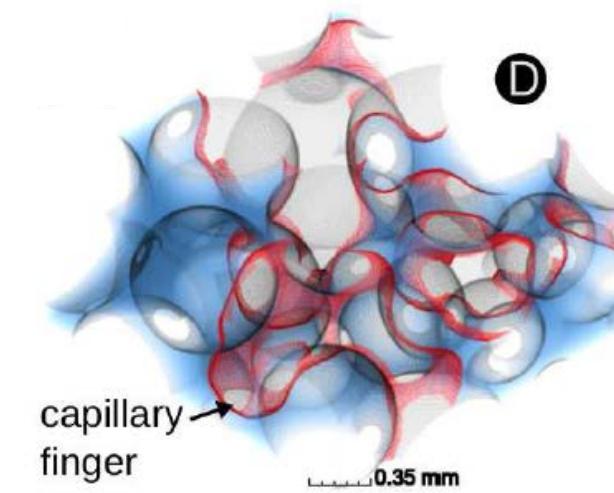
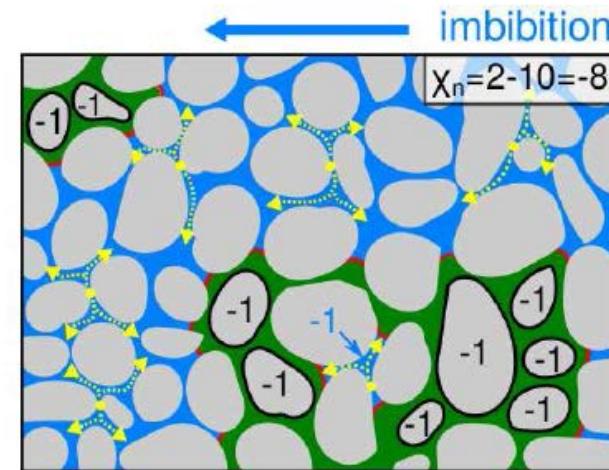
# The Discovery of the 4<sup>th</sup> State Variable



Drainage =  
maintaining connectivity  
(avoid forming loops)



Imbibition =  
Snap-off  $\rightarrow$  formation of  
clusters and loops



# Proof by Direct Numerical Simulation

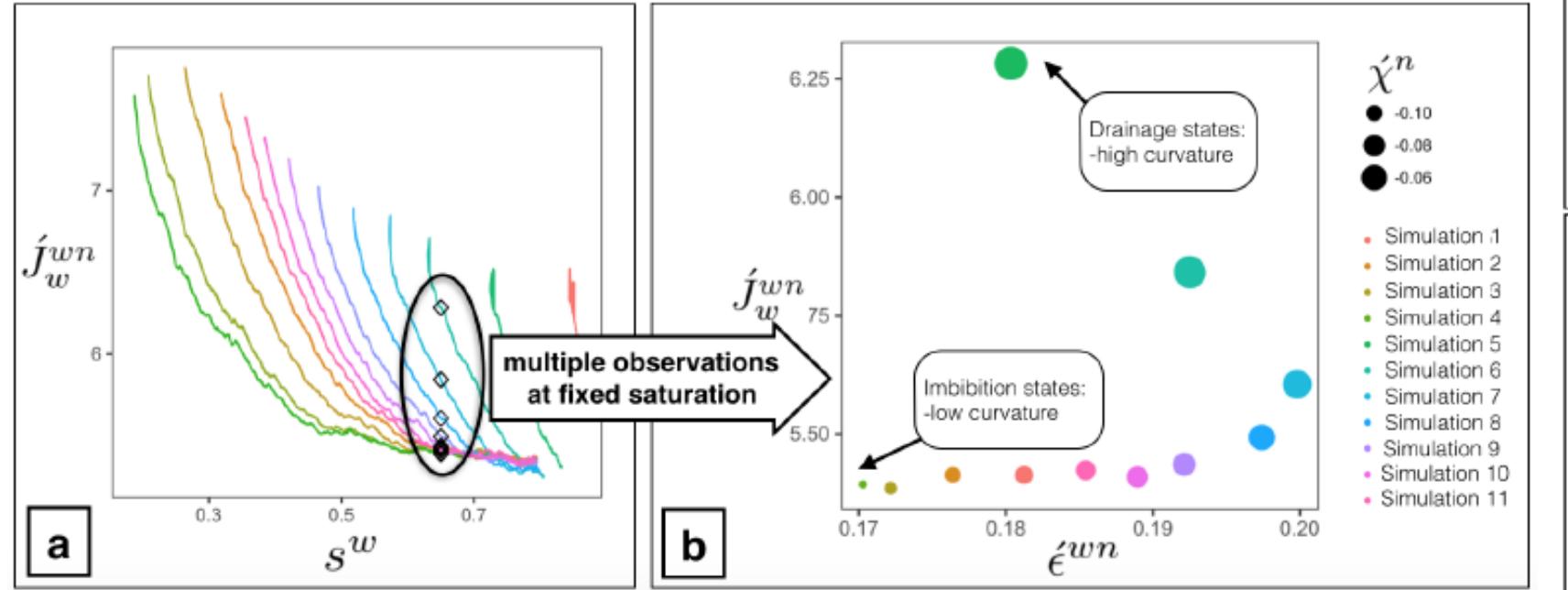
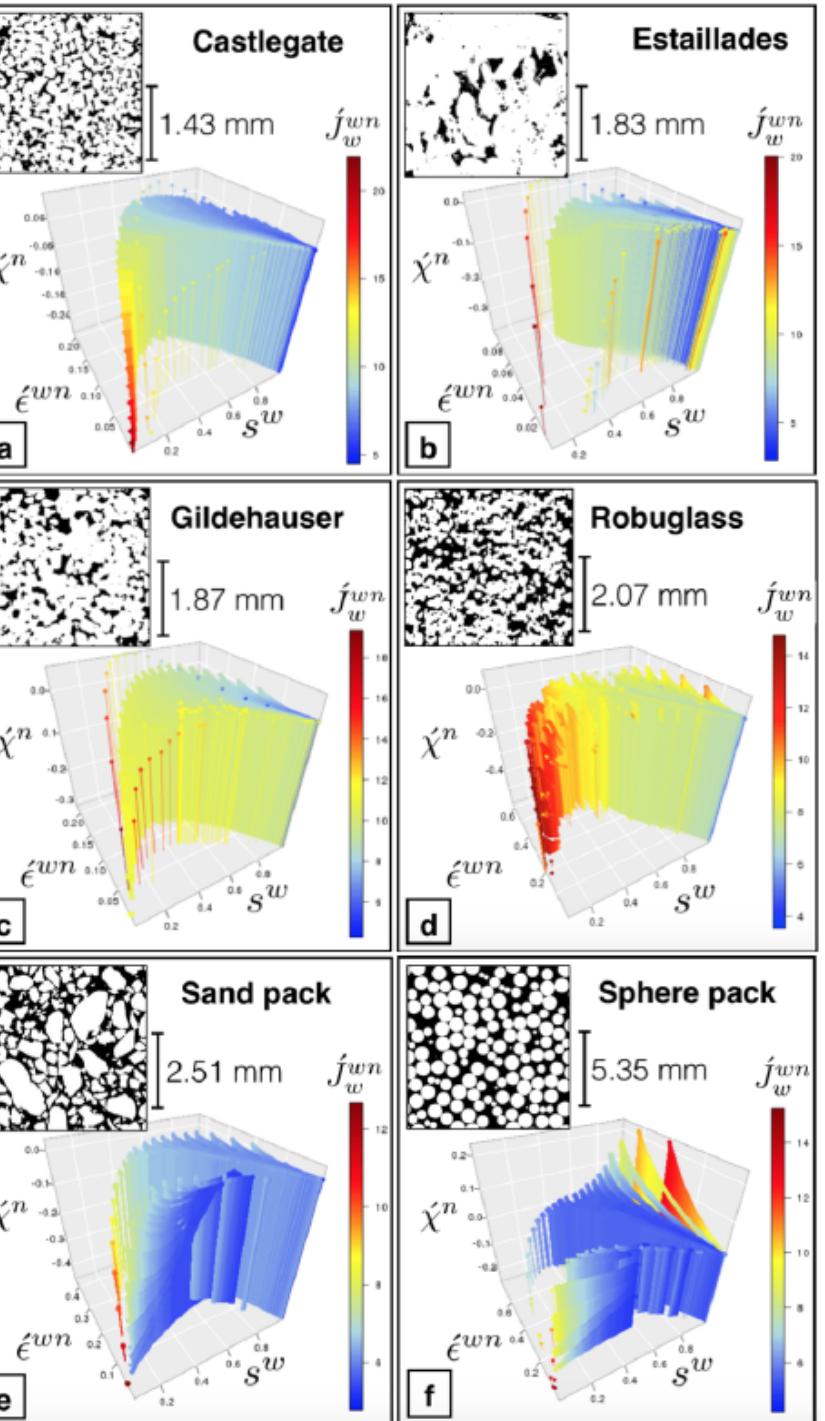


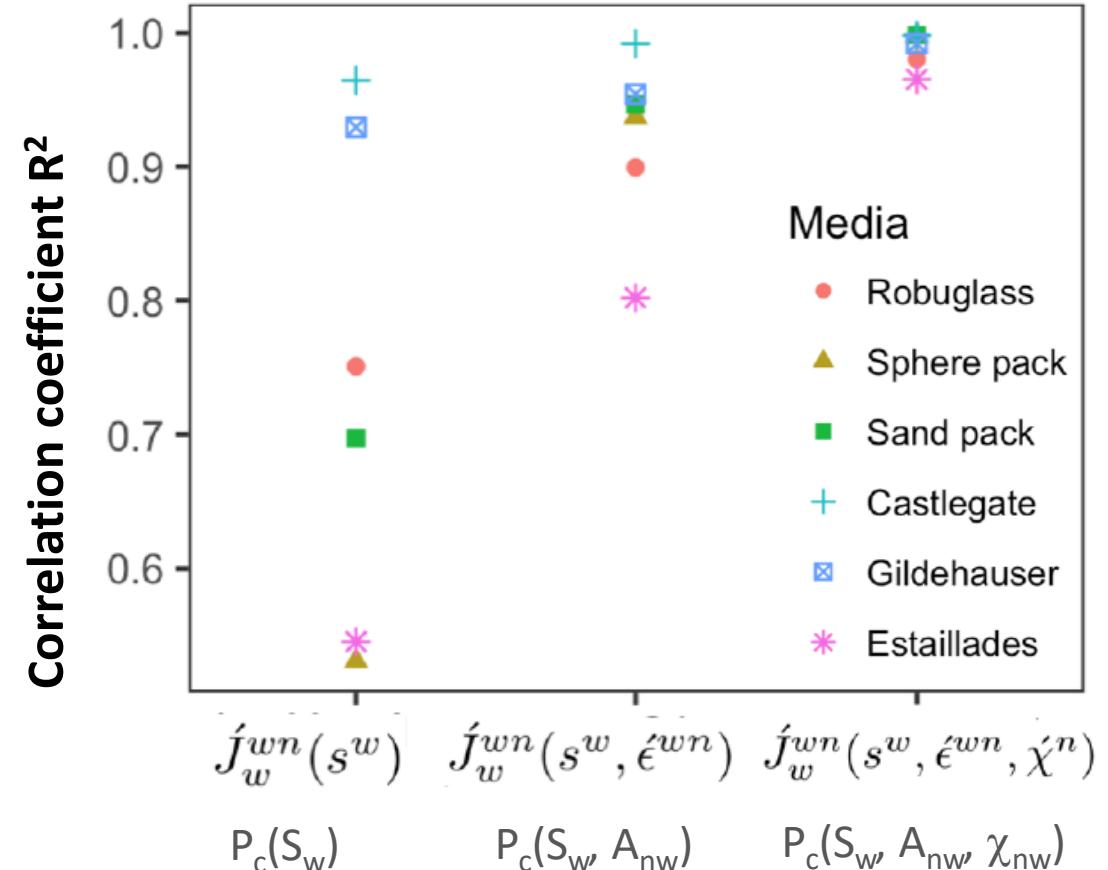
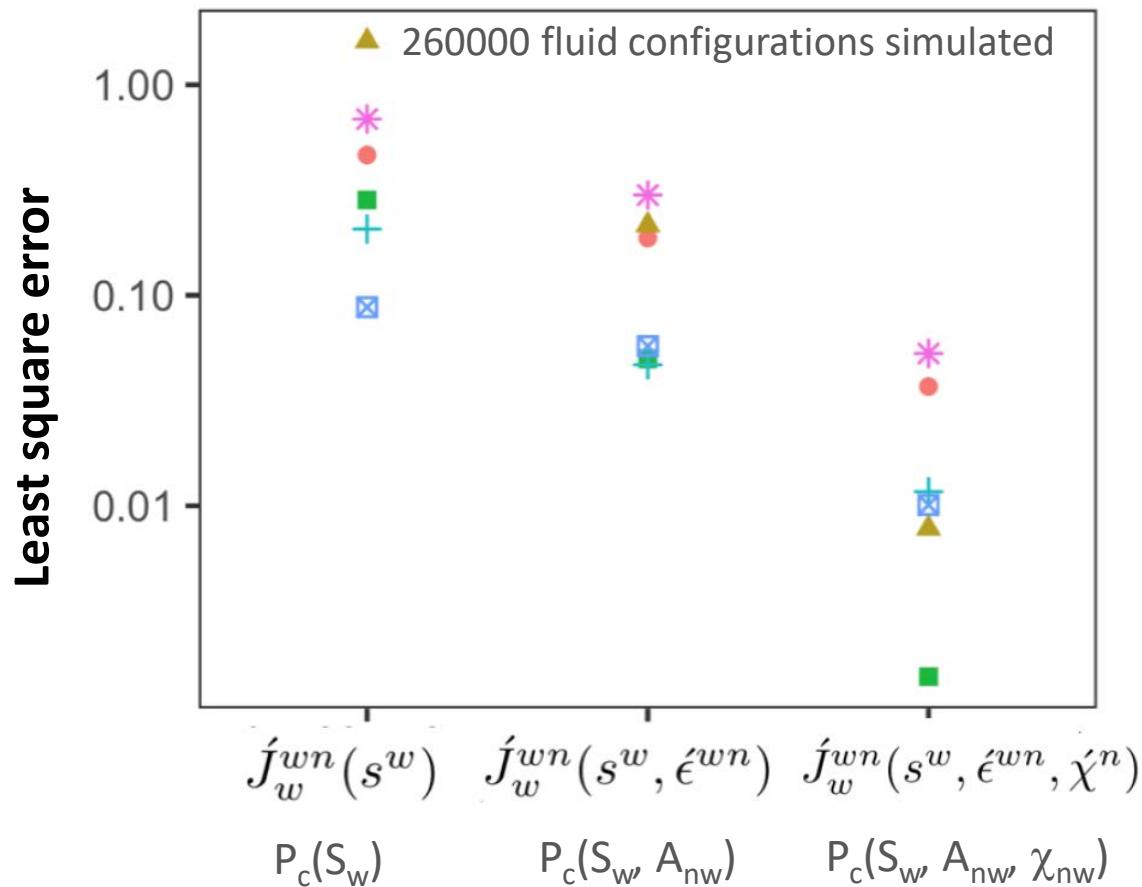
FIG. 6. Traditional models assume that the macroscale capillary pressure is a function only of the saturation of the wetting fluid. Two-fluid displacement simulations within a sand pack show that: (a) At fixed saturation,  $s^w$ , the relative mean curvature  $j_w^{wn}$  can attain many possible values depending on the system history; (b)  $j_w^{wn}(s^w, \epsilon^{wn})$  is non-unique for  $s^w = 0.65$ .



Media	$\epsilon$	$D$ (mm)	Size (voxels)	Sim.	Config.	
Castlegate	0.205	0.111	$512 \times 512 \times 512$	A,D	23,123	
Estaillades	0.111	0.124	$834 \times 834 \times 556$	A,B,D	23,599	
Gildehauser	0.188	0.133	$852 \times 852 \times 569$	A,B,D	38,788	
Robuglass	0.345	0.173	$988 \times 988 \times 598$	A,B,D	49,515	
Sand pack	0.376	0.368	$512 \times 512 \times 512$	A,C,D	64,650	
Sphere pack	0.369	1.00	$900 \times 900 \times 900$	A,C,D	59,341	

# Proof by Direct Numerical Simulation

McClure et al. Phys. Rev. Fluids, 2018

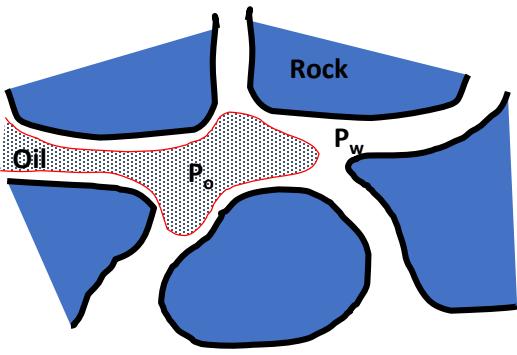


Saturation and Saturation + Interfacial Area are not sufficient to fully parameterize hysteresis (error > 10%)  
→ Saturation + Interfacial Area + Euler Characteristic: error < 10% → full set of state variables

# Upscaling from Pore to the Darcy Scale

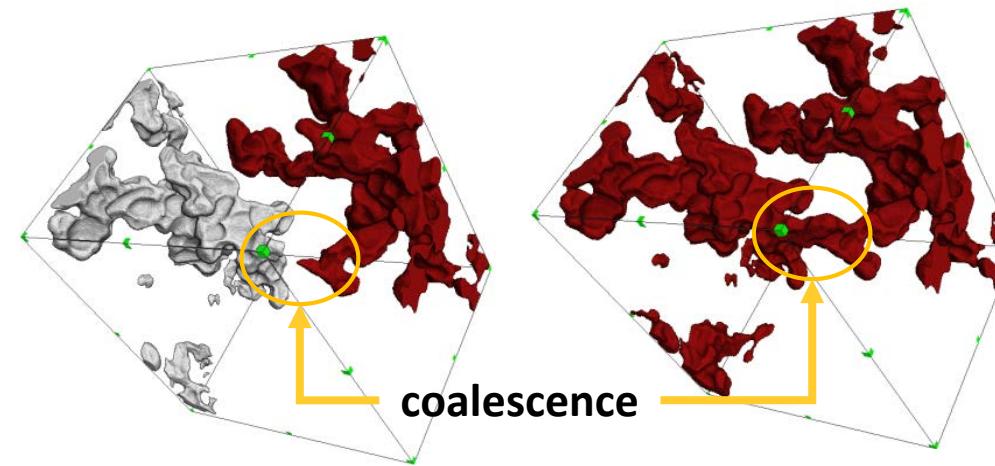
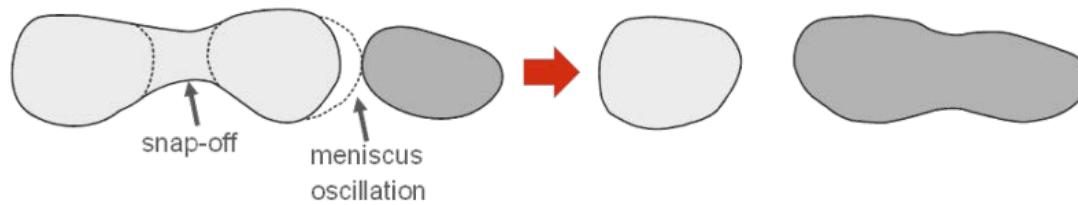
## "Pore scale"

Single pores  
continuum oil & water phases  
Single interfaces



## "Cluster scale"

Non-wetting phase clusters,  
Cooperative dynamics  
Coalescence & breakup  
→ **Changing topology**

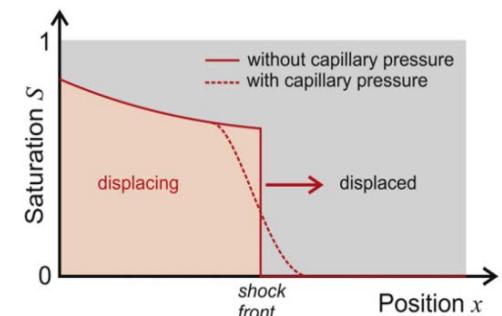


**State variables:  
Minkowski functionals**

## "Darcy scale"

Continuum mechanics:  
porosity, permeability, saturation  
phenomenological description

$$v_i = -k_{r,i} \frac{K}{\mu_i} \frac{dp_i}{dx}$$



cm

m

μm

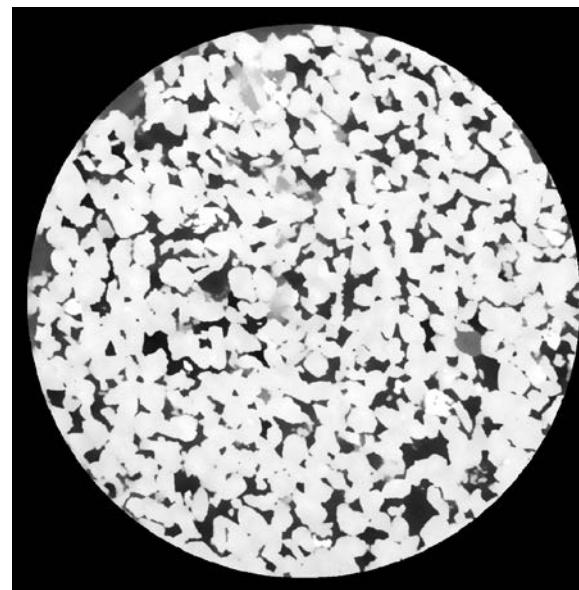
mm

# How to Compute - Software Packages

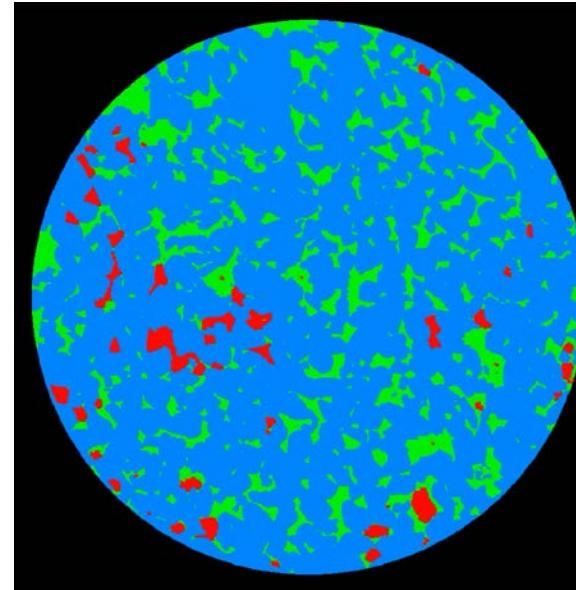
0. Image processing
1. Avizo
2. ImageJ and FIJI: BoneJ plugin
3. Python – scikit-image (currently only 2D)
4. Matlab
5. Quantim ([https://www.ufz.de/export/data/2/94413\\_quantim4\\_ref\\_manual.pdf](https://www.ufz.de/export/data/2/94413_quantim4_ref_manual.pdf))
6. Boundary and connectivity issues
7. Dragonfly/DeepRocks – compute  $\chi$  on network

# Image Processing

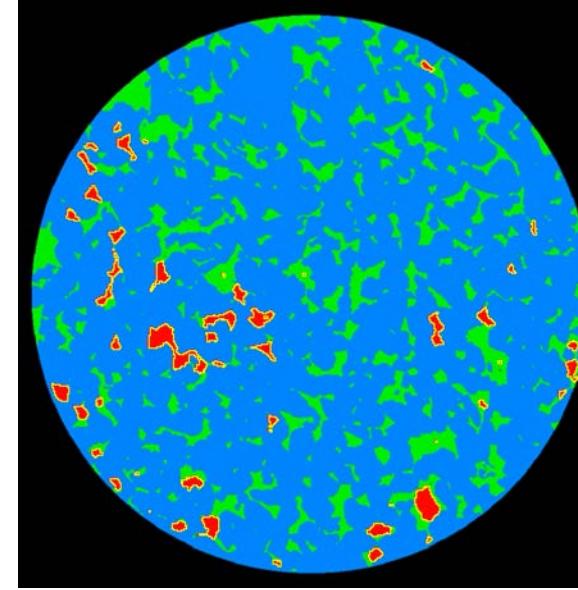
$\mu\text{CT}$  image  $\xrightarrow{\text{segmentation}}$  Segmented  $\xrightarrow{\text{select phase}}$  Oil Phase  $\xrightarrow{\text{object analysis}}$  Oil clusters



16 bit grey level  
(potentially filter)

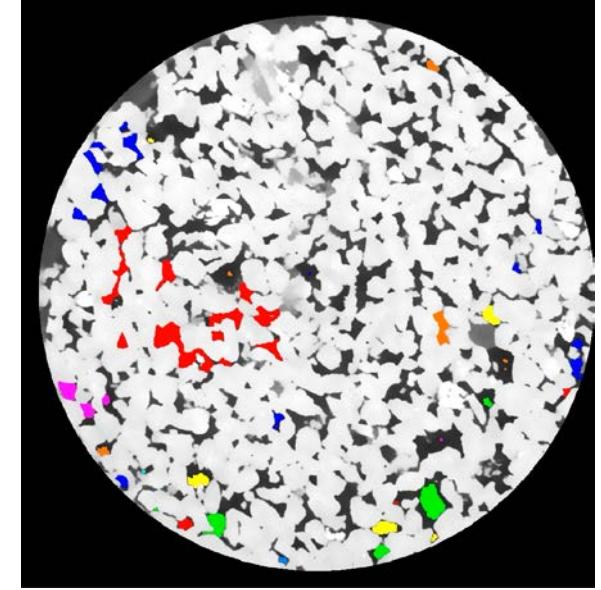


Segmented image  
Blue = rock  
Red = oil  
Green = gas



For each **phase** compute

- Volume
- Interfacial area
- Mean curvature
- Euler characteristic



For each **cluster** compute

# Avizo: Mean Curvature

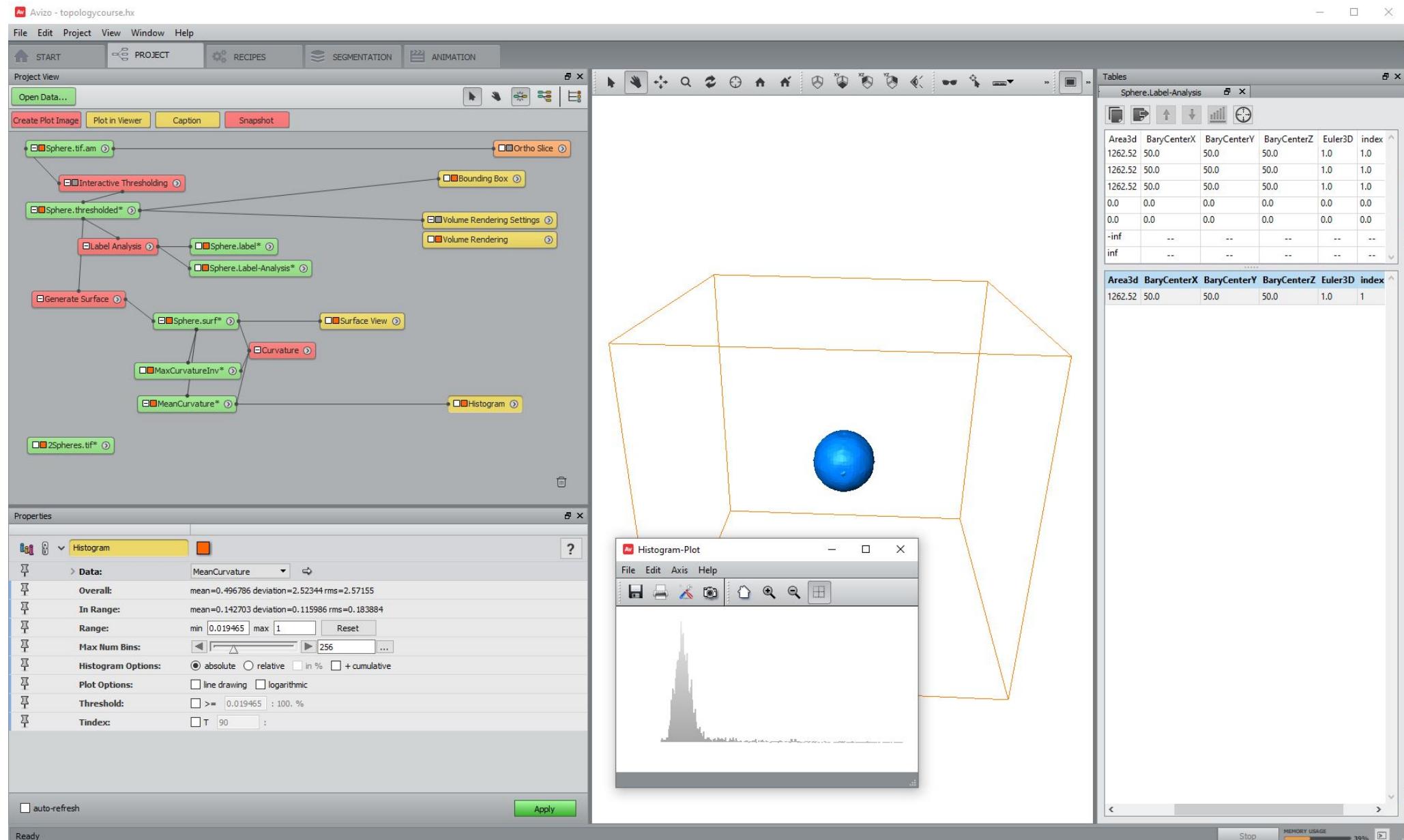
segmentation

→ Label analysis

- Volume
- Area
- Euler char.

→ Generate surface  
(spline interpolation)

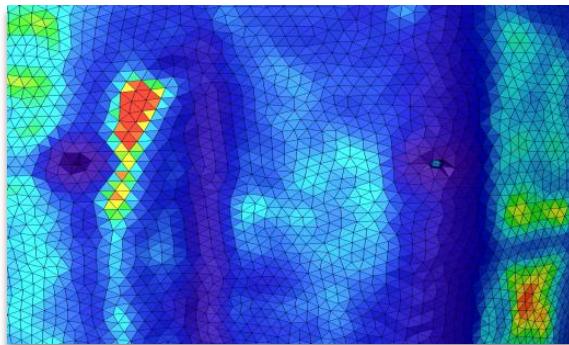
- Mean curvature
- Gaussian curvature



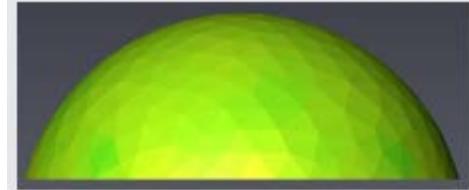
# Avizo : Surface meshes

- Surface property calculation

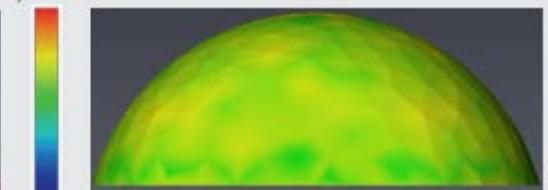
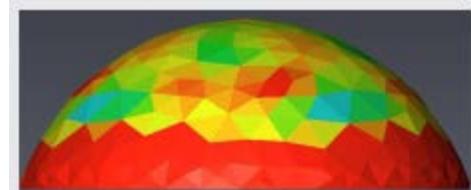
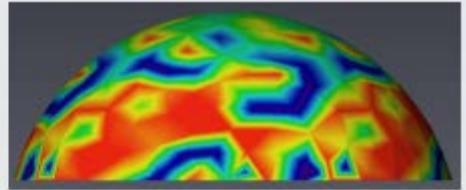
- Curvature
- Distance
- Roughness
- Thickness



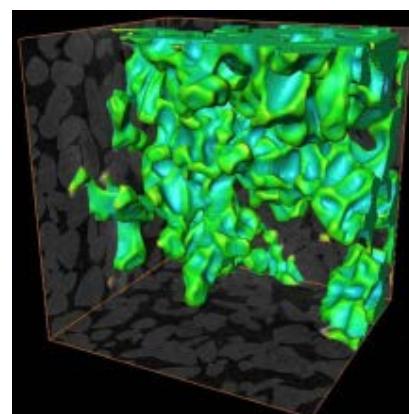
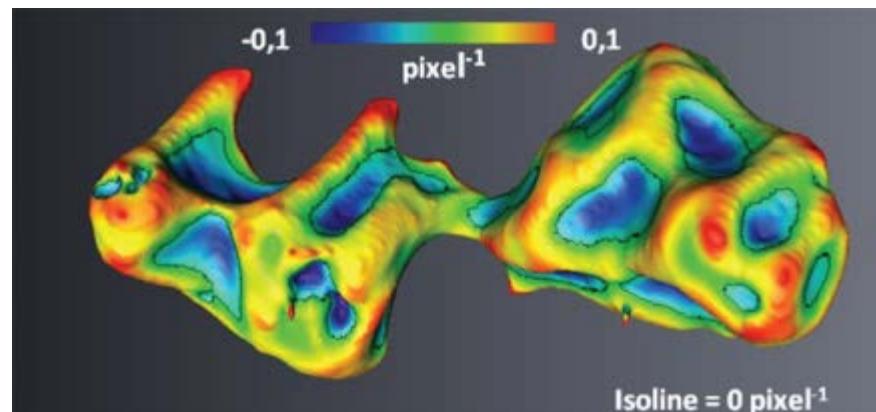
Surface-based Approach:



Voxel-based Approach:

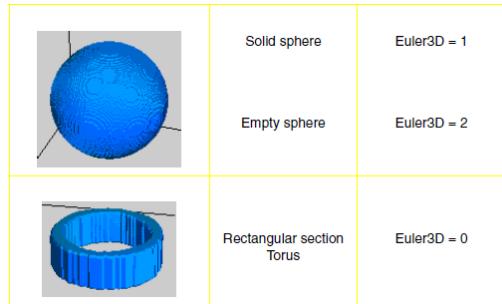


SCA2012-55



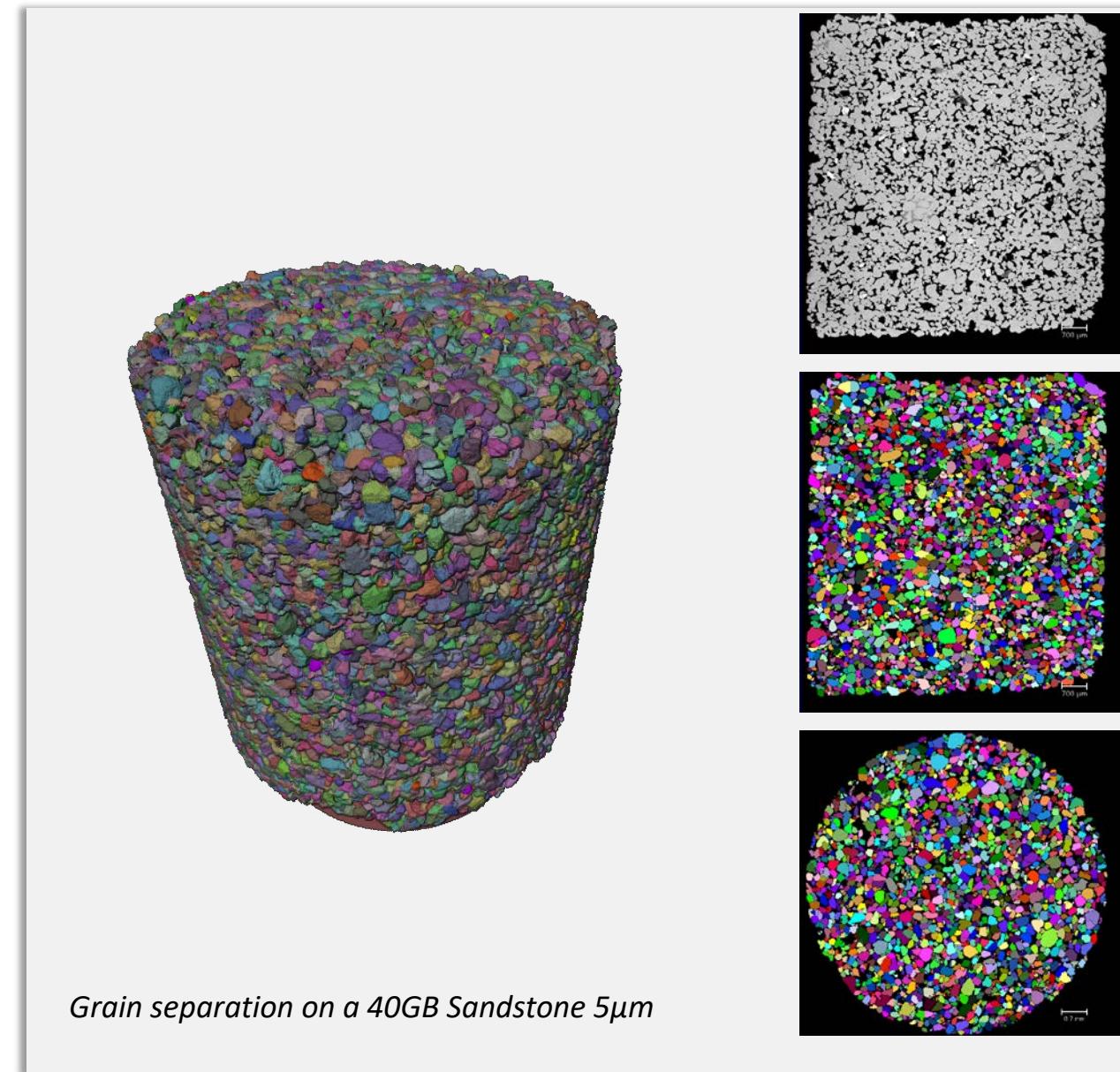
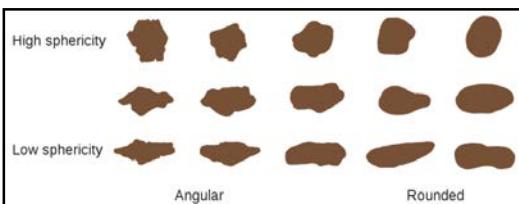
# Avizo : Object Analysis

- Analysis
  - Shape factor
  - Equivalent diameter
  - **Volume, area**
  - **Euler**
  - Orientation
  - Roundness
  - Sphericity
  - Rugosity
  - Crofton

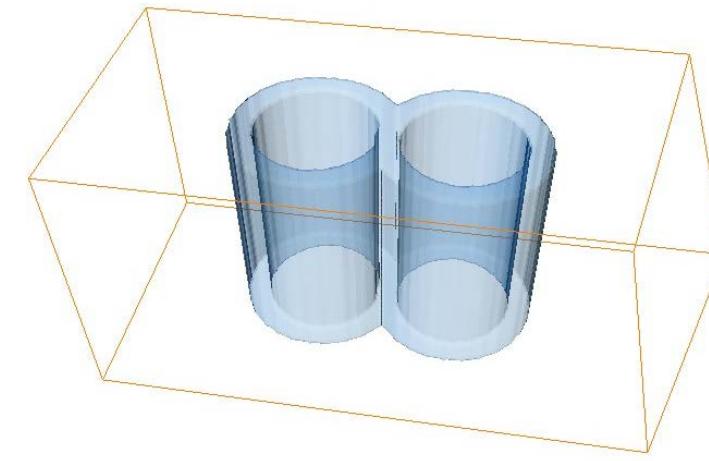
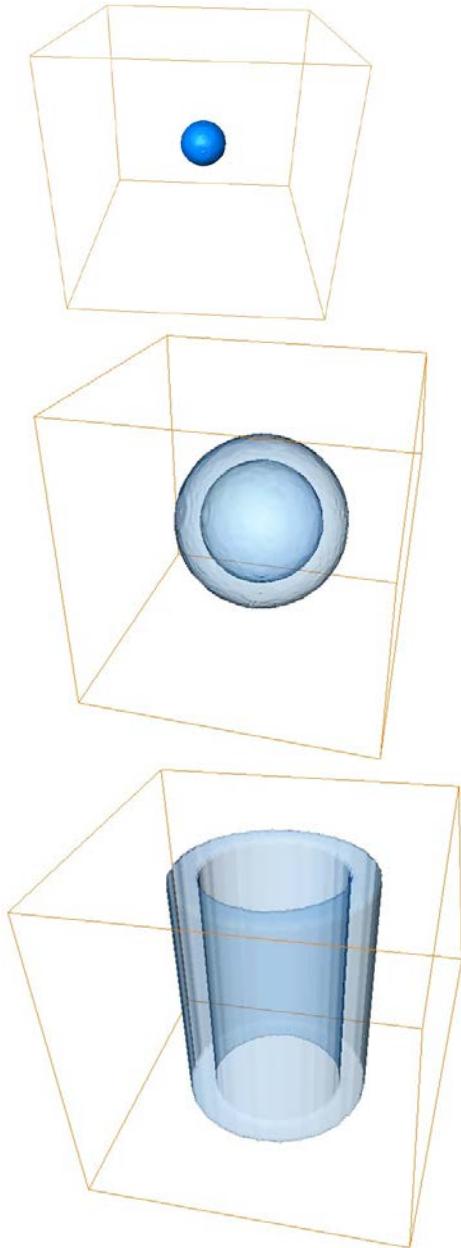


Name	Picture	Volume	Area	Sphericity
<b>Platonic Solids</b>				
tetrahedron		$\frac{\sqrt{2}}{12} s^3$	$\sqrt{3} s^2$	$\left(\frac{\pi}{6\sqrt{3}}\right)^{\frac{1}{2}} \approx 0.671$
cube (hexahedron)		$s^3$	$6 s^2$	$\left(\frac{\pi}{6}\right)^{\frac{1}{2}} \approx 0.806$
octahedron		$\frac{1}{3} \sqrt{2} s^3$	$2\sqrt{3} s^2$	$\left(\frac{\pi}{3\sqrt{3}}\right)^{\frac{1}{2}} \approx 0.846$
dodecahedron		$\frac{1}{4} (15 + 7\sqrt{5}) s^3$	$3\sqrt{25 + 10\sqrt{5}} s^2$	$\left(\frac{(15 + 7\sqrt{5})^2 \pi}{12(25 + 10\sqrt{5})}\right)^{\frac{1}{2}} \approx 0.910$
icosahedron		$\frac{5}{12} (3 + \sqrt{5}) s^3$	$5\sqrt{3} s^2$	$\left(\frac{(3 + \sqrt{5})^2 \pi}{60\sqrt{3}}\right)^{\frac{1}{2}} \approx 0.939$

- Classification

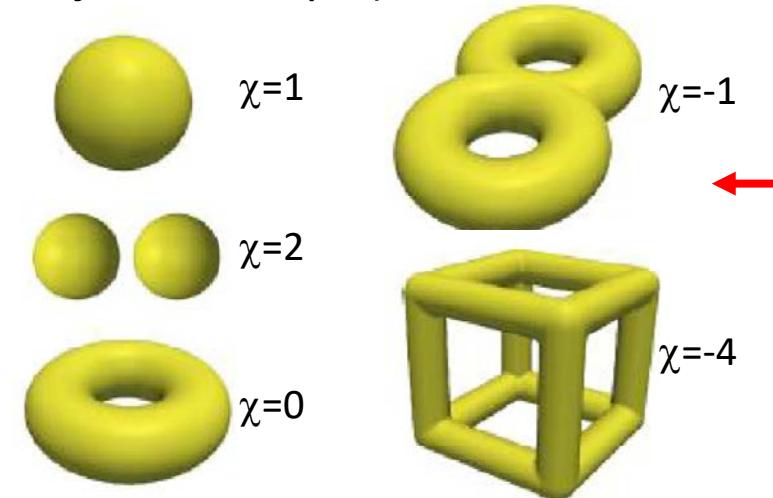


# Avizo: Euler Characteristic



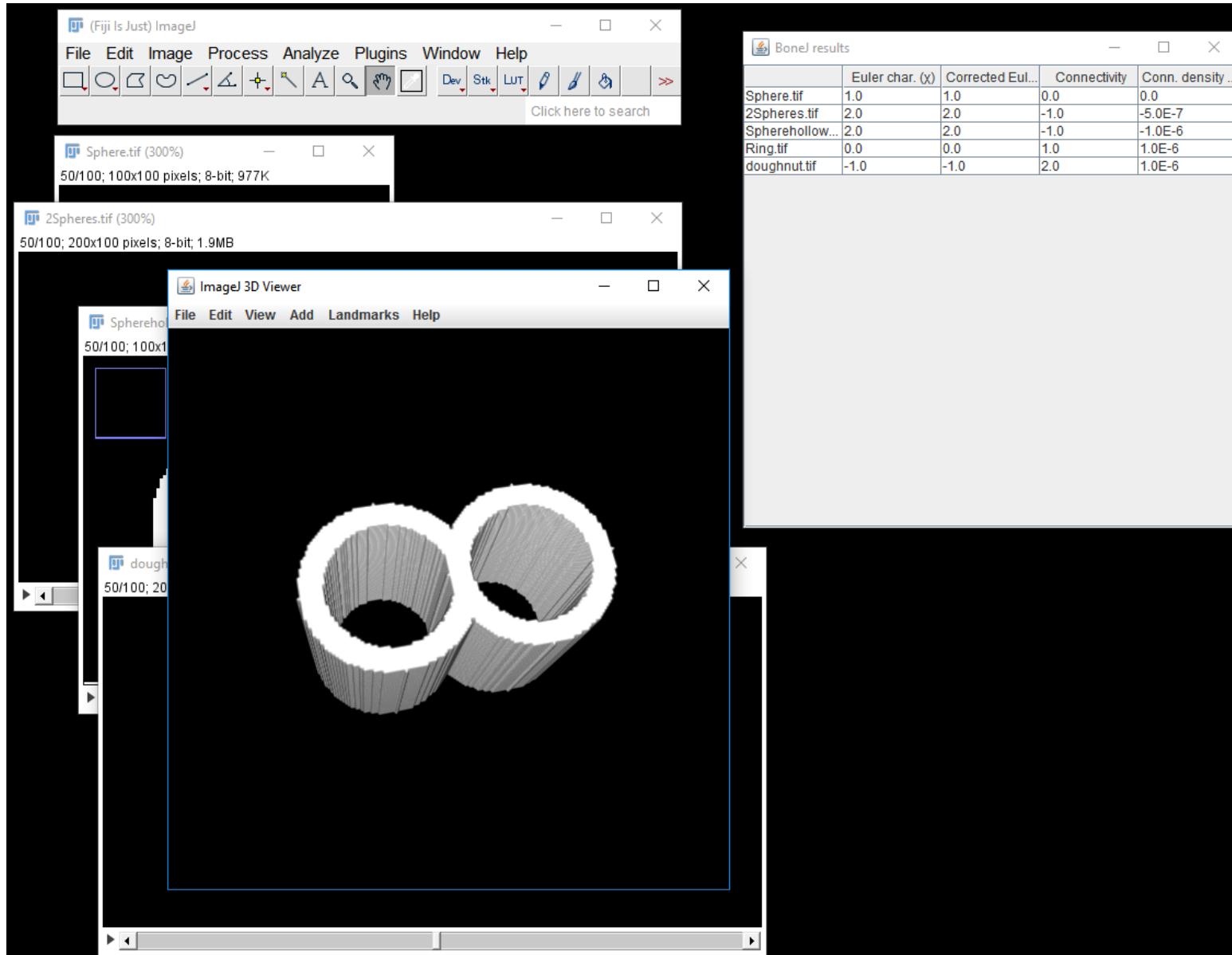
	Volume	Area	BarX	Bary	Barz	Euler
Sphere	4169	1262.52	50	50	50	1
2spheres	4169	1262.52	50	50	50	1
	4169	1262.52	150	50	50	1
sum	4170	1263.52				2
hollow sphere	72982	16953.9	50	50	50	2
ring	110284	27631.2	50	50	50	0
doughnut	203899	49396.3	99.9872	50	50	-1

Euler Characteristic  $\chi =$   
Objects – Loops (+ Inclusions)

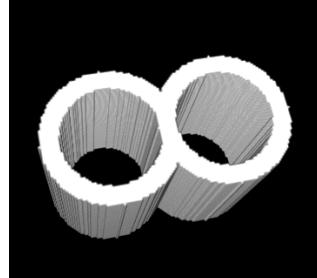
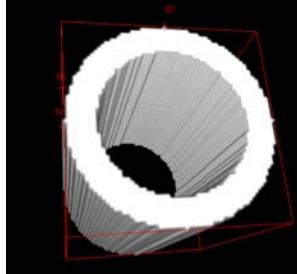
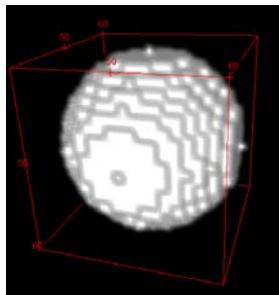
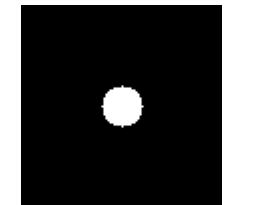


# ImageJ and FIJI: BoneJ plugin (OpenSource)

ImageJ - <https://imagej.nih.gov/ij/>  
FIJI - <https://fiji.sc/>  
BoneJ - <http://bonej.org/>



# ImageJ and FIJI: BoneJ plugin (OpenSource)



	Euler char. ( $\chi$ )	Corrected Euler ( $\chi + \dots$ )	Connectivity	Conn. density (pixel $^3$ )	Surface area (pixel $^2$ )
Sphere.tif	1.0	1.0	0.0	0.0	1263.9727775132994
2Spheres.tif	2.0	2.0	-1.0	-5.0E-7	2527.9455550266375
Spherehollow.tif	2.0	2.0	-1.0	-1.0E-6	18358.24357045593
Ring.tif	0.0	0.0	1.0	1.0E-6	29282.273575259434
doughnut.tif	-1.0	-1.0	2.0	1.0E-6	52570.373993201676

Euler Characteristic  $\chi =$   
Objects – Loops (+ Inclusions)



$\chi=1$



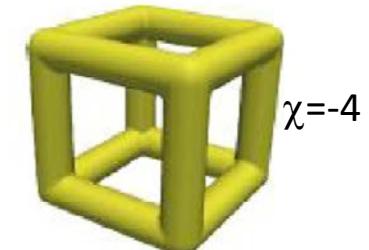
$\chi=-1$



$\chi=2$



$\chi=0$



$\chi=-4$

# Python

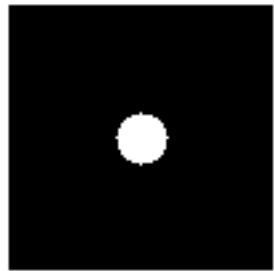


scikit-image  
image processing in python

<https://scikit-image.org/docs/dev/api/skimage.measure.html>

Many functions unfortunately currently only in 2D !

Euler = 1



Euler = 0



Euler = -1



Euler Characteristic  $\chi$  =  
Objects – Loops (+ Inclusions)



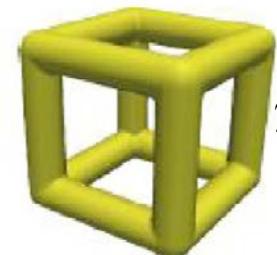
$\chi=1$



$\chi=-1$



$\chi=2$

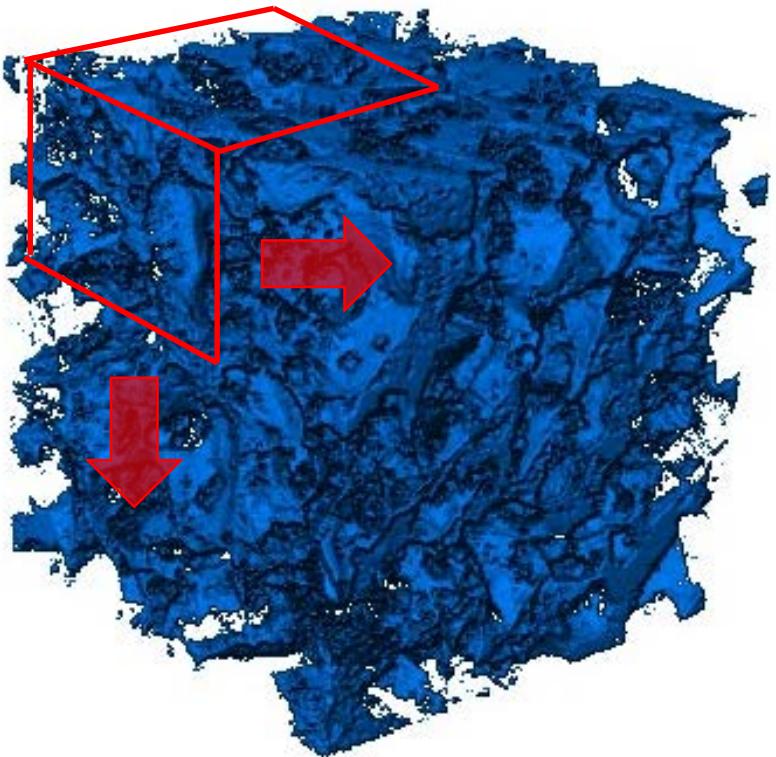


$\chi=-4$

Herring, AWR 2013

```
17 import matplotlib.pyplot as plt
18 from skimage import io
19 from skimage.measure import label, regionprops
20
21 filenamelist=['Sphere.tif','Ring.tif','doughnut.tif']
22
23
24 fig = plt.figure()
25 i=1
26 for imagename in filenamelist:
27     image = io.imread(imagename)
28     xdim, ydim, zdim = image.shape
29     label_img = label(image[int(xdim/2),:,:])
30     regions = regionprops(label_img)
31     for props in regions:
32         el = props.euler_number
33 #         ar = props.area
34     ax = fig.add_subplot(1,3,i)
35     ax.imshow(image[int(xdim/2),:,:], cmap=plt.cm.gray)
36     ax.set_title('Euler = '+str(el), loc="left")
37     ax.set_xticks(())
38     ax.set_yticks(())
39     ax.axis('off')
40     i=i+1
41
42 plt.show()
```

# MATLAB



<https://github.com/mattools/matImage>

Legland, D.; Kiêu, K. & Devaux, M.-F. Computation of Minkowski measures on 2D and 3D binary images.  
Image Anal. Stereol., 2007, 26, 83-92

```
for w_idx = [900]      % ROI size  
    w_surface = [];  
    w_labels = [];  
    i_loc = ((mask_x - w_idx)/2)+1;  
    j_loc = ((mask_y - w_idx)/2)+1;  
    k_loc = ((mask_z - w_idx)/2)+1;  
    roi = FinallImage(i_loc:i_loc+w_idx-1, j_loc:j_loc+w_idx-1,  
                      k_loc:k_loc+w_idx-1);  
    [Euler_roi, Euler_labels]= imEuler3d(roi);  
    % Implements Euler number codes  
    w_Euler = vertcat(w_Euler, Euler_roi);  
    w_labels = vertcat(w_labels, Euler_labels);  
    file_name = [num2str(w_idx), 'Filename_900.mat'];  
    save(file_name, 'w_Euler', 'w_labels')  
end
```

# Boundary and Connectivity Issues

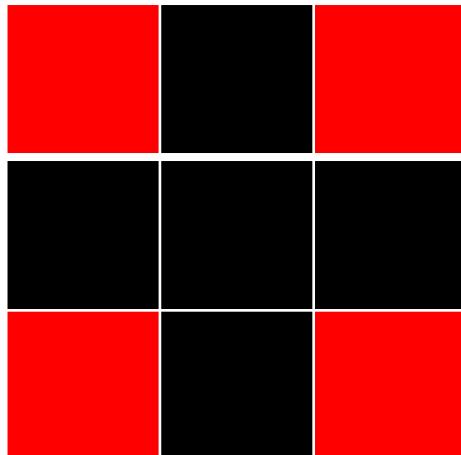
ROI size	Euler3D calculated by Avizo	Euler3D calculated with MATLAB codes
200	-5.07E+02	-5.62E+02
300	-2.70E+03	-2.69E+03
400	-6.76E+03	-6.80E+03

Do these two objects connect?

YES → Euler = 1

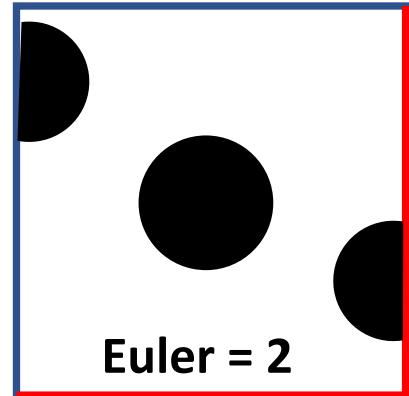
No → Euler = 2

Pixel Connectivity

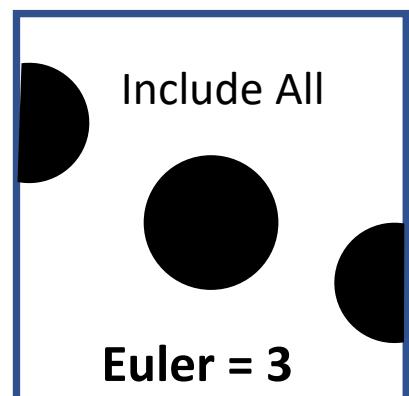


## Boundary Issues

Include



Exclude

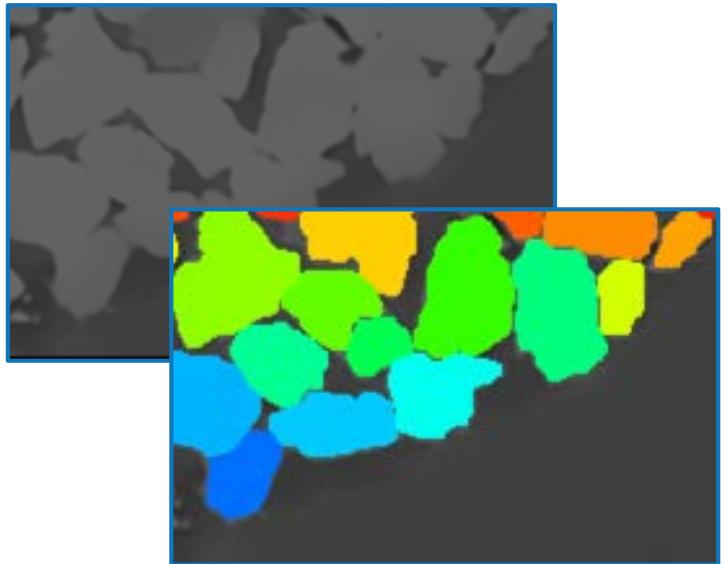


Include All

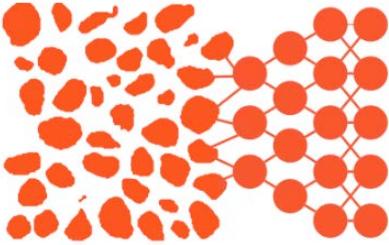
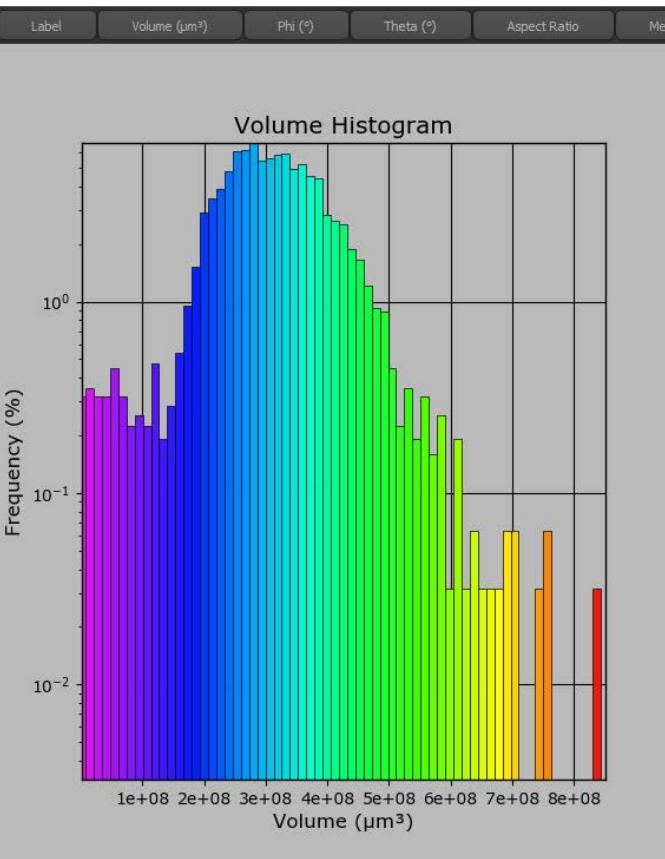
# Measurements on Segmented Voxels

## Color and Sort by measurement

- Volume
- Surface area
- Aspect Ratio
- Phi
- Theta
- Position (x,y,z)
- Intensity (mean, min, max, stdev)



Grain separation and labeling by Deep Watershed

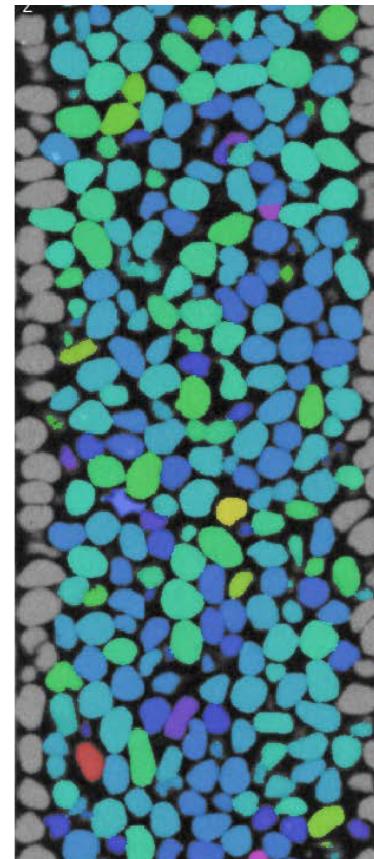


DEEP  
ROCKS

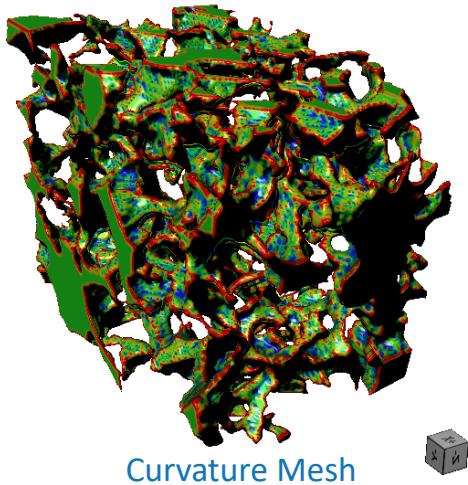
Whole Core  
Medical CT, Core photography

Plugs and Cuttings  
Micro-CT, Nano-CT, FIB-SEM

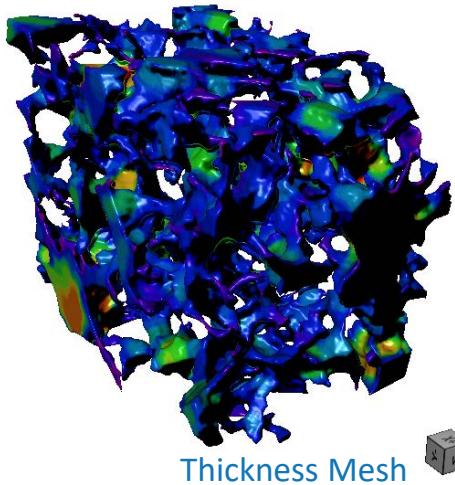
Thin Section  
LM, SEM, TEM, CL, Mineralogy



Courtesy of  
TheObjects



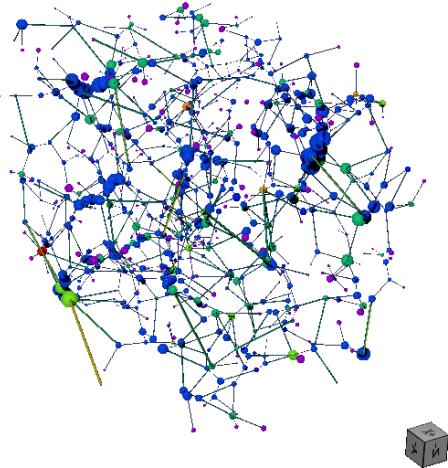
Curvature Mesh



Thickness Mesh

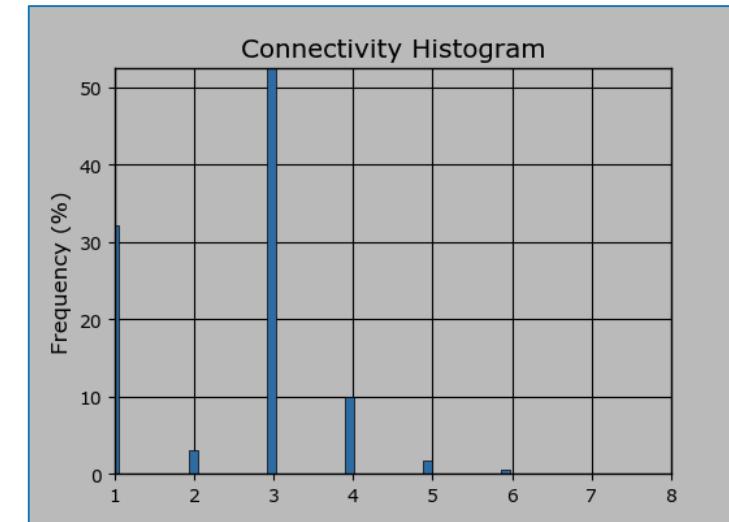
### Mesh operations

- Smoothing
- Decimation
- Thickness
- **Mean Curvature**
- **Gaussian Curvature**



### Pore network

- Nodes scaled by pore body radius
- Nodes colored by connectivity index
- Edges colored by edge length
- **$\chi = -212$  (Euler characteristic)**



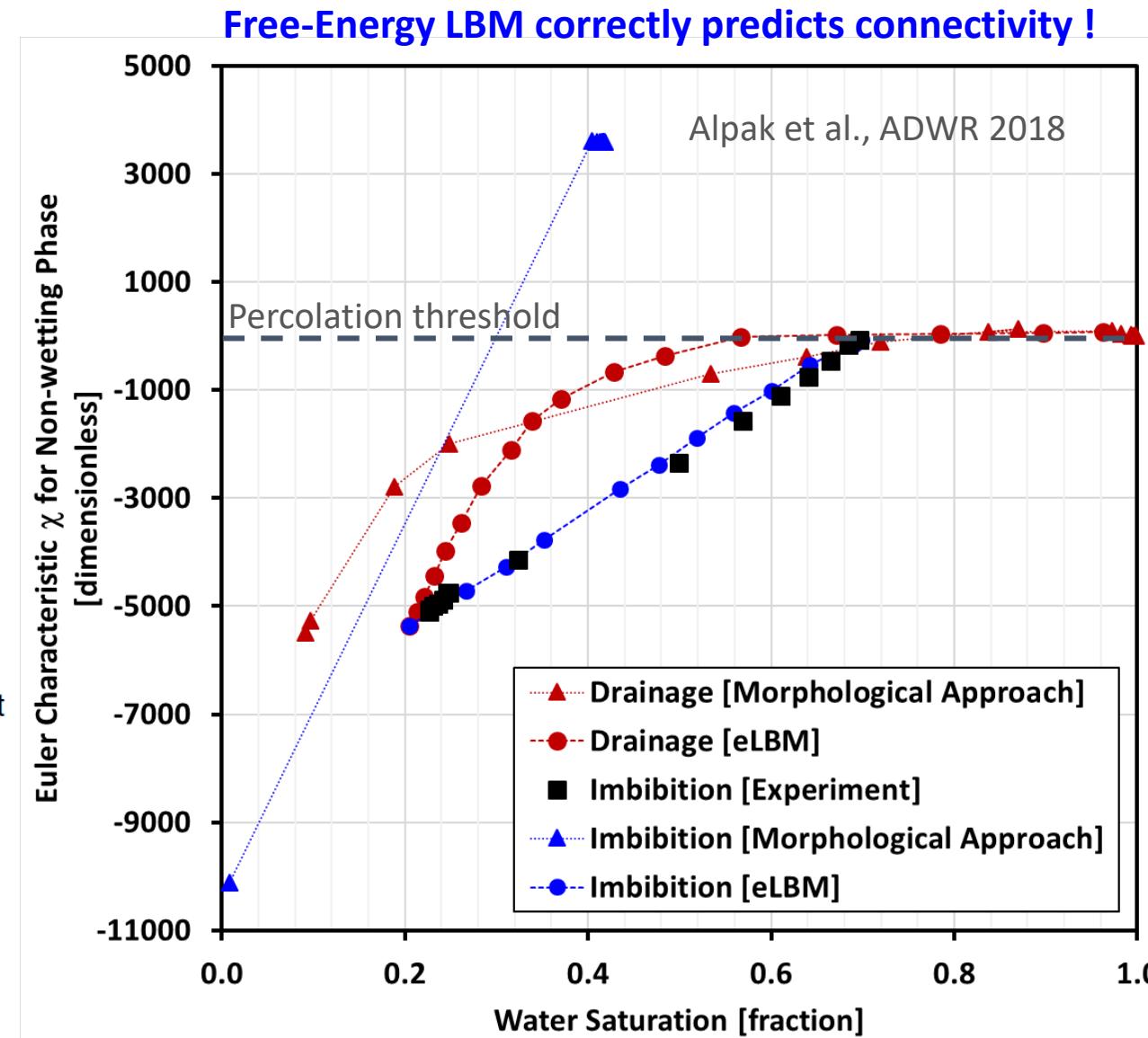
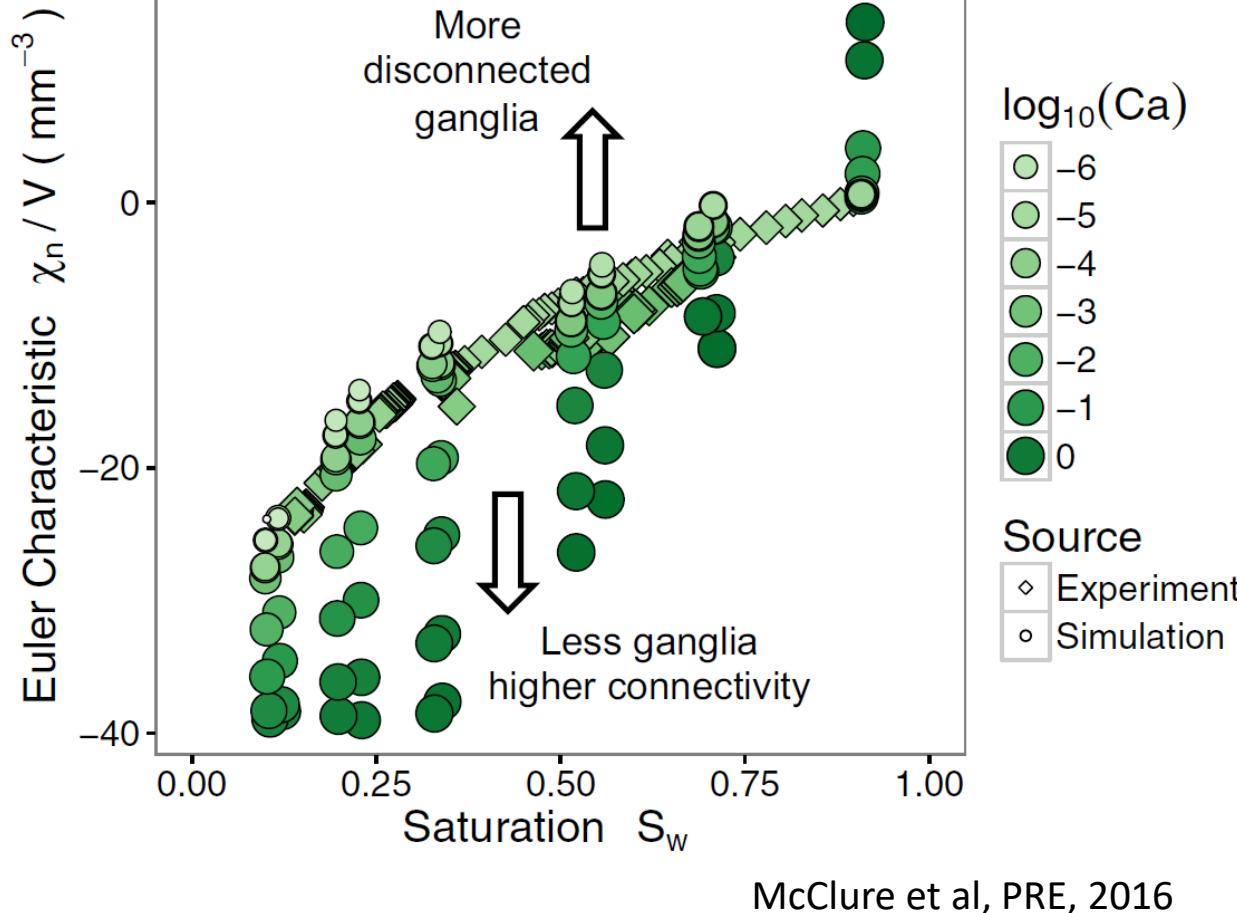
Courtesy of TheObjects

# Applications

1. Digital Rock: Validation of pore scale simulators: relative permeability
2. new route to relative permeability
3. Phase connectivity inferred from resistivity measurements
4. Phase connectivity in critical gas saturation
5. Hysteresis models for Darcy-scale flow
6. Wettability: Description of contact angle as deficit curvature
7. Wettability: bi-continuous interfaces in intermediate/mixed-wet rock
8. Petrology mineral analysis for interpreting relative permeability
9. Permeability in fracture networks

# Application: Validation of Pore Scale Simulation

LBM fractional flow simulation: match with exp. data

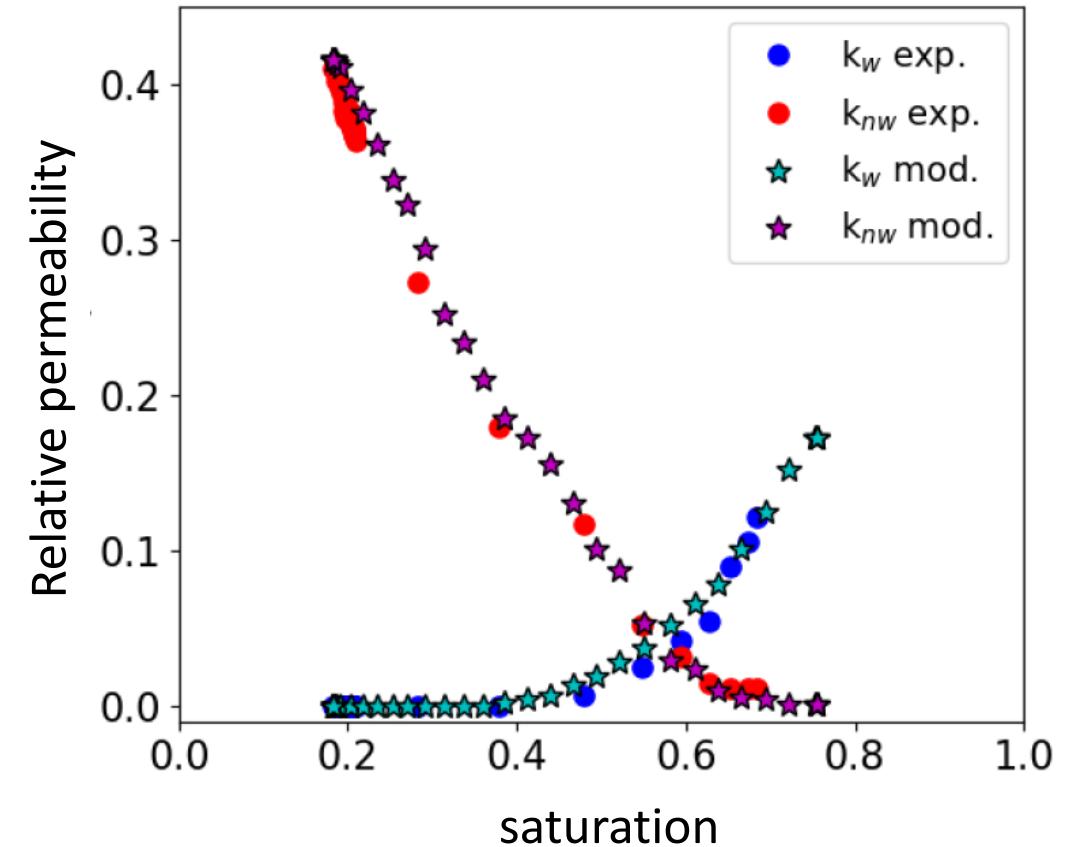
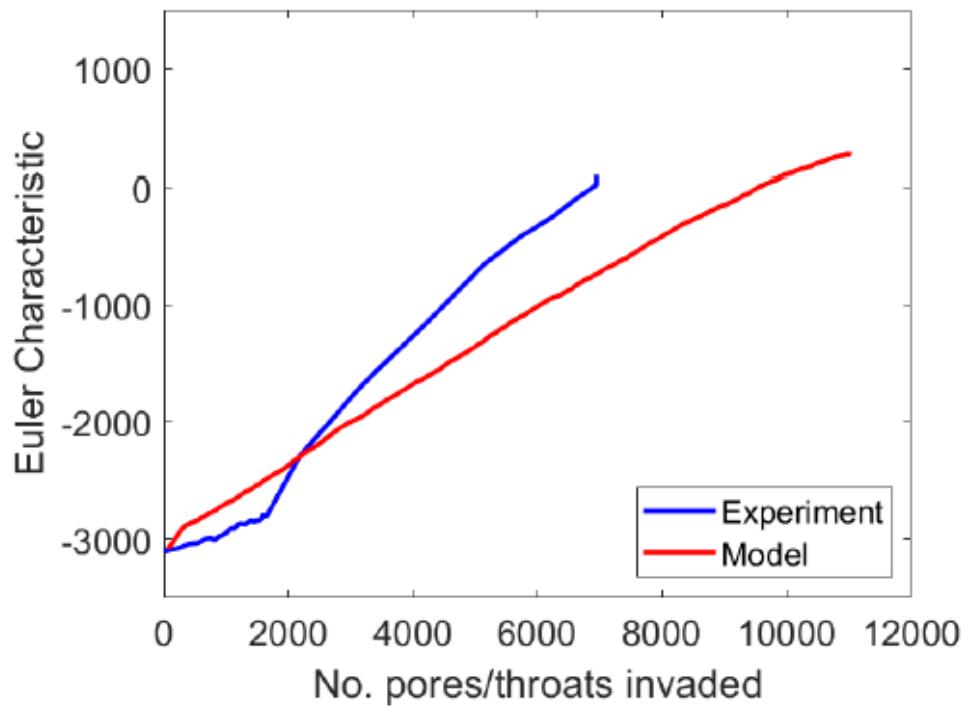


# Application: Validation of Pore Scale Simulation

Model: quasi-static Pore Network Model (Ruspini et al. 2018)

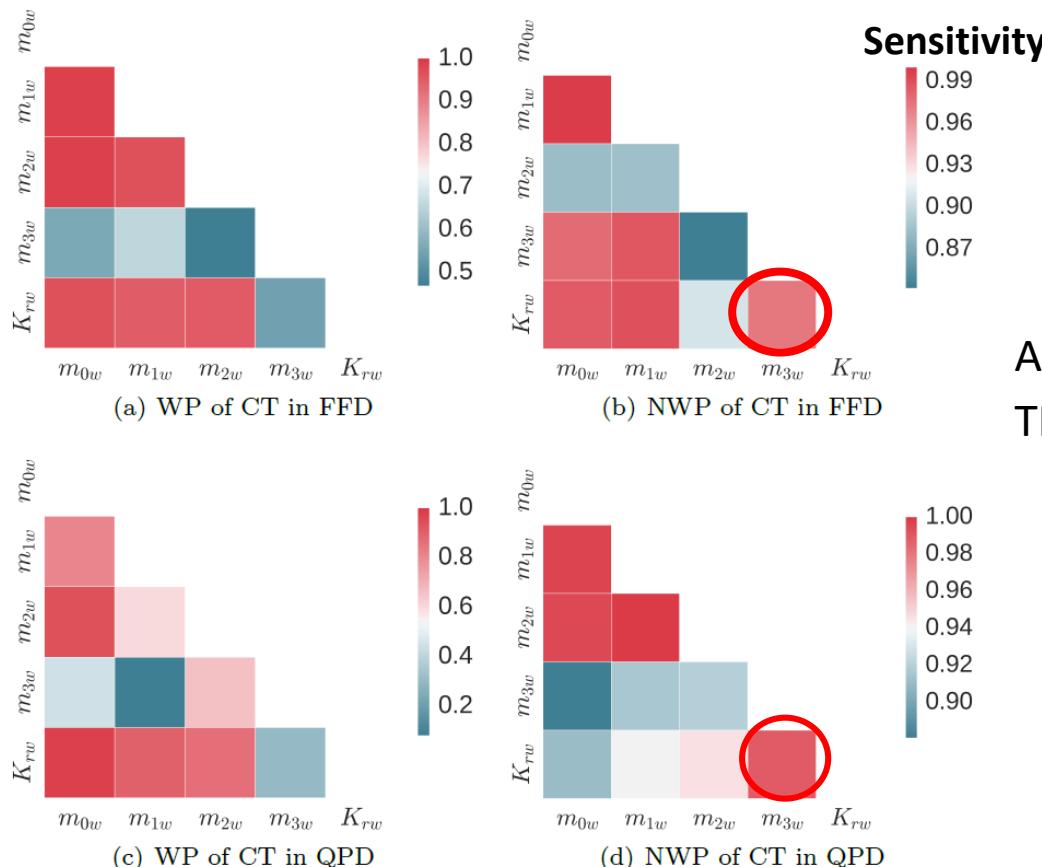
Rock: Gildehauser sandstone (similar to Bentheimer)

Imbibition



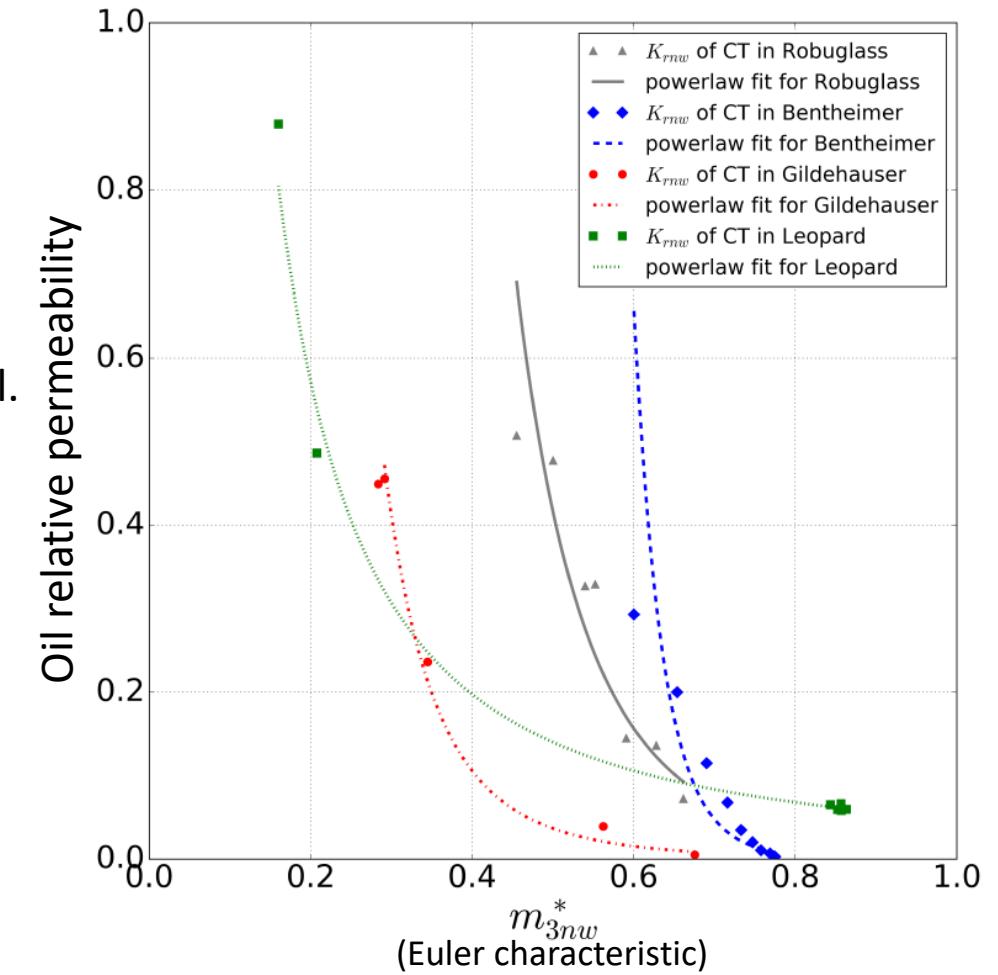
# Application: new route to relative permeability

Non-wetting phase relative permeability has high sensitivity with Euler characteristic

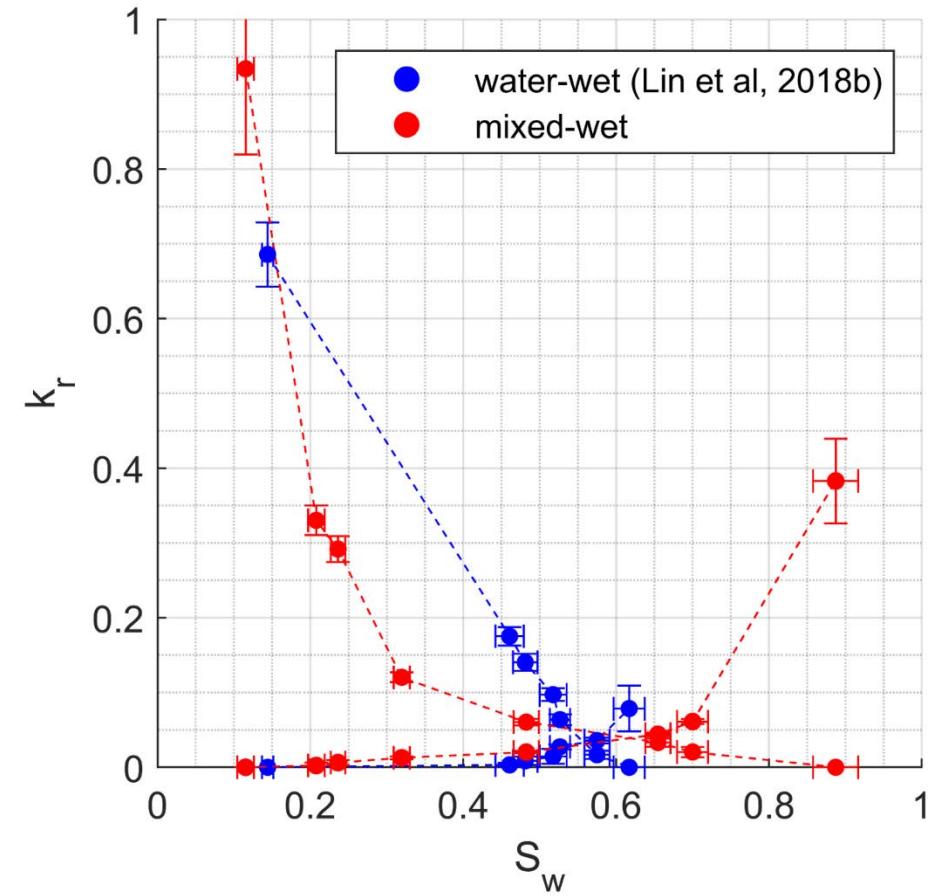


Armstrong et al.  
TIPM 2017

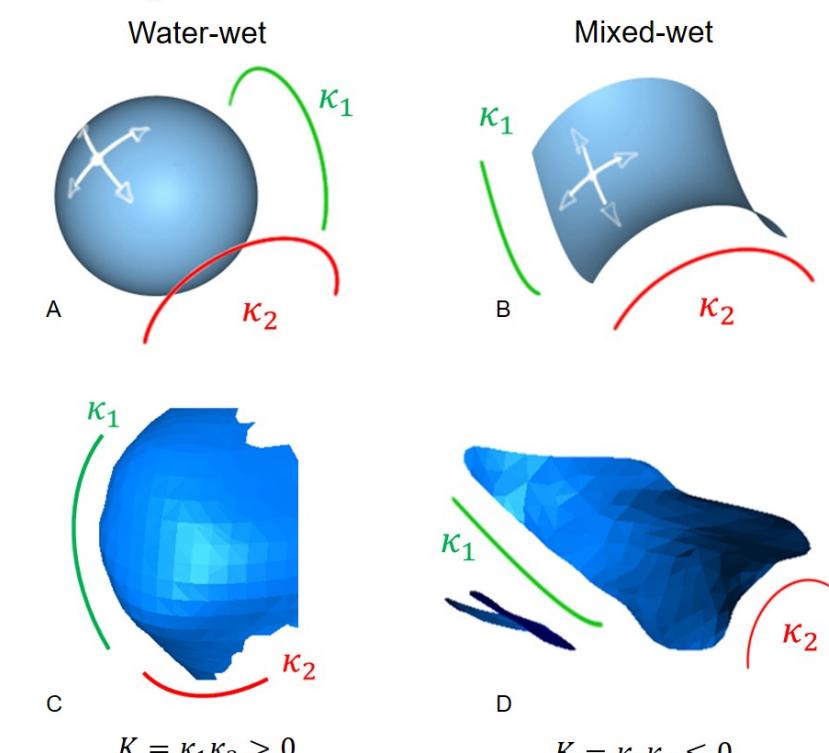
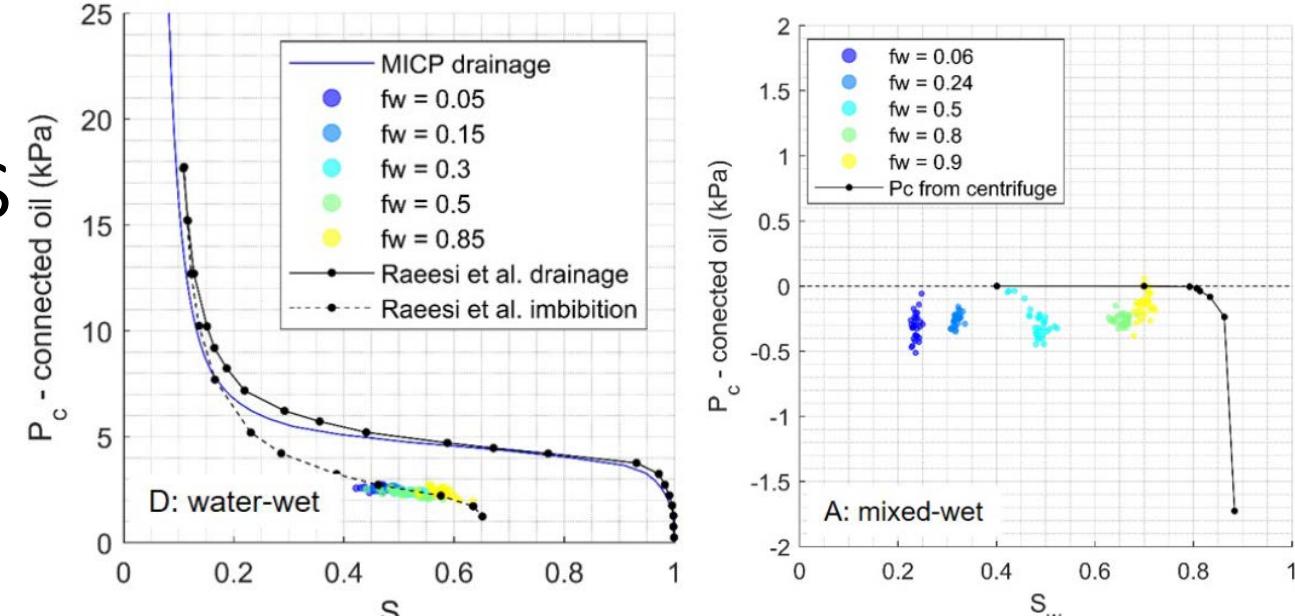
Non-wetting phase relative permeability is simple  
Power law function of Euler characteristic



# Bi-Continuous Interfaces



→ bi-continuous interfaces with high connectivity

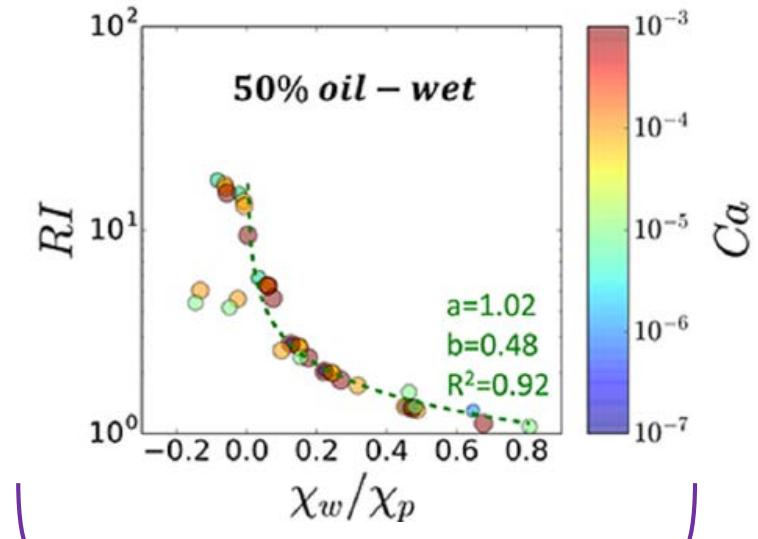
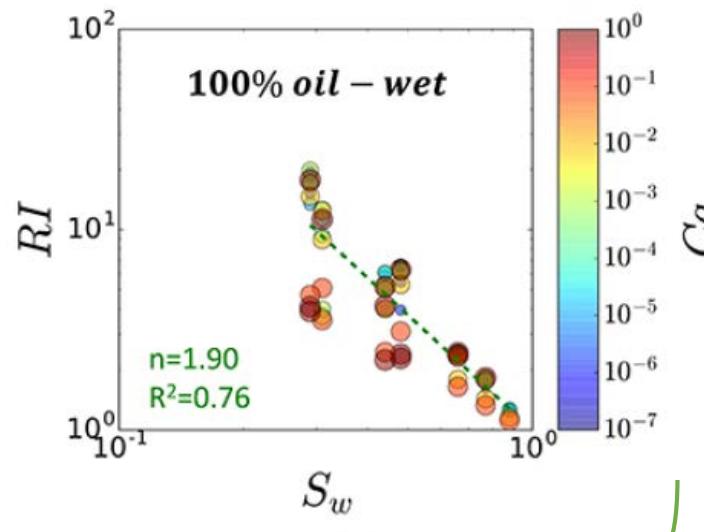
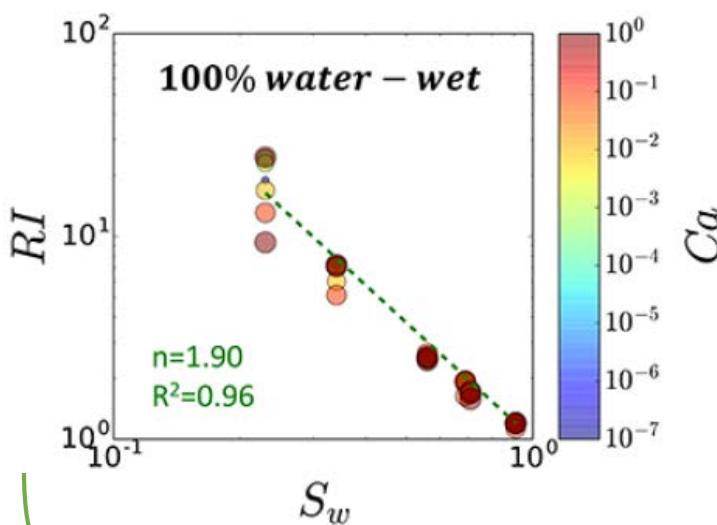


# Resistivity Index and Topology

[Liu et al. 2018]

Percolation Theory:

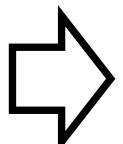
$$\xi \sim (p - p_c)^\beta$$



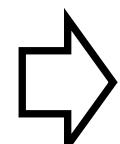
Percolation parameter = Saturation

Percolation parameter = Topology

Models comparing percolation parameters for various wetting conditions.

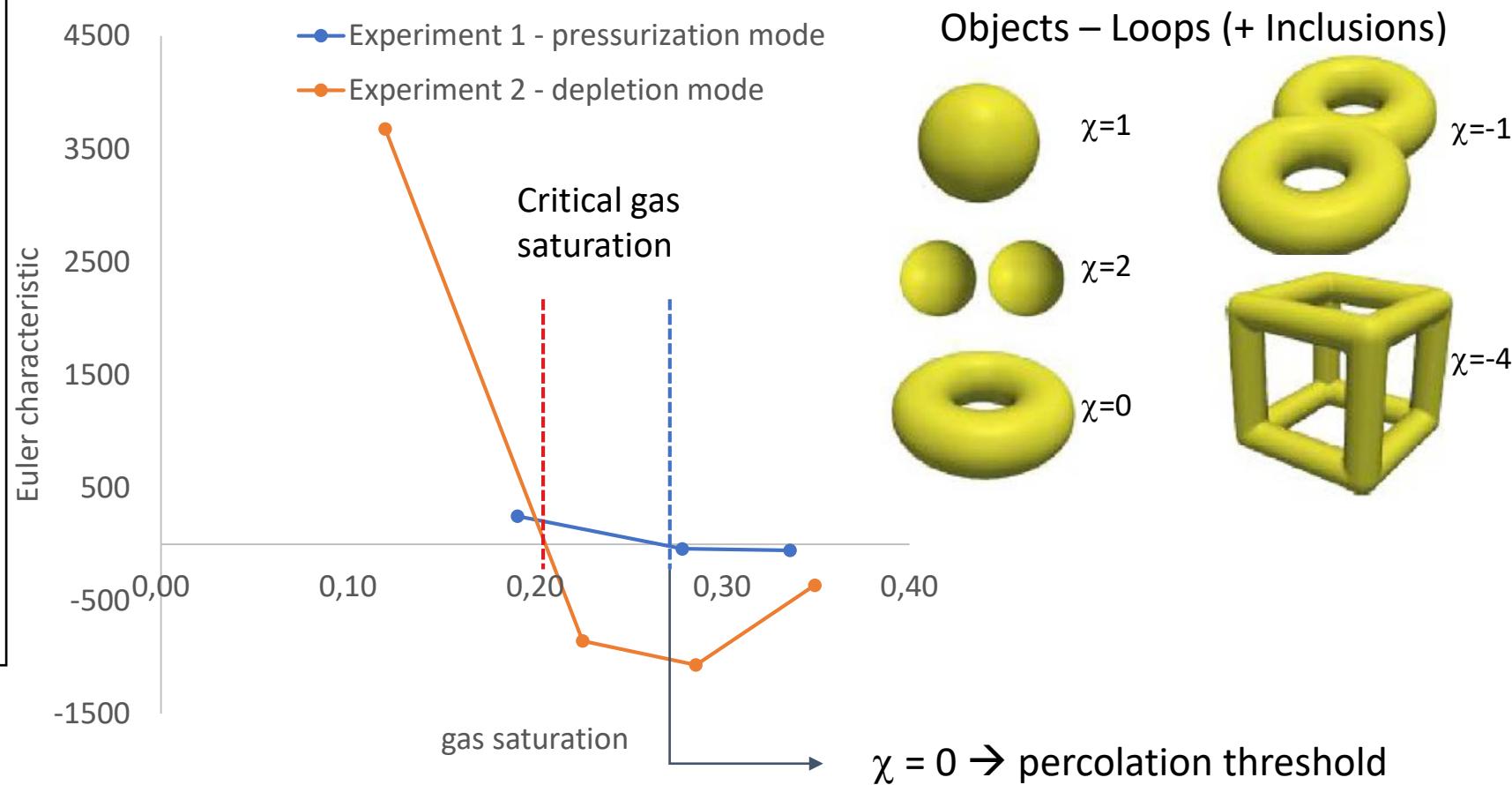


Wettability	$lg(RI) \sim lg(S_w)$	$lg(RI) \sim lg(\chi_w/\chi_p)$
100% water-wet	0.98	0.22
50% water-wet	0.97	0.97
50% oil-wet	0.88	0.96
100% oil-wet	0.82	0.90



Euler characteristic inferred from RI

# Phase Connectivity in Critical Gas Saturation



→ See SCA021 on Wednesday, 11:30

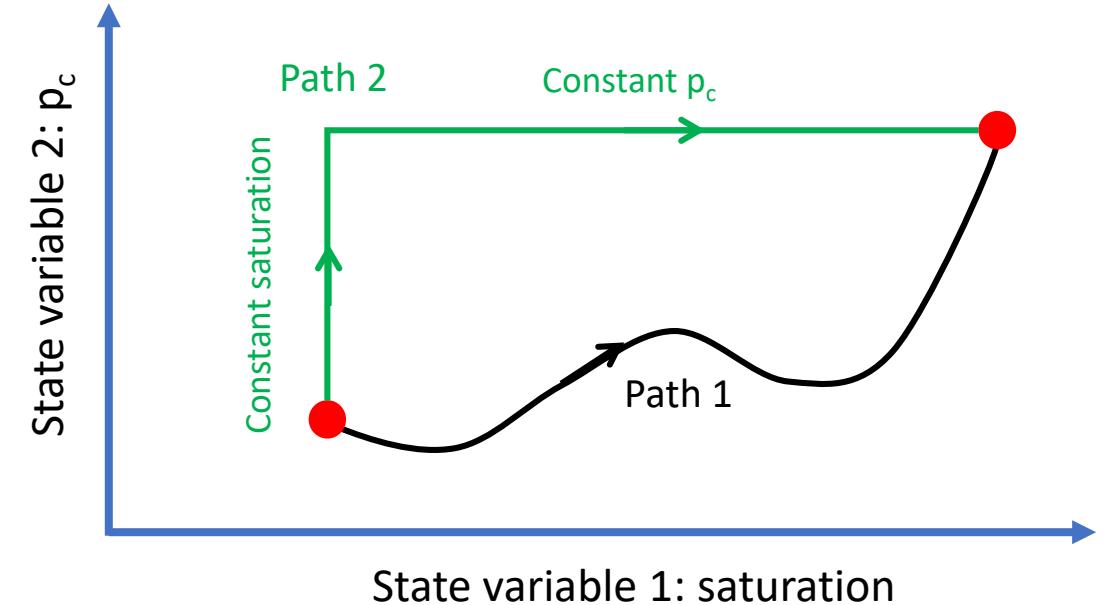
# Application: Relative Permeability Hysteresis Modelling

Advantage of state-variable description: path-independence

1. Can express any parameter, e.g.  $k_r$  as total differential of state variables

$$dk_r = \underbrace{\frac{\partial k_r}{\partial S} dS + \frac{\partial k_r}{\partial \hat{\chi}} d\hat{\chi}}_{\text{Phase distribution}} + \underbrace{\frac{\partial k_r}{\partial I} dI}_{\text{Wettability}} + \underbrace{\frac{\partial k_r}{\partial N_{CA}} dN_{CA}}_{\text{Capillary number}} + \underbrace{\frac{\partial k_r}{\partial \lambda} d\lambda}_{\text{Rock structure}}$$

SPE-182655 (Penn State group)



2. Can choose different paths to measure  $k_r$

- e.g. branch 1: constant saturation (steady-state, constant fractional flow  $F_w$ )  
branch 2: constant  $p_c$  (porous plate)

# Application: Relative Permeability Hysteresis Modelling

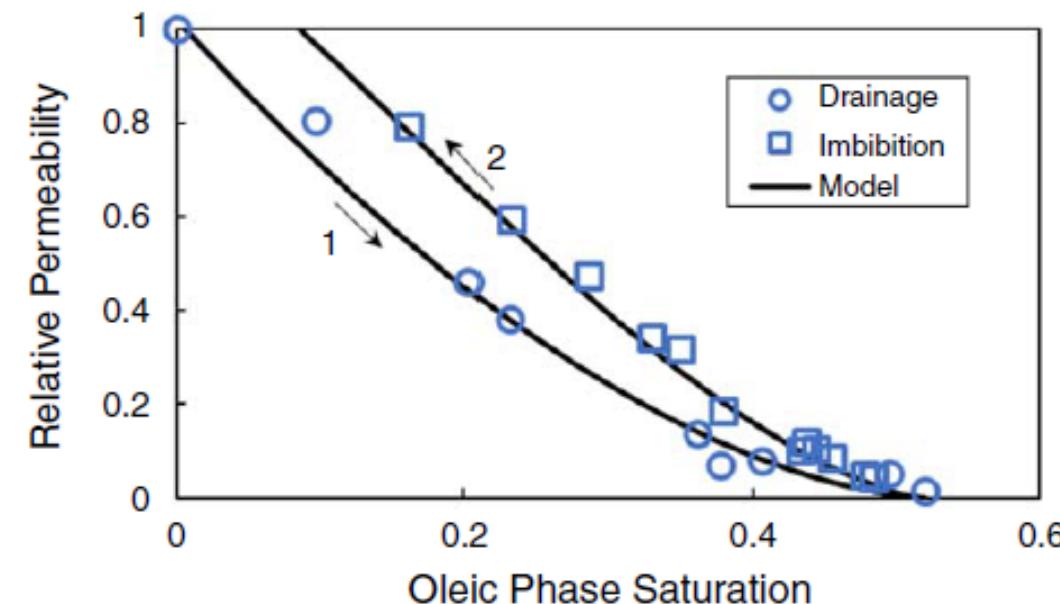
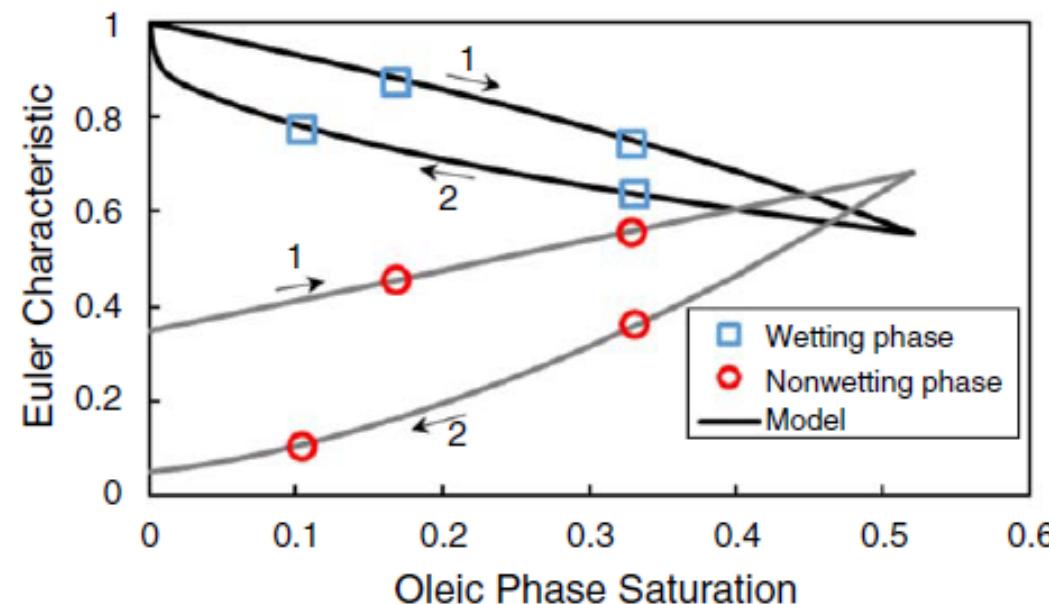
Relative permeability as an

Equation of State (EOS)

SPE-182655 (R. T. Johns)

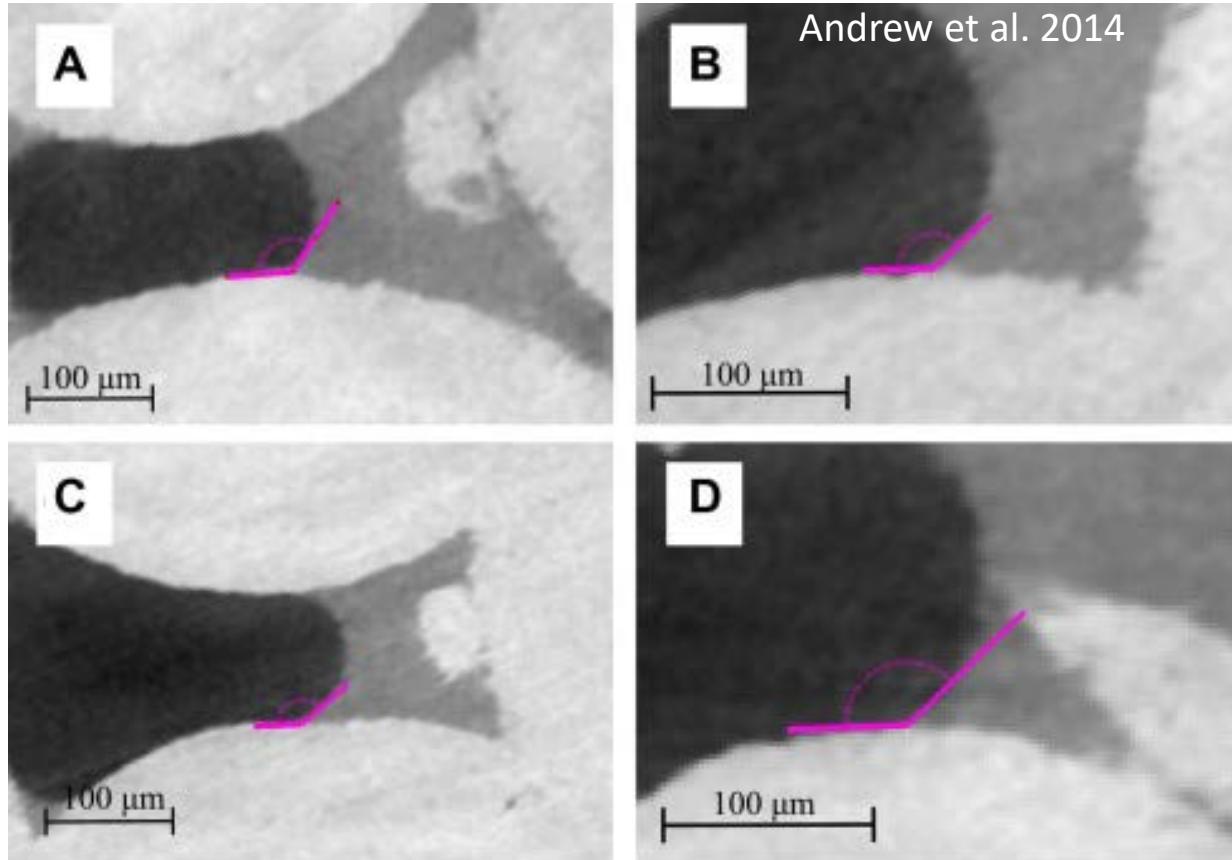
$$dk_r = \underbrace{\frac{\partial k_r}{\partial S} dS + \frac{\partial k_r}{\partial \hat{\chi}} d\hat{\chi}}_{\text{Phase distribution}} + \underbrace{\frac{\partial k_r}{\partial I} dI}_{\text{Wettability}} + \underbrace{\frac{\partial k_r}{\partial N_{CA}} dN_{CA}}_{\text{Capillary number}} + \underbrace{\frac{\partial k_r}{\partial \lambda} d\lambda}_{\text{Rock structure}}$$

$$\begin{array}{c|cc} & \frac{\partial S}{\partial t} > 0 & \frac{\partial S}{\partial t} < 0 \\ \hline \frac{\partial \hat{\chi}}{\partial S} & \alpha_x \left( \frac{\hat{\chi} - 1}{S - 1} \right) & \frac{1}{C_x (\hat{\chi} S)^{n_x}} \end{array}$$



# Application: Wettability

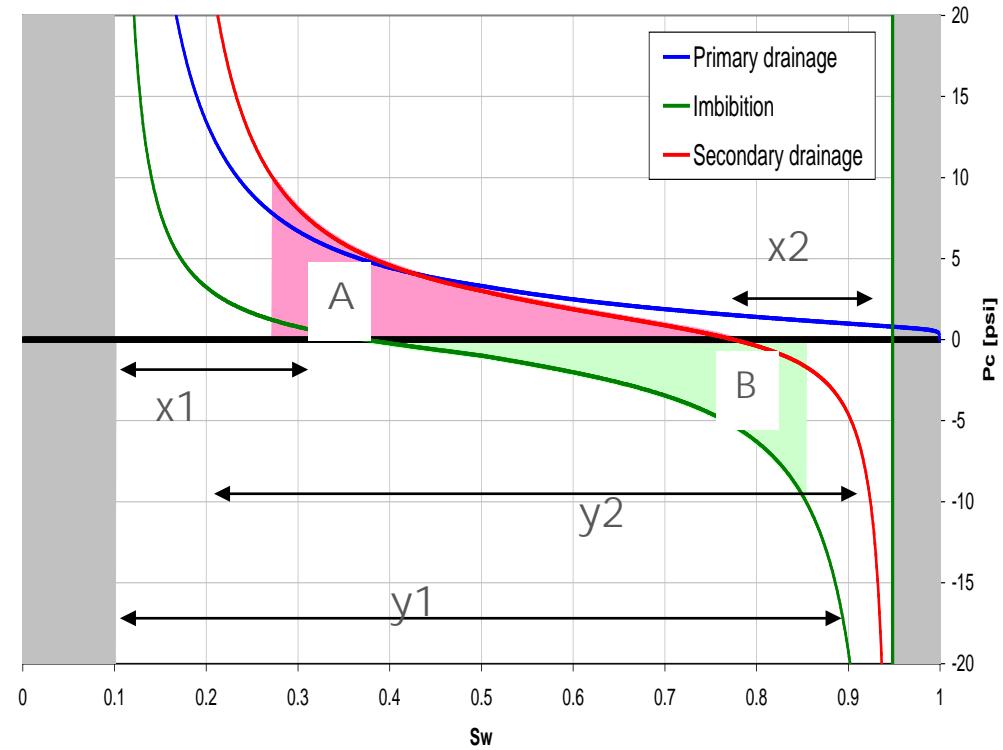
Pore scale



Effective contact angle  $\neq$  intrinsic contact angle  
(Intrinsic contact angle  $\sim$  surface energy)

Darcy-SCAL scale

via capillary pressure curve



# Contact Angle from Deficit Curvature

Gauss-Bonnet Theorem

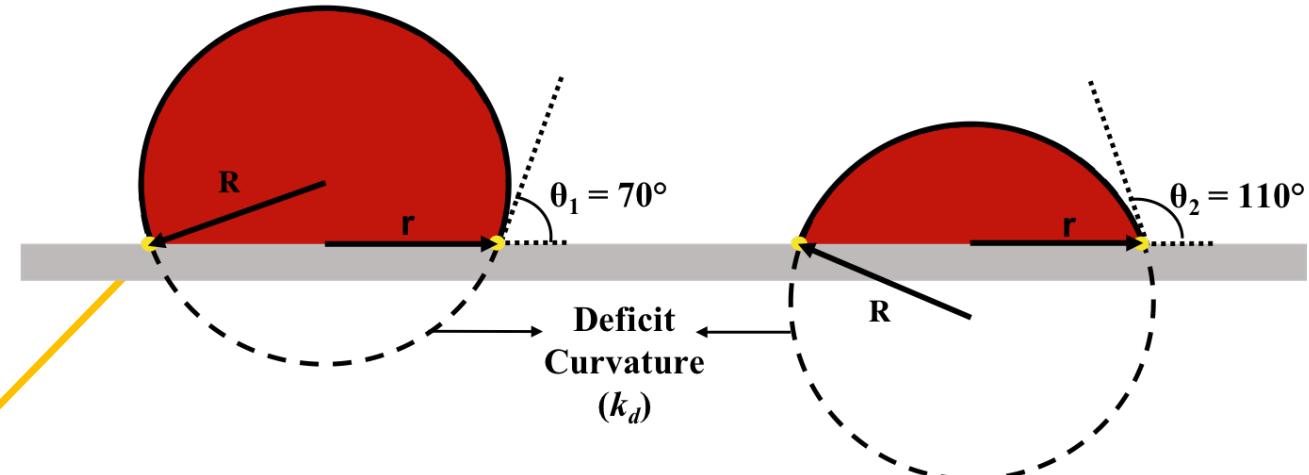
$$\int_S [1/(r_1 r_2)] ds + \boxed{\int_B k_g dl} = 2\pi\chi(S)$$

Total geodesic curvature  
along the contact line

Deficit Curvature,  $k_d$

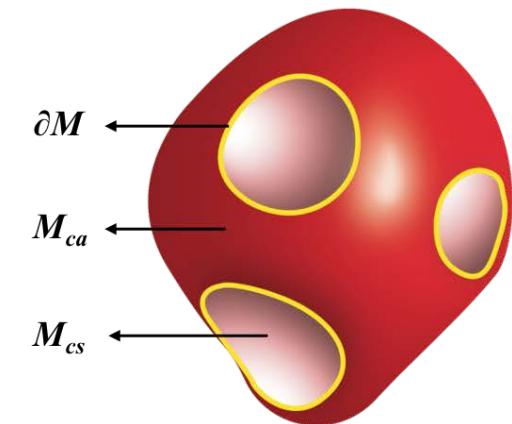
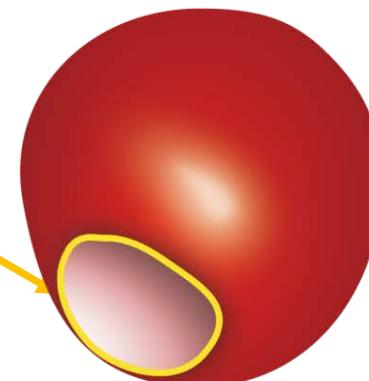
Macroscopic  
contact angle

$$\theta_{macro} = \frac{k_d}{4N_c}$$



Sessile Drop

3D Oil Cluster



# Contact Angle = Deficit Curvature

Gauss Bonnet Theorem

$$\text{Gaussian curvature} \quad \text{Geodesic curvature}$$

$$2\pi\chi(M) = \int_M \kappa_T dS + \int_{\partial M} \kappa_g dC$$



For each region:

Fluid-Fluid

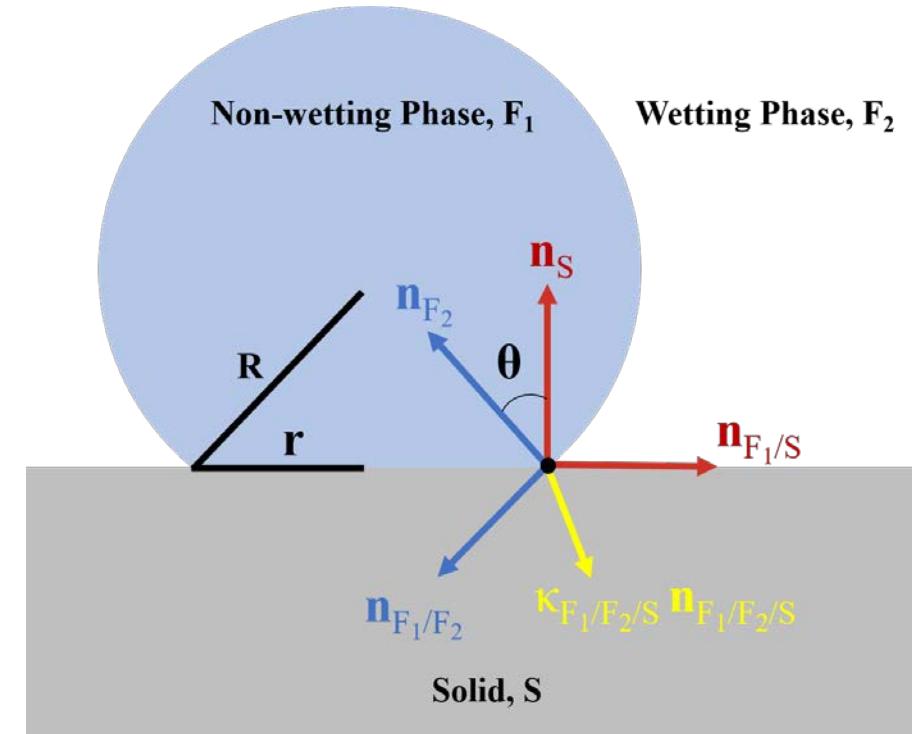
$$2\pi\chi(M_{ab}) = \int_{M_{ab}} \kappa_T dS + \int_{\partial M_{ab}} \kappa_g dC,$$

Fluid-solid

$$2\pi\chi(M_{bs}) = \int_{M_{bs}} \kappa_T dS + \int_{\partial M_{bs}} \kappa_g dC.$$



$$\begin{aligned} 4\pi\chi(C) &= 2\pi\chi(M_{ab}) + 2\pi\chi(M_{bs}) \\ &= \int_{M_{ab}} \kappa_T dS + \int_{M_{bs}} \kappa_T dS + \int_{\partial M} (\kappa_{g_{ab}} + \kappa_{g_{bs}}) dC \end{aligned}$$



# Contact Angle = Deficit Curvature

$$k_d = 4\pi\chi(C) - \kappa_{ab}A_{ab} - \kappa_{bs}A_{bs}$$



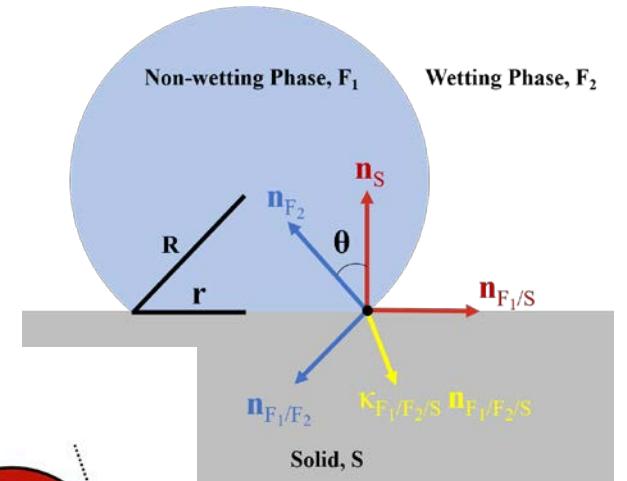
$$4\pi\chi(C) = 2\pi\chi(M_{ab}) + 2\pi\chi(M_{bs})$$

with

$$\kappa_{ab} = \frac{1}{A_{ab}} \int_{M_{ab}} \kappa_T dS,$$

$$\kappa_{bs} = \frac{1}{A_{bs}} \int_{M_{bs}} \kappa_T dS,$$

$$k_d = \int_{\partial M} (\kappa_{g_{ab}} + \kappa_{g_{bs}}) dC = \int_{\partial M} \kappa_{abs} \mathbf{n}_{abs} \cdot [\mathbf{n}_S \sin \theta + \mathbf{n}_{bs} (1 - \cos \theta)] dC$$



For  $N_c$  contacts with solid

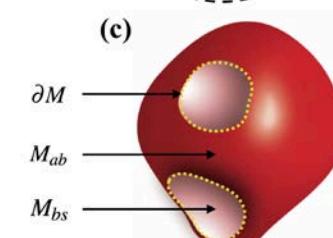
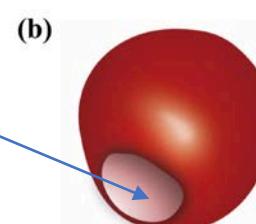
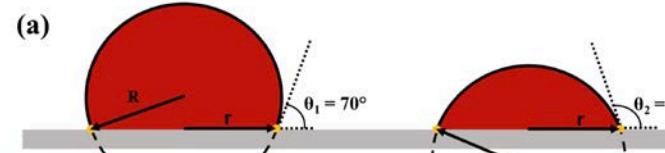
$$k_d = \sum_{j=1}^{N_c} \int_{\partial M_j} \kappa_{abs} \mathbf{n}_{abs} \cdot [\mathbf{n}_S \sin \theta + \mathbf{n}_{bs} (1 - \cos \theta)] dC$$



Macroscopic contact angle

$$\theta^{macro} = \frac{k_d}{4N_c}$$

Number  
of contacts  
 $N_c$



# Contact Angle = Deficit Curvature

Macroscopic contact angle  $\theta^{macro} = \frac{k_d}{4N_c}$

For 1 sessile droplet:

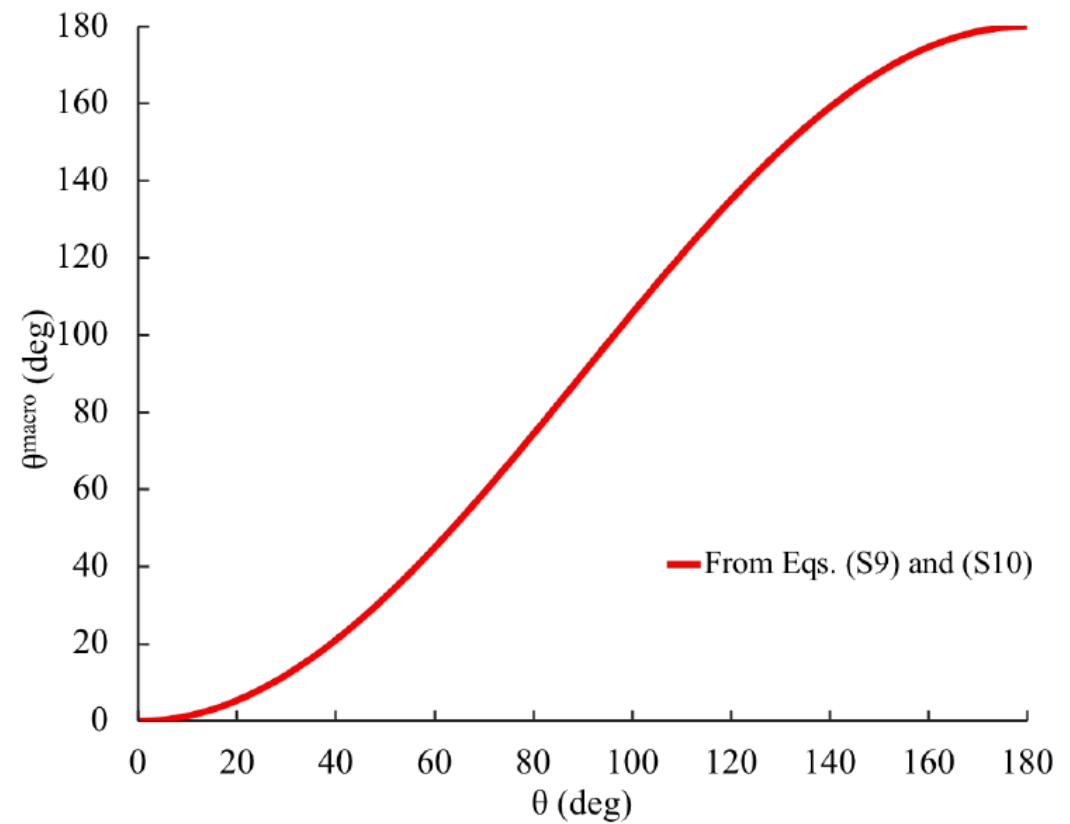
$$\kappa_{gab} = -\sqrt{1 - (r/R)^2}/r \quad \kappa_{gbs} = 1/r. \quad r = R \sin \theta.$$



$$\begin{aligned} k_d &= \int_{\partial M} \frac{1}{r} \left( 1 - \sqrt{1 - (r/R)^2} \right) dC \\ &= 2\pi \left( 1 - \sqrt{1 - (R \sin \theta / R)^2} \right) \\ &= 2\pi(1 - \cos \theta). \end{aligned}$$



$$\theta^{macro} = \frac{\pi(1 - \cos \theta)}{2}$$



# Contact Angle = Deficit Curvature

Armstrong et al. *under review*

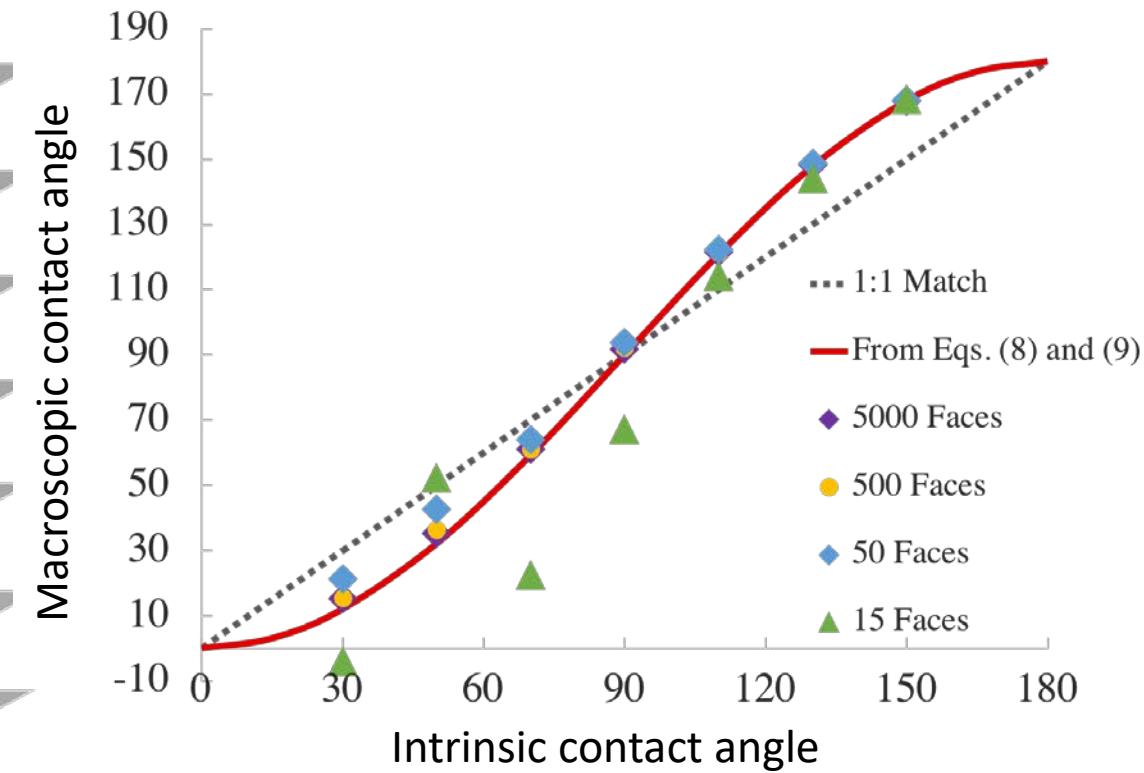
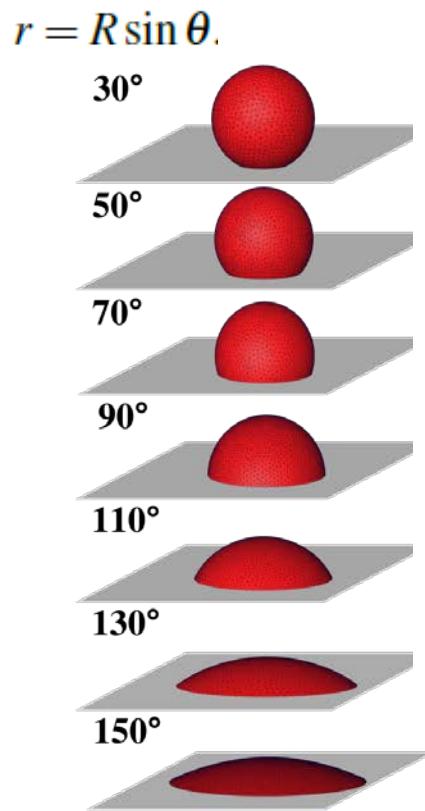
$$\text{Macroscopic contact angle} \quad \theta^{macro} = \frac{k_d}{4N_c}$$

For 1 sessile droplet:

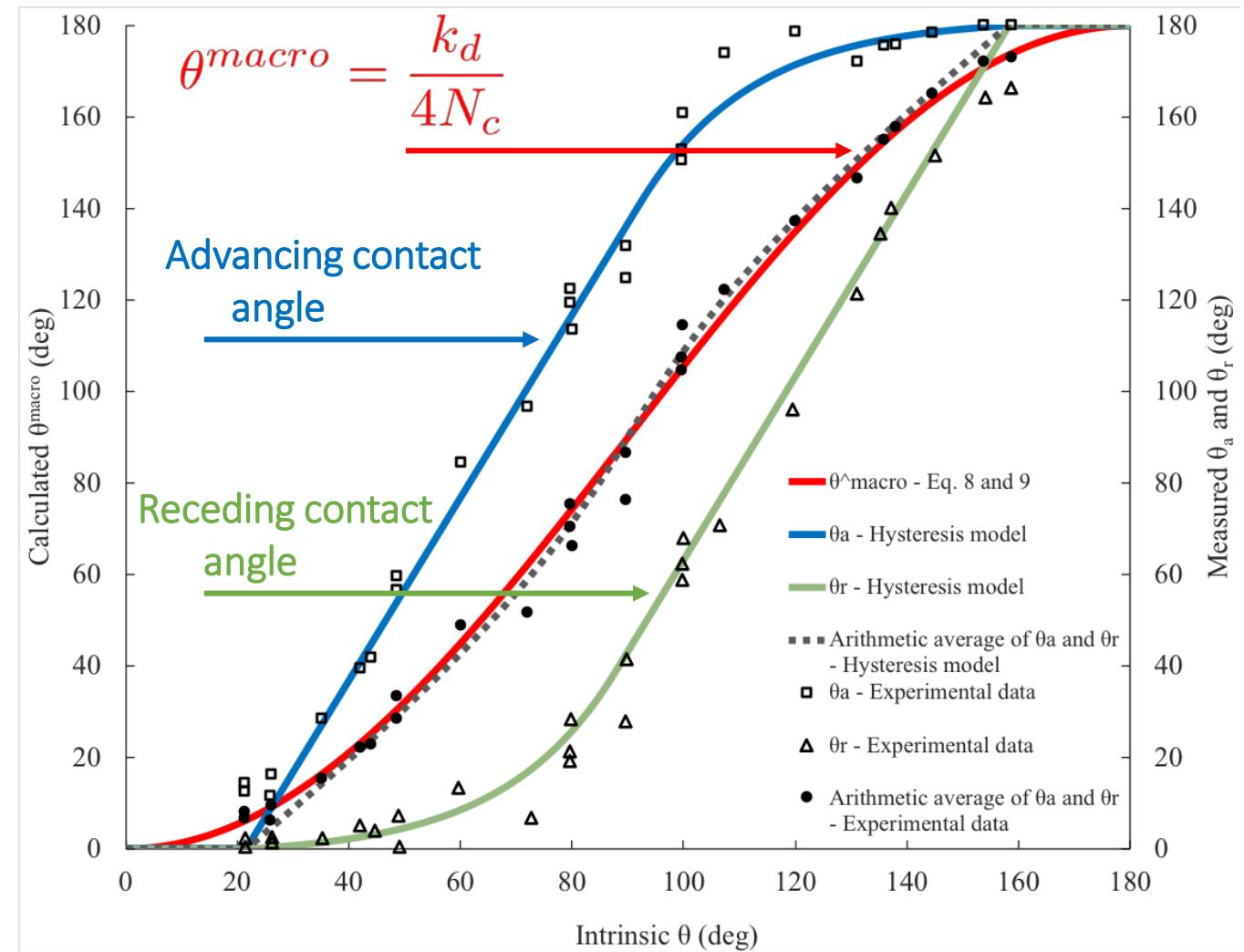
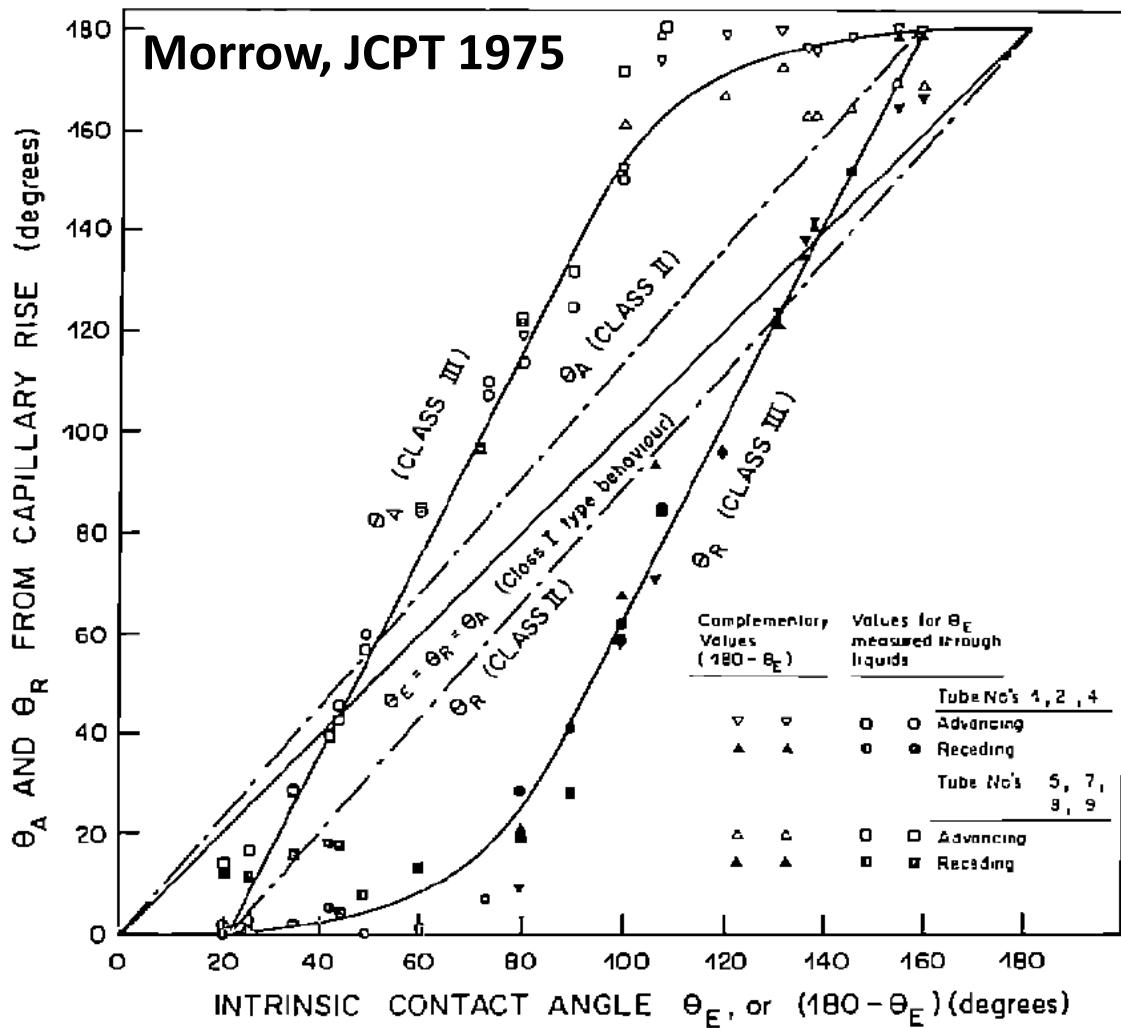
$$\kappa_{gab} = -\sqrt{1-(r/R)^2}/r \quad \kappa_{gbs} = 1/r.$$

$$\begin{aligned} k_d &= \int_{\partial M} \frac{1}{r} \left( 1 - \sqrt{1 - (r/R)^2} \right) dC \\ &= 2\pi \left( 1 - \sqrt{1 - (R \sin \theta / R)^2} \right) \\ &= 2\pi(1 - \cos \theta). \end{aligned}$$

$$\theta^{macro} = \frac{\pi(1 - \cos\theta)}{2}$$

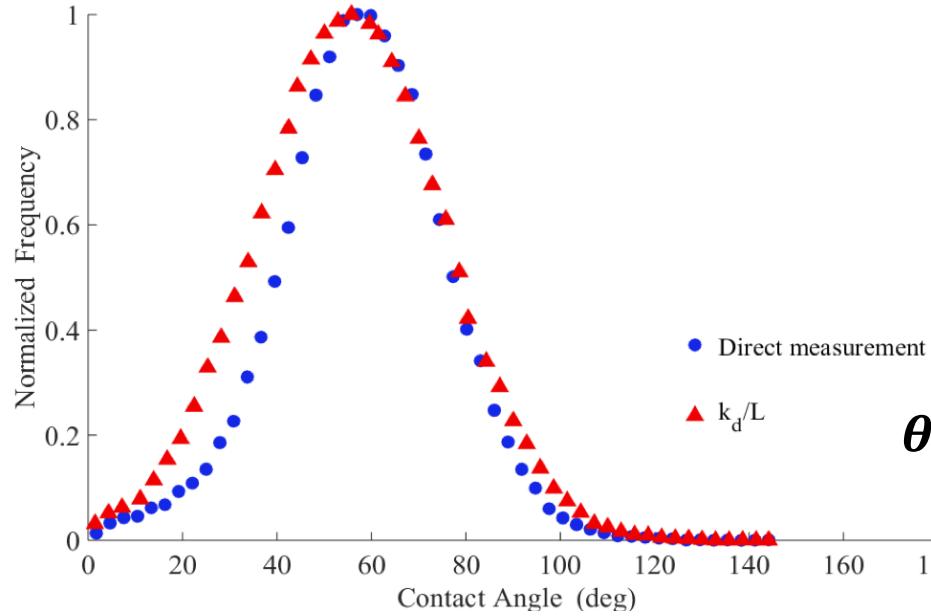


# Validation: Advancing and Receding Contact angle

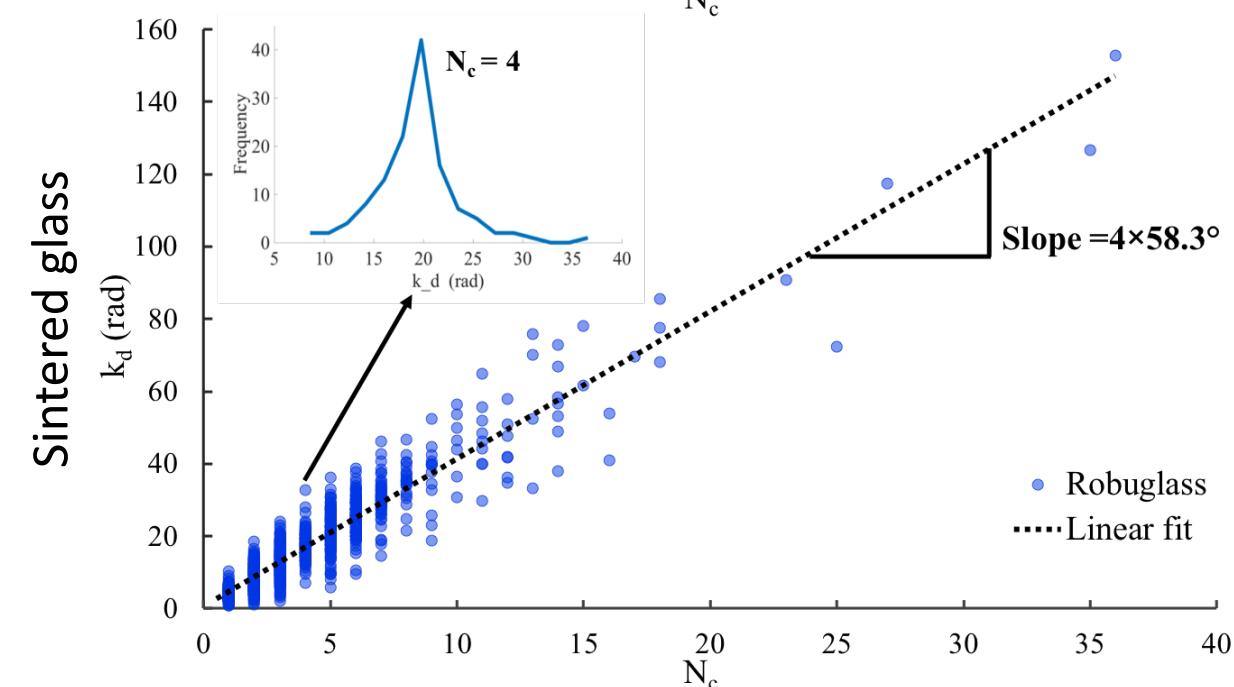
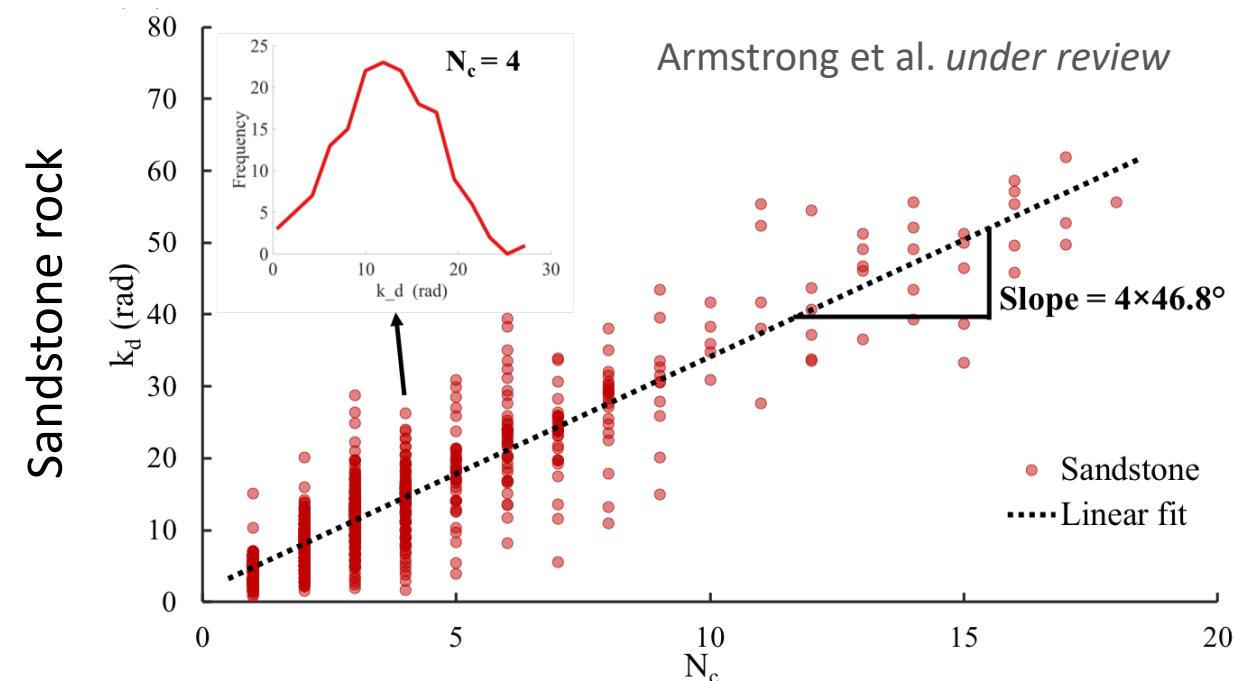
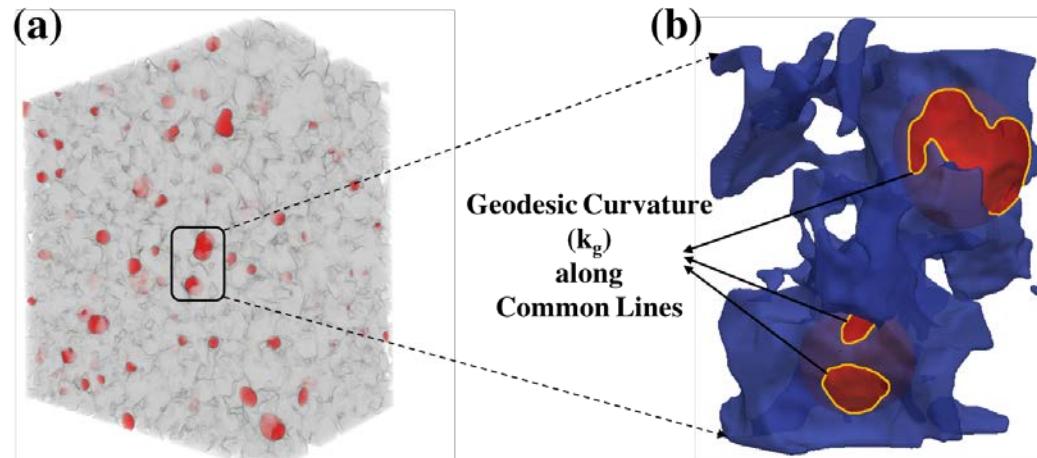


# Predicting Contact Angle from oil cluster distribution & $p_c(S_w)$

→ SCA009 (Tuesday, 9:30)



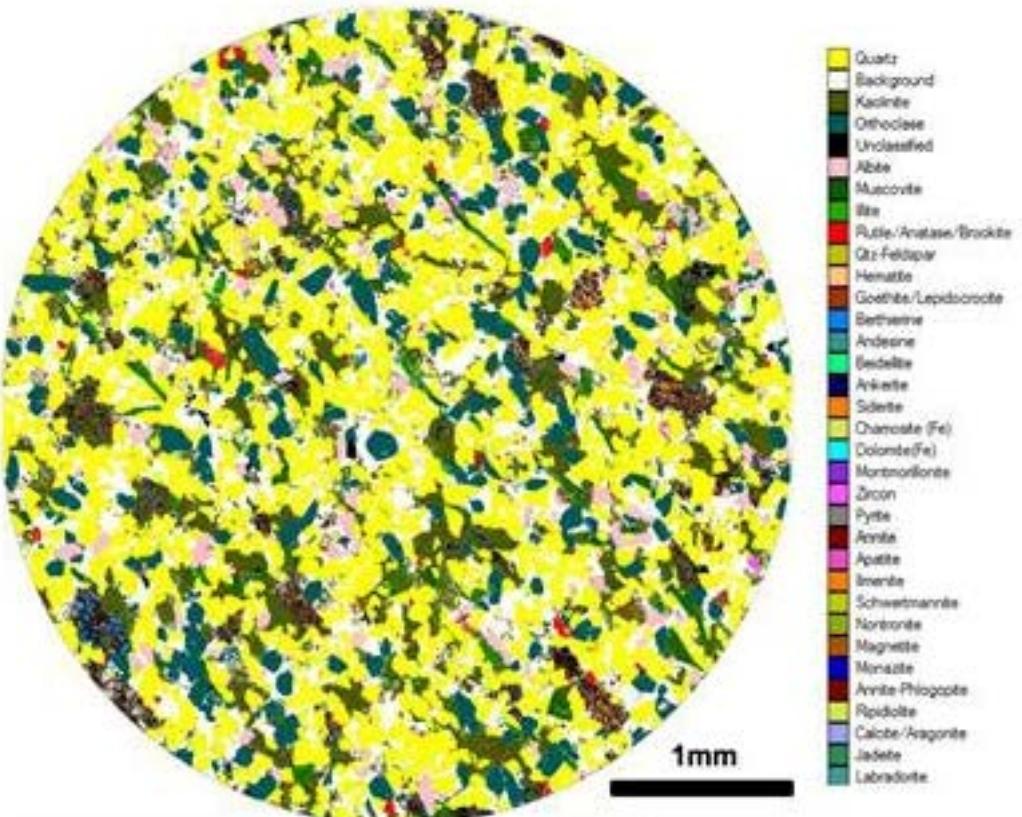
$$\theta_{macro} = \frac{\kappa_d}{4N_c}$$



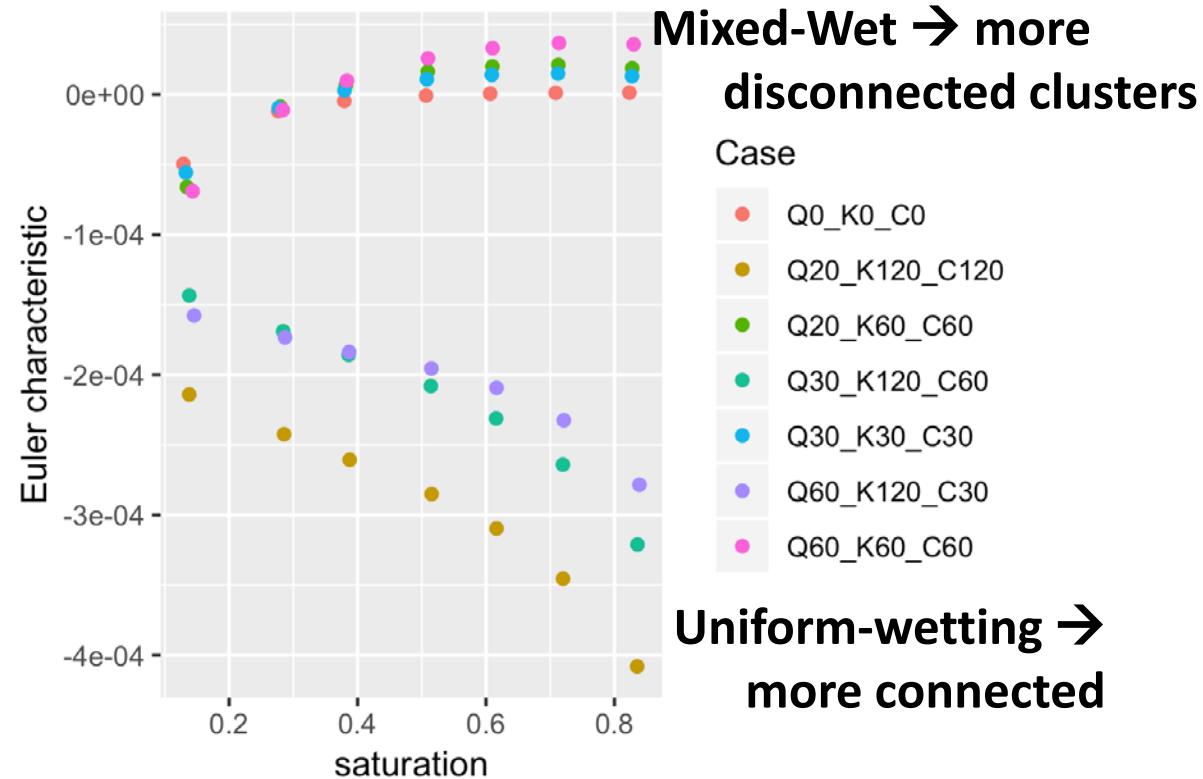
Armstrong et al. under review

# Minerology, Wettability and Topology

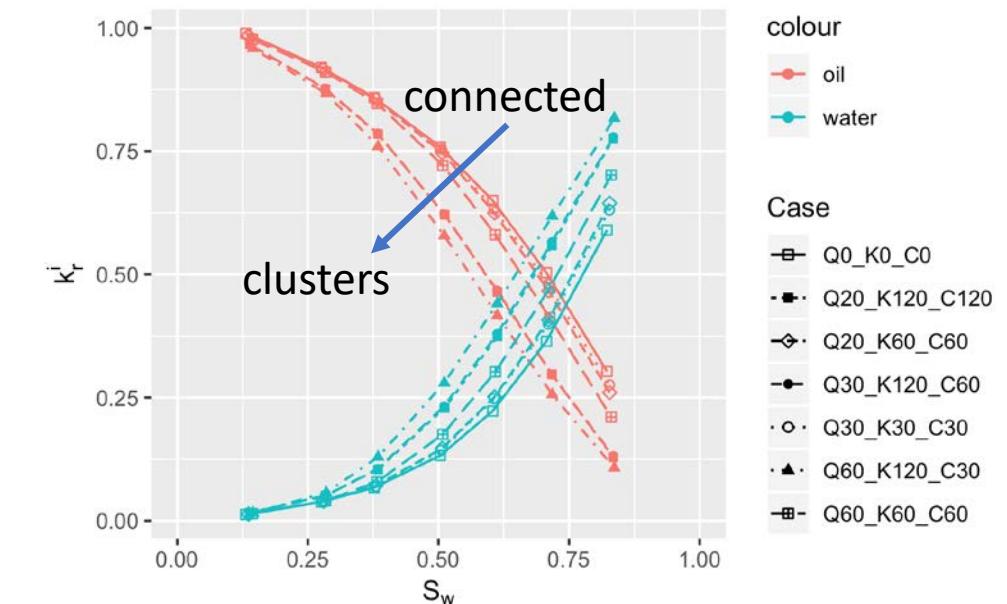
QEMSCAN DATA, Pore-Scale Minerology



Euler characteristic quantifies impact of wettability on the connectivity of the oil phase



Uniform-wetting → more connected



# Fractured Porous Media

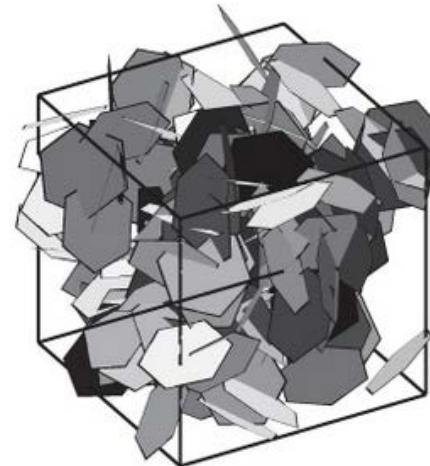
Original works of Adler (Fractured Porous Media)

PRL, Scholz et al. 2012

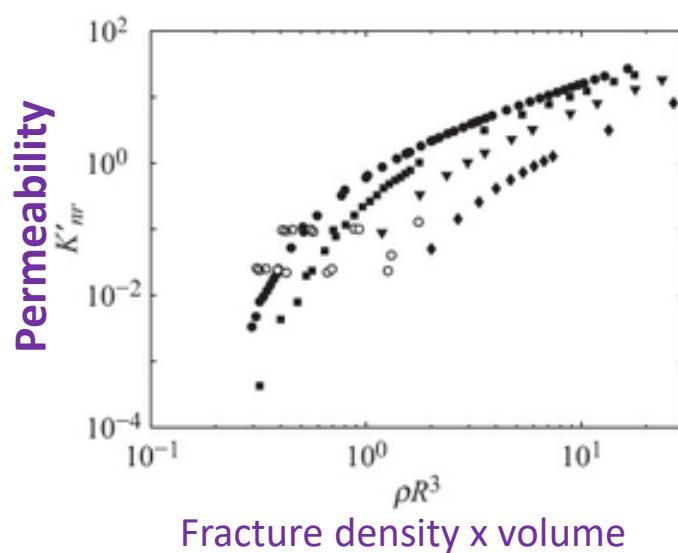
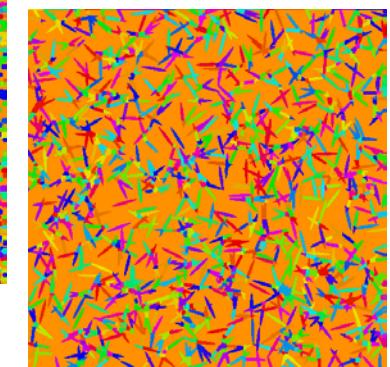
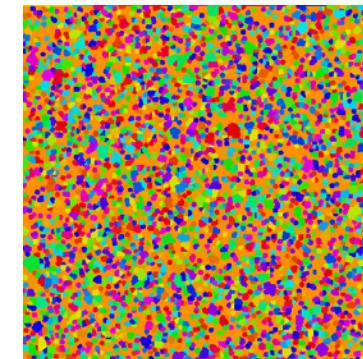
Dimensionless Density,  $\rho'$

A measure of the connectivity of the fracture network

Fracture Network

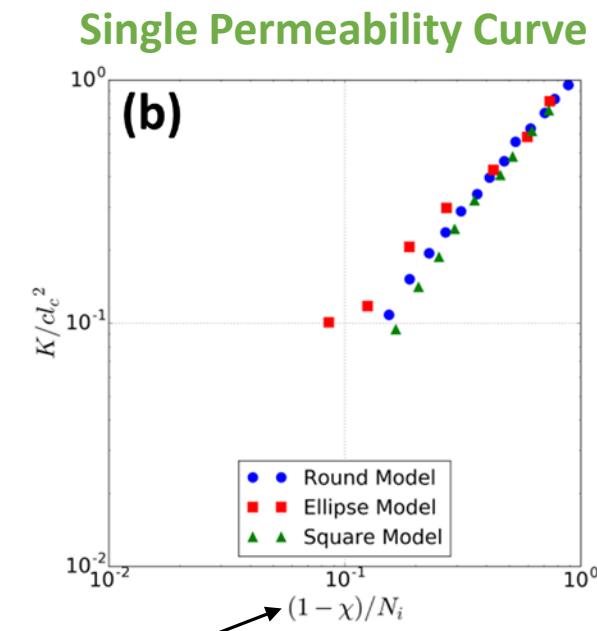
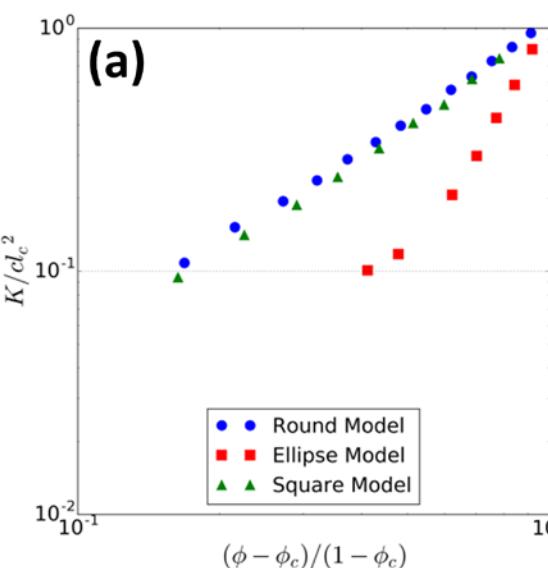
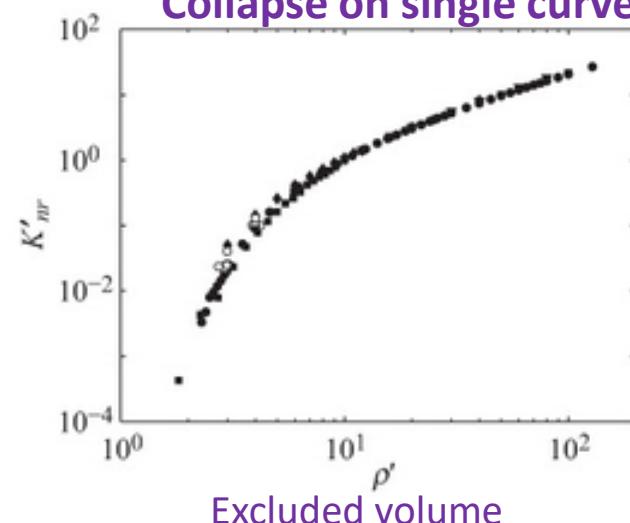


2D Porous Structures



$$\text{Number of Loops, } B_1 = [(\rho'/2) - 1]\rho'$$

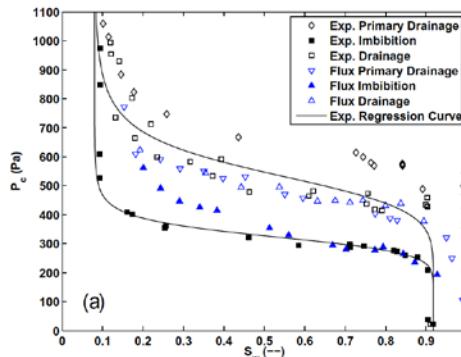
Collapse on single curve



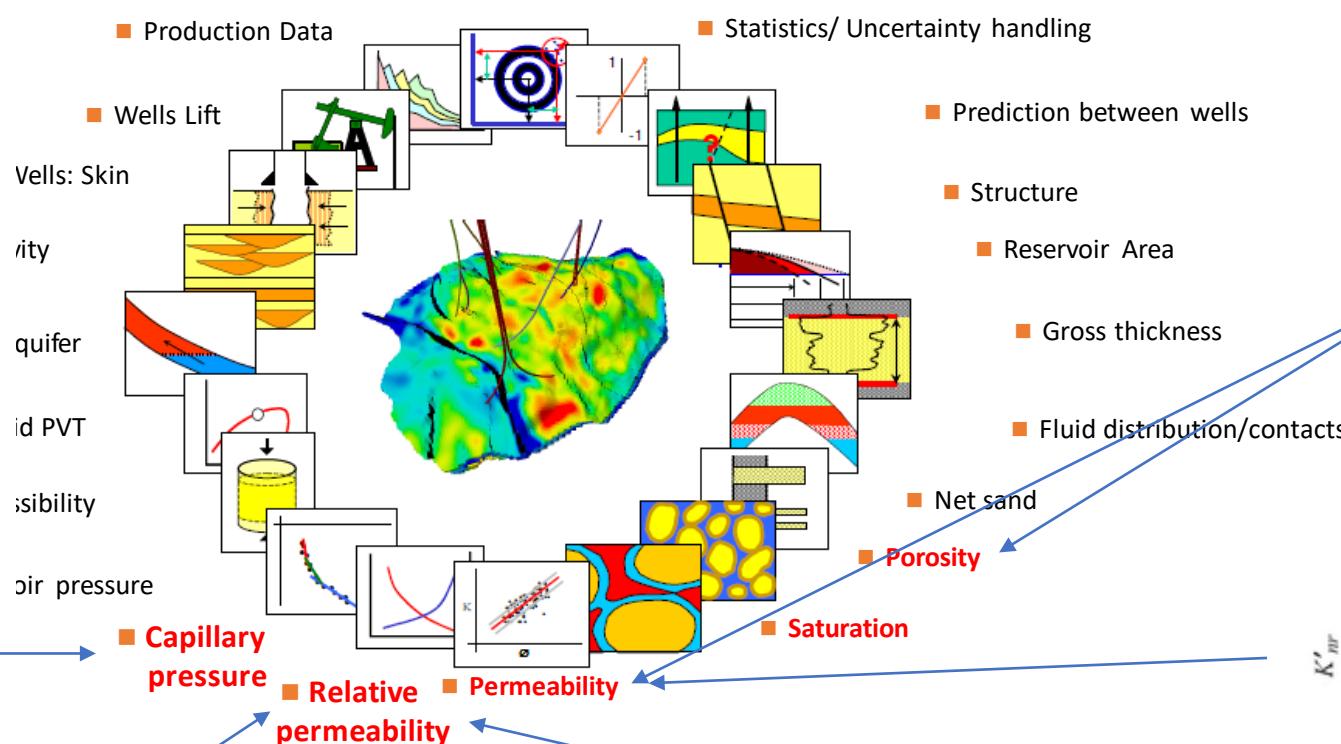
[LOOPS + OBJECTS] / OBJECTS

# Conclusions

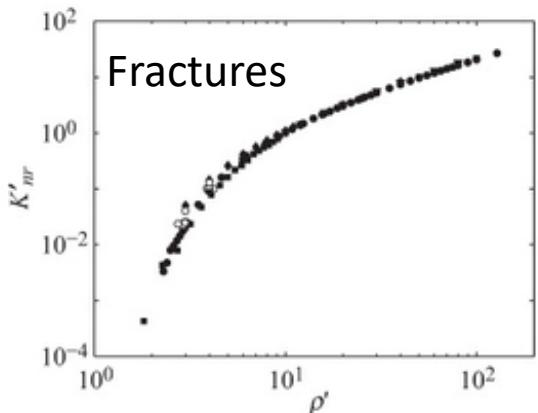
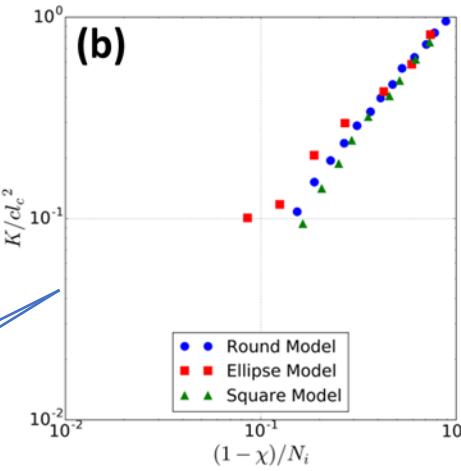
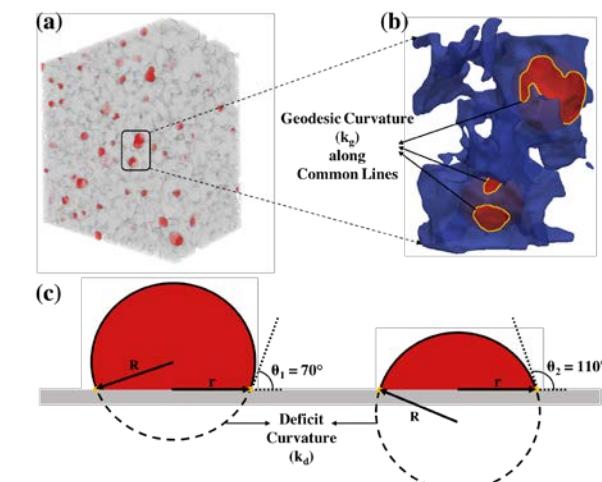
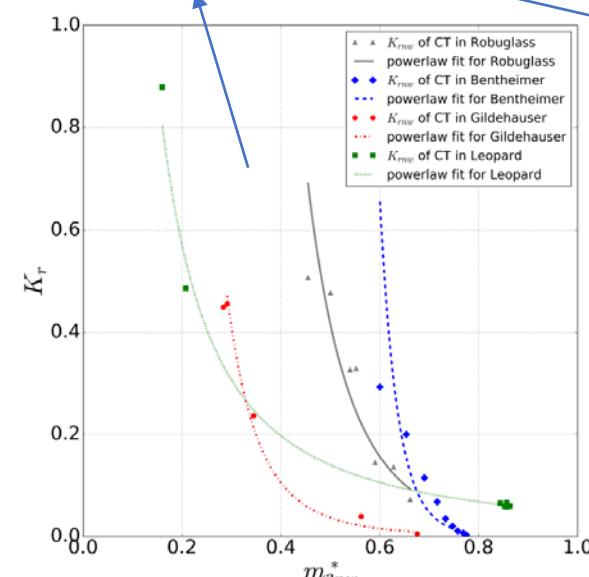
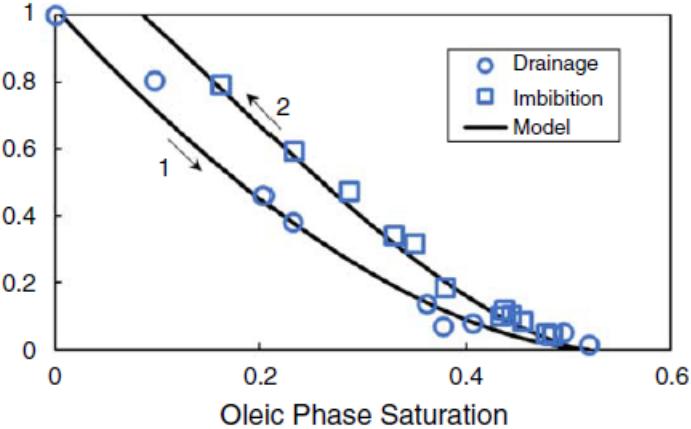
Pc hysteresis



Vels: Skin  
ivity  
quifer  
id PVT  
sibility  
air pressure



Reperm hysteresis



# Key Literature

## Review Papers and Textbooks

- Porous Media Characterization Using Minkowski Functionals: Theories, Applications and Future Directions  
Transport in Porous Media, 2018 *in press* doi:10.1007/s11242-018-1201-4
- J. Ohser, F. Mücklich, Statistical analysis of microstructures in materials science, Wiley, 2000.
- Klaus R. Mecke, Dietrich Stoyan, Statistical Physics and Spatial Statistics. The Art of Analyzing and Modeling Spatial Structures and Pattern Formation, Lecture Notes in Physics, Springer, 2000.

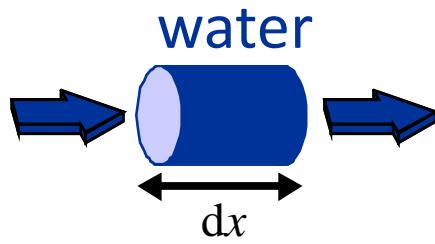
## Research Papers

- C. H. Arns, M. A. Knackstedt, K. Mecke, 3D Structural Analysis: Sensitivity of Minkowski Functionals.  
Journal of Microscopy 240, 181-196, 2010.
- H.J. Vogel, U. Weller, S. Schlüter, Quantification of Soil Structure Based on Minkowski Functions,  
Computers & Geosciences 36, 126-1251, 2010.
- Herring et al. Advances in Water Resources 62, 47-58, 2013.
- S. Schlüter, S. Berg, M. Rücker, R. T. Armstrong, H.-J. Vogel, R. Hilfer and D. Wildenschild,  
Pore scale displacement mechanisms as a source of hysteresis for two-phase flow in porous media  
Water Resources Research 52(3), 2194-2205 2016.
- J. E. McClure, R. T. Armstrong, M. A. Berrill, S. Schlüter, S. Berg, W. G. Gray, C. T. Miller  
A geometric state function for two-fluid flow in porous media, Phys. Rev. Fluids 3(8), 084306, 2018.
- Z. Liu, A. Herring, A. Sheppard, C. Arns, S. Berg, R. T. Armstrong, Morphological characterization of two-phase flow using  
X-ray microcomputed tomography flow-experiments, Transport in Porous Media 118(1), 99-117, 2017.

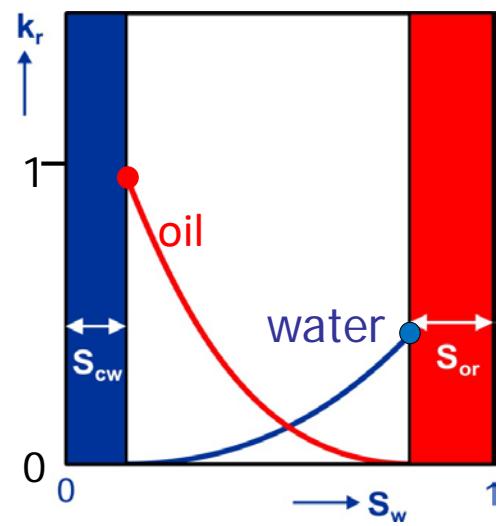
# Backup

# 2-Phase Flow in Porous Media

Single-phase



relative permeability



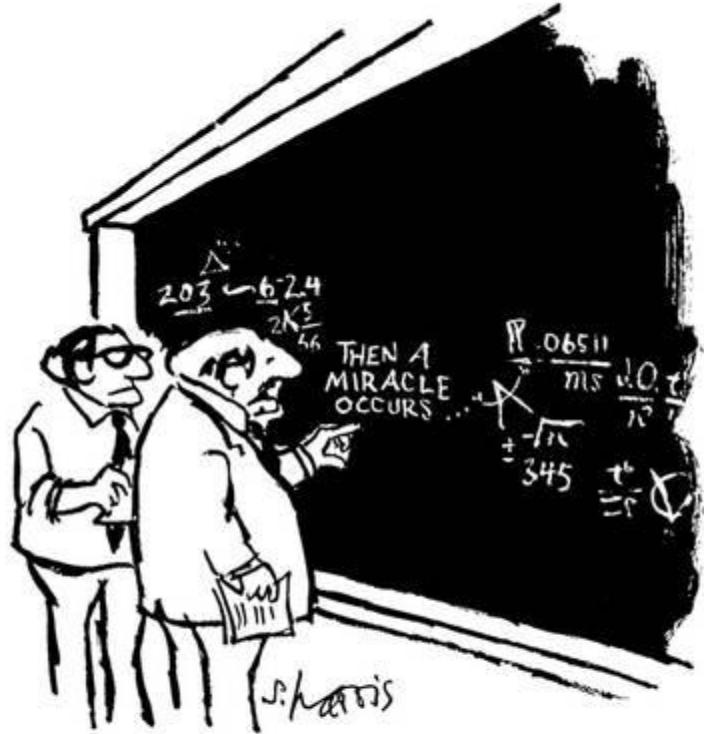
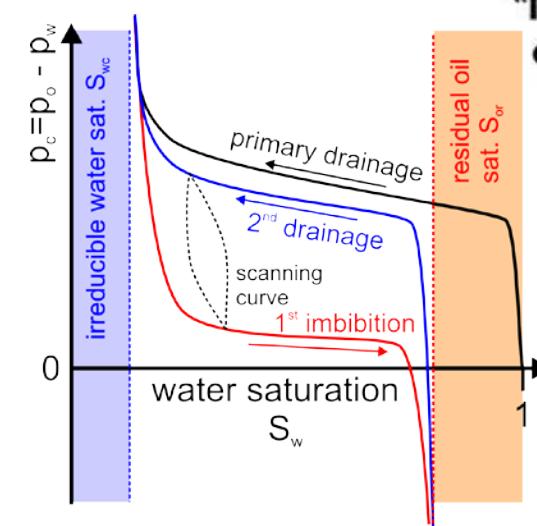
Darcy's law

$$v_{Darcy} = -\frac{K}{\mu} \frac{dp}{dx}$$

Phenomenological  
extension of Darcy's law  
i.e. not a law anymore

$$v_i = -k_{r,i} \frac{K}{\mu_i} \frac{dp_i}{dx} \quad k_{r,i} = k_{r,i}(S_w) \quad p_c = p_o - p_w = p_c(S_w)$$

Capillary pressure



"I think you should be more explicit here in step two."

# So far this has been sufficient, but

When trying to augment SCAL by Digital Rock, it is important to

- correctly classify the problem

And it becomes inevitable to understand what relative permeability actually is, i.e. face the

- Upscaling from Pore to Darcy Scale challenge

More conceptually.

# The State Variables of Capillarity

**Hadwiger's theorem:** unique characterization of 3D objects by 4 Minkowski functionals

$m_0$  = volume (saturation)

$$M_0^n = \lambda(\Omega_n) = \int_{\Omega_n} dr$$

$m_1$  = interfacial area

$$M_1^n = \lambda(\Gamma_n) = \int_{\Gamma_n} dr$$

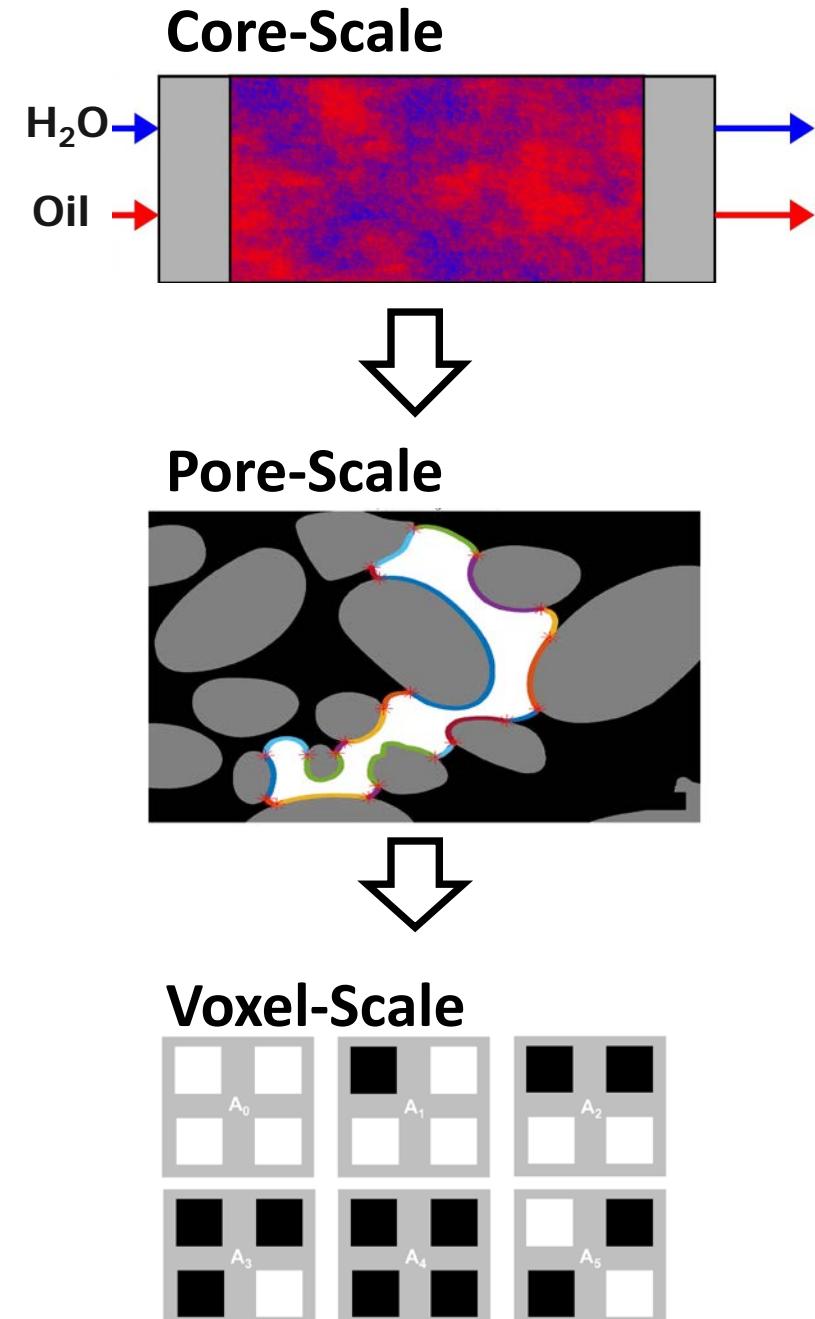
$m_2$  = mean curvature (cap. pressure)

$$M_2^n = \int_{\Gamma_n} \left( \frac{1}{R_1} + \frac{1}{R_2} \right) dr$$

$m_3$  = integral curvature =  $2\pi\chi$

$$M_3^n = \int_{\Gamma_n} \frac{1}{R_1 R_2} dr .$$

McClure et al. Phys. Rev. Fluids, 2018



# Capillary Pressure vs. Saturation

## Capillary Pressure

$$P^c = \gamma J_w^{wn}$$

- $P_c$ - $S_w$  is essentially a geometric definition (or statement)
- Saturation does not uniquely define the geometrical state
- Steiners formula suggests that all four MF are required for a unique definition

## MF in terms of macro-scale parameters

$$M_0^n = \epsilon^n V$$

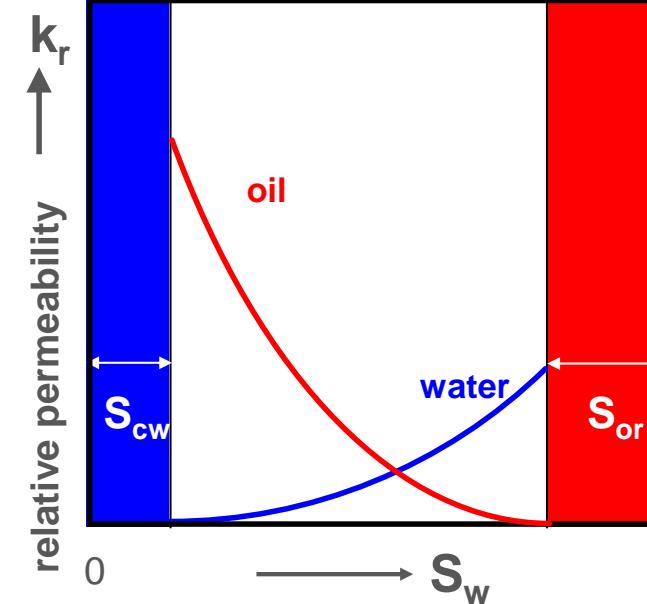
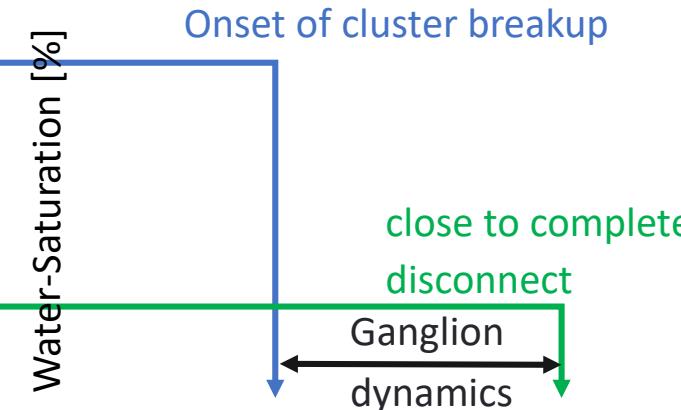
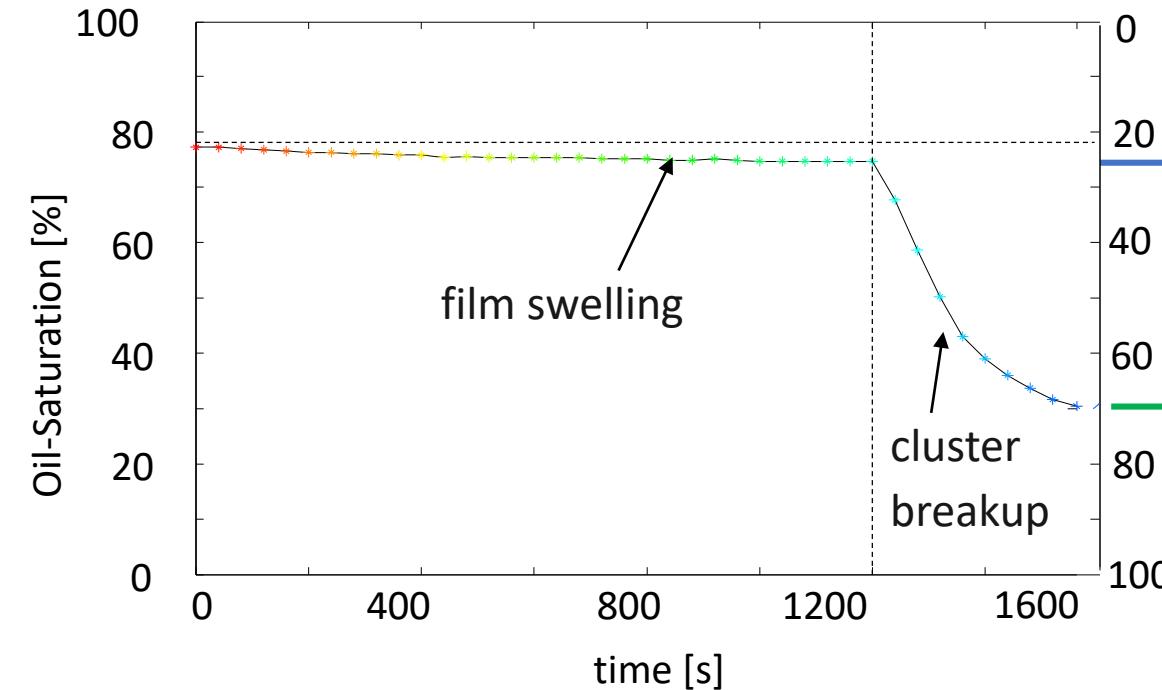
$$M_1^n = (\epsilon^{wn} + \epsilon^{ns})V$$

$$M_2^n = (J_w^{wn} \epsilon^{wn} + J_s^{ns} \epsilon^{ns})V$$

$$M_3^n = \chi^n$$

# Cluster Dynamics Introduces Topological Changes

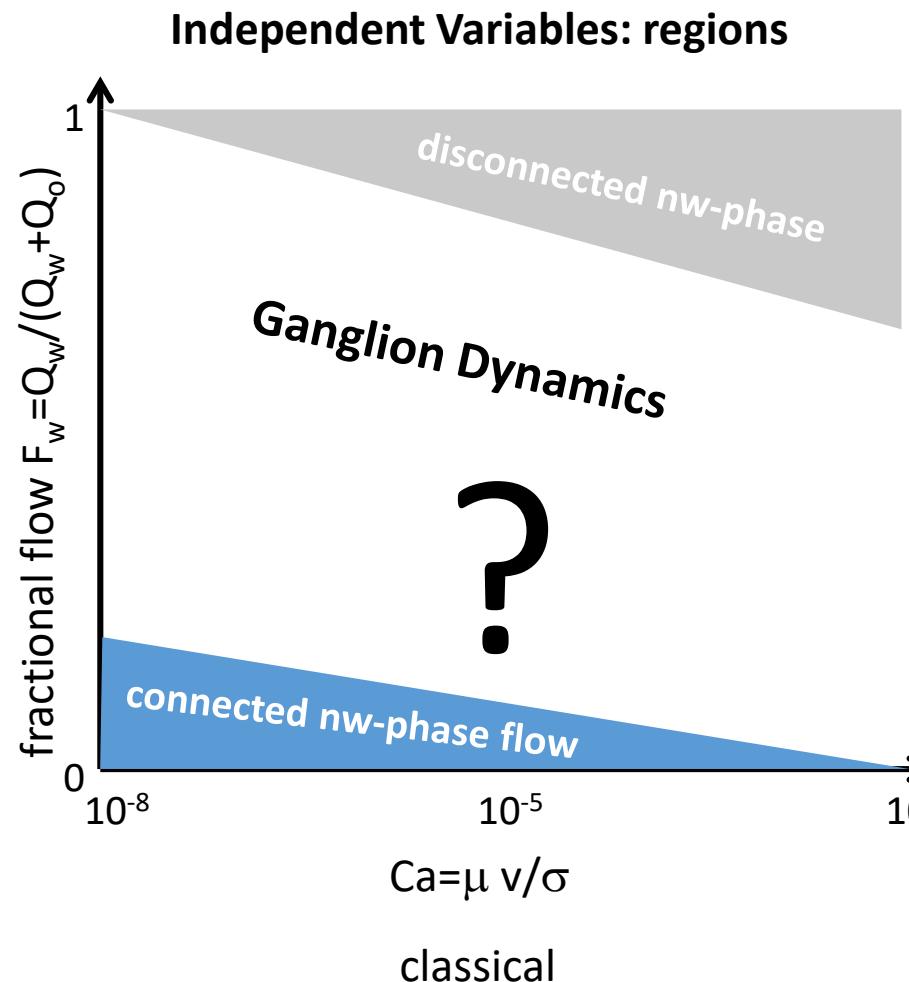
[Rücker et al., GRL, 2015]



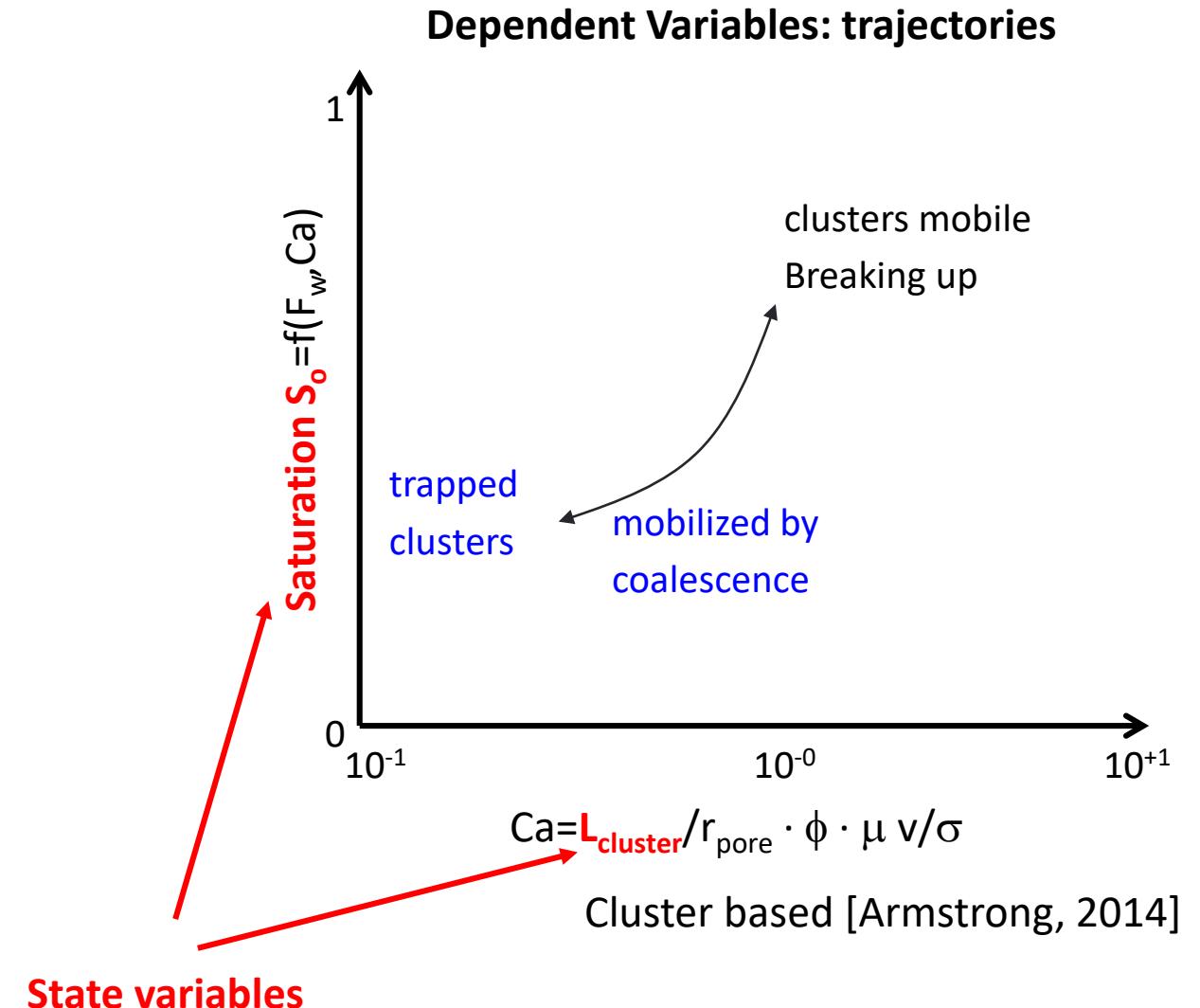
- Onset of cluster breakup at relatively oil saturation
- Oil clusters remain mobile over large saturation range
- Oil clusters immobilized / trapped close to  $S_{or}$

Co-existence of connected pathway flow and ganglion dynamics over most of the mobile saturation range

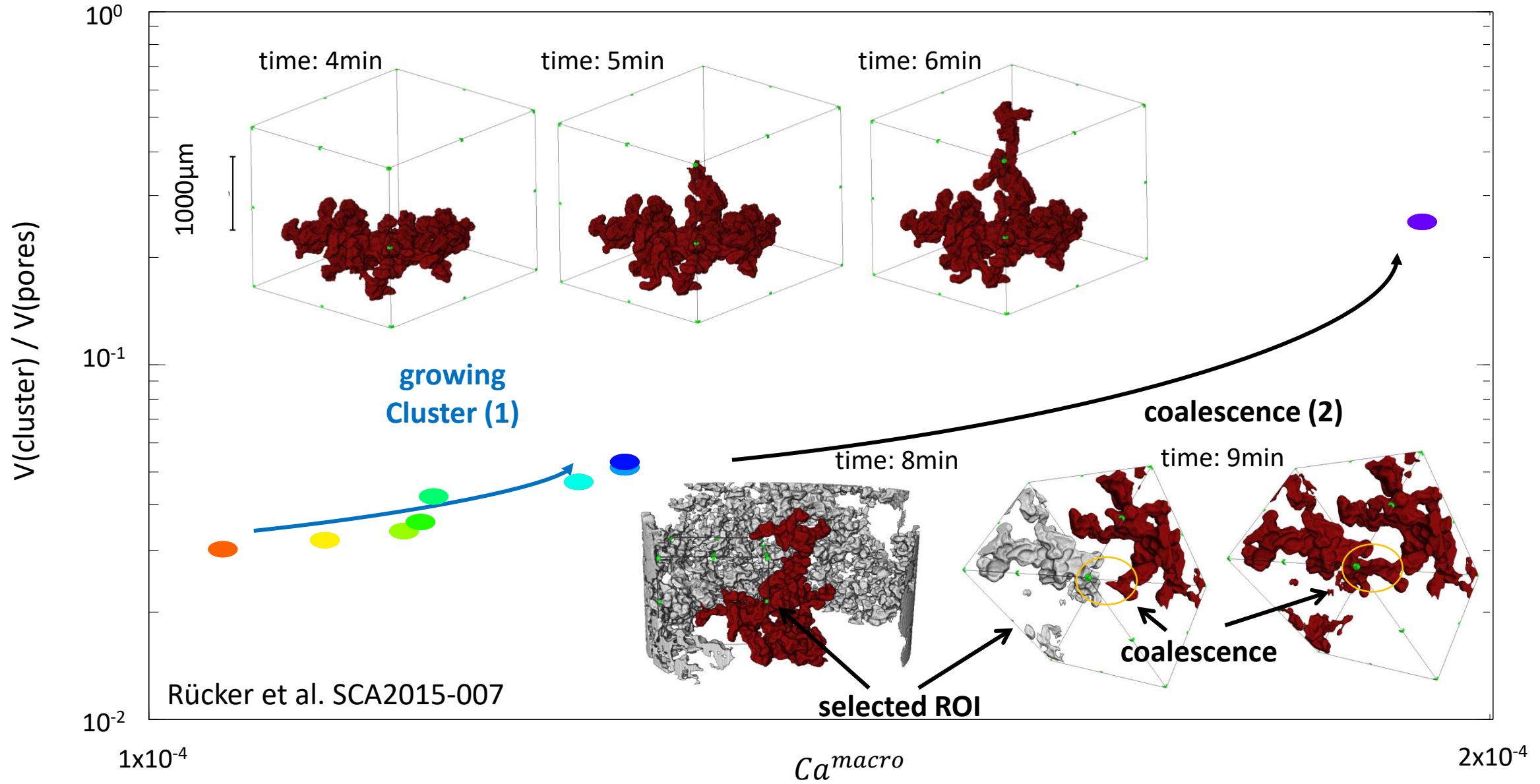
# Characterization of Flow Regimes: Phase Diagrams ...



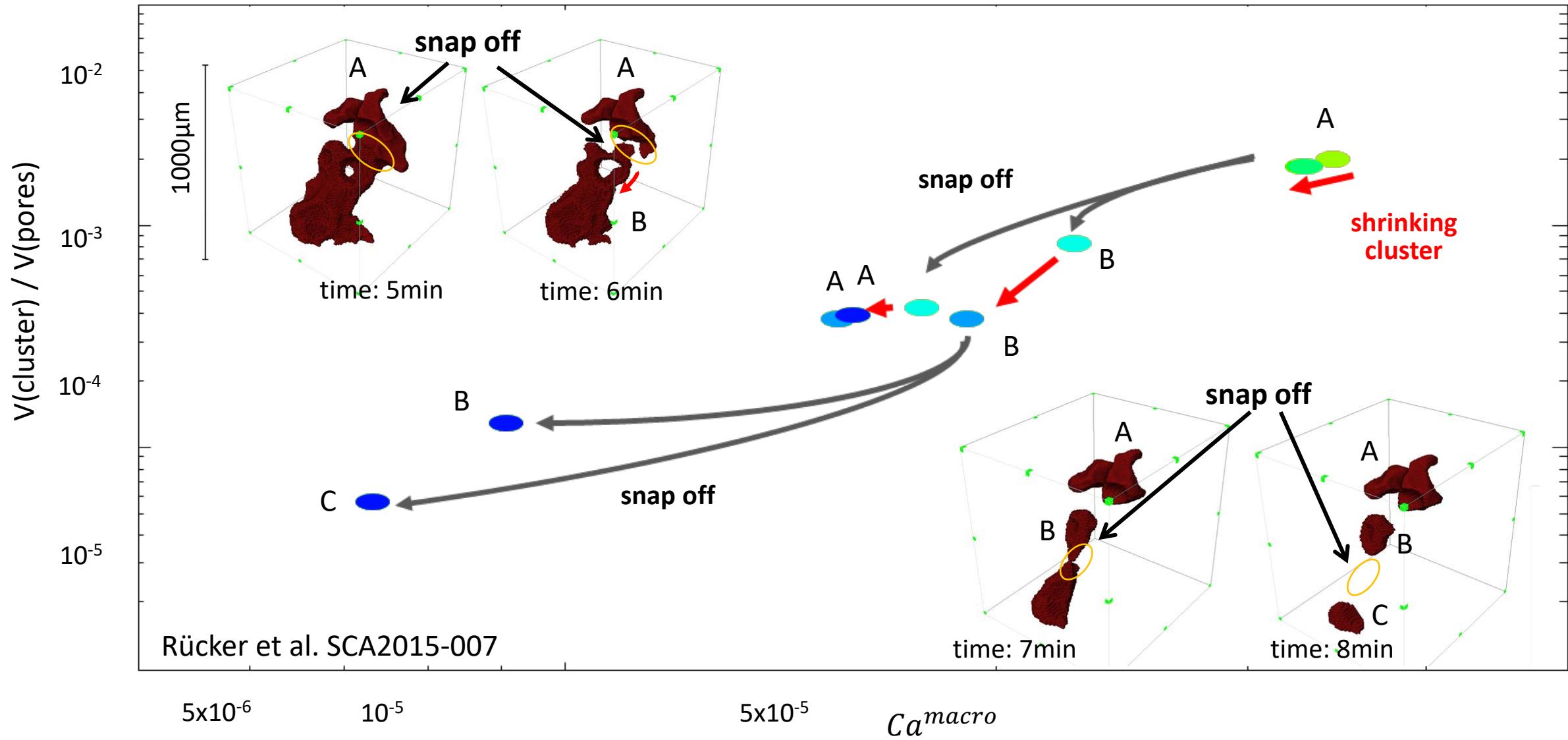
(at the time we did not know yet that we should have used  $A_{nw}$ )



# Clusters: Growing and Coalescence → $\chi$ decreases



# Clusters: Break-up by Snap-off $\rightarrow \chi$ increases



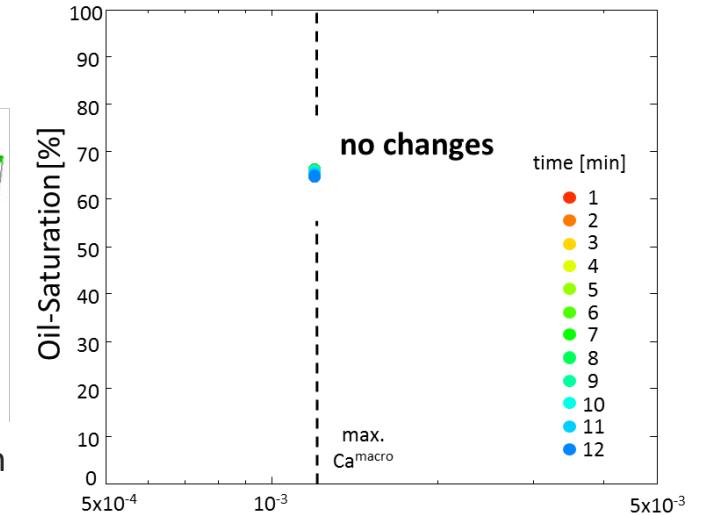
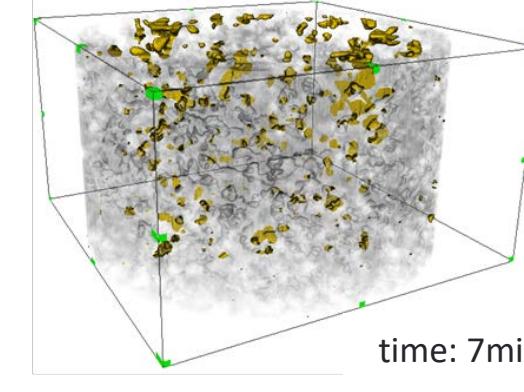
# Characterization of Flow Regimes: Phase Diagrams ...

Connected pathway flow

time: 3min

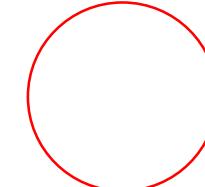
time: 5min

time: 7min

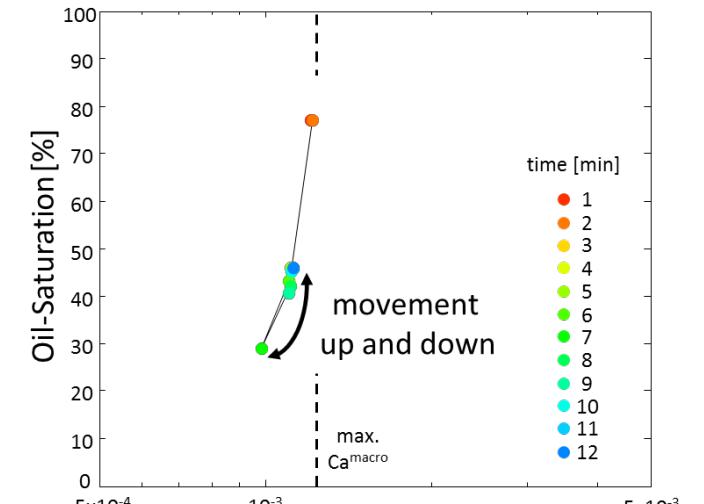
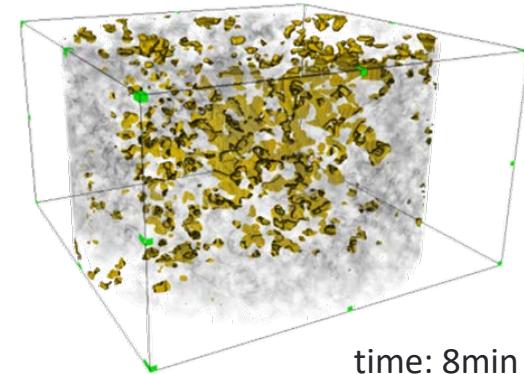


Ganglion dynamics

time: 6min



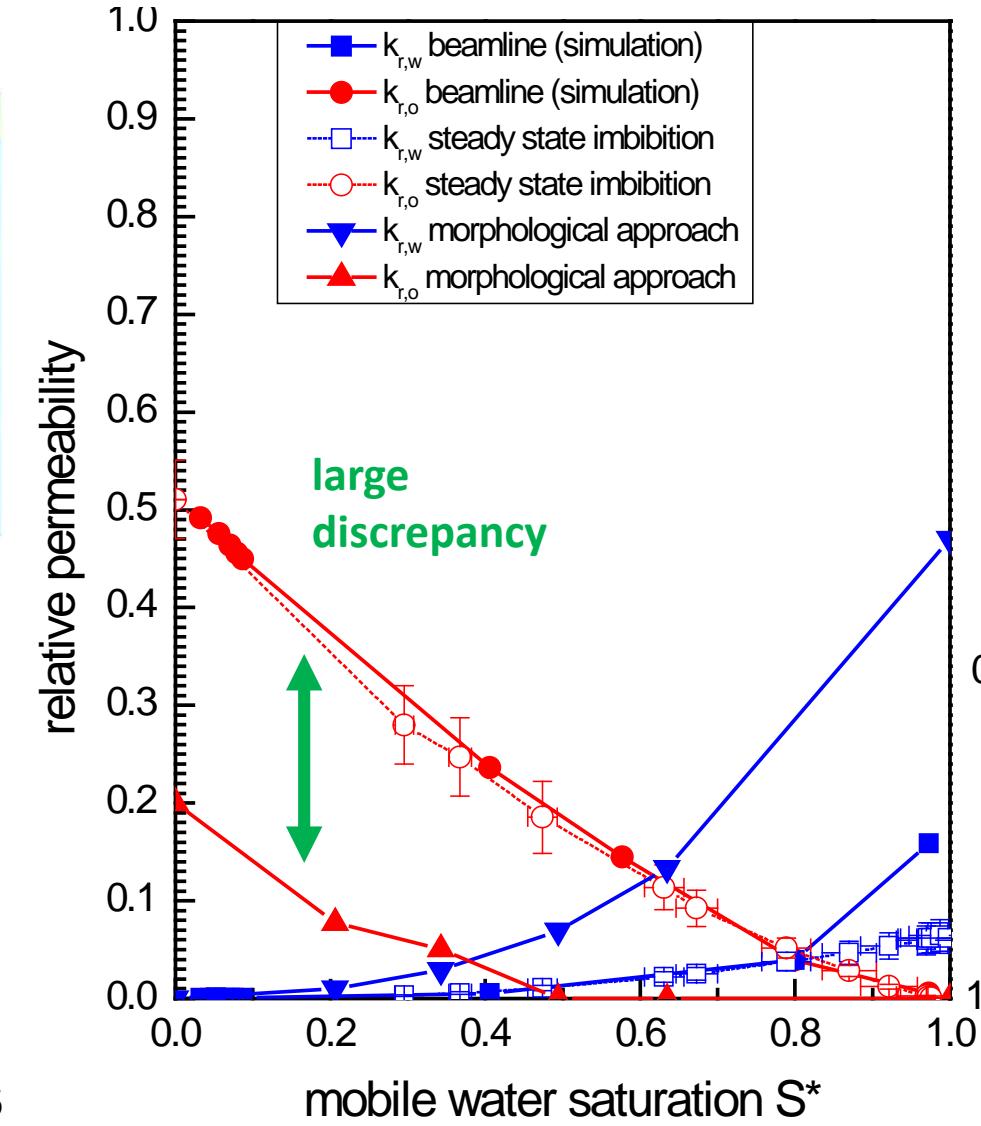
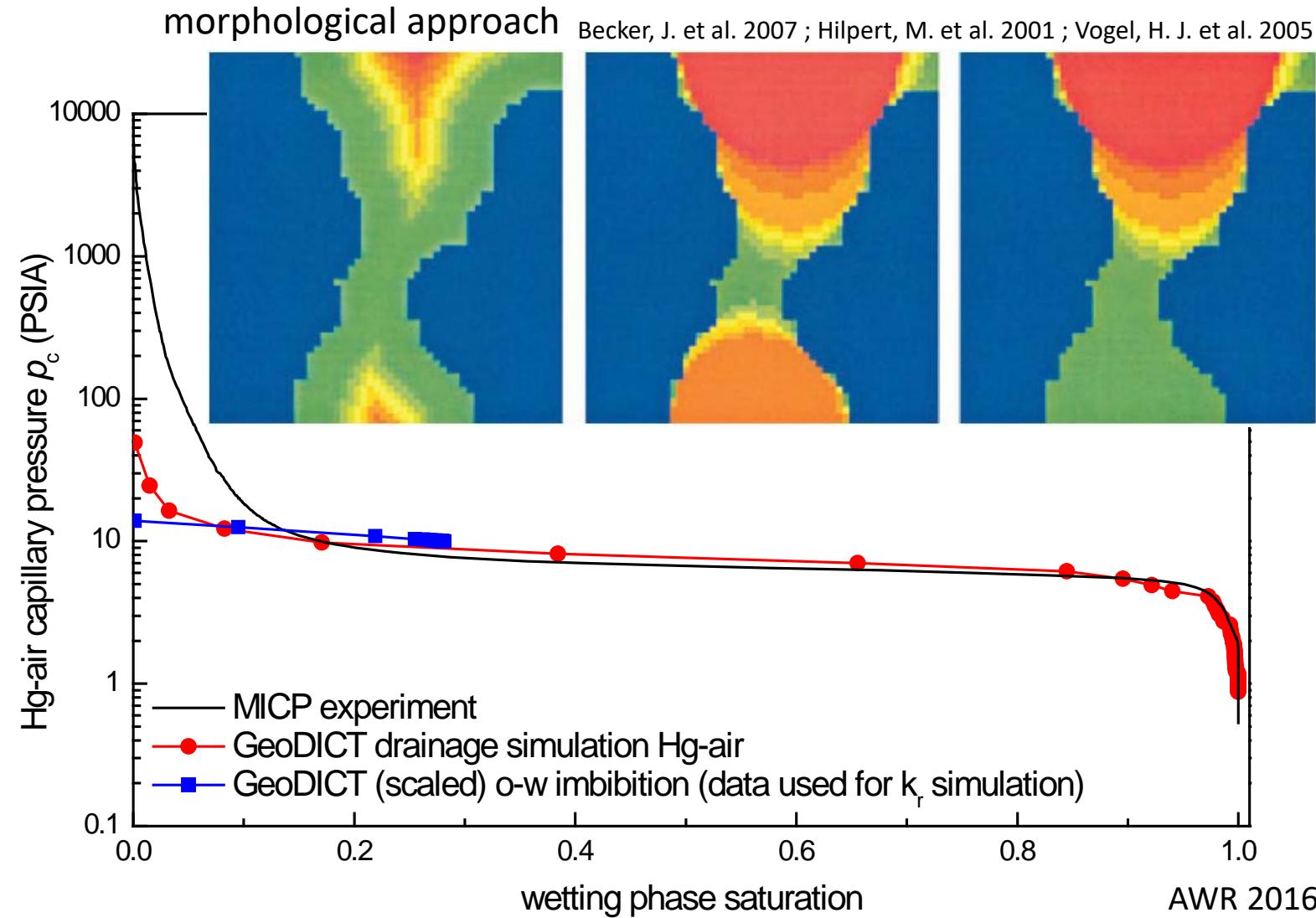
time: 7min



$$\text{Ca}_o^{\text{macro}} = \frac{l^{cl}}{r^p} \varphi \frac{\mu_w v}{\gamma_{wnw}},$$

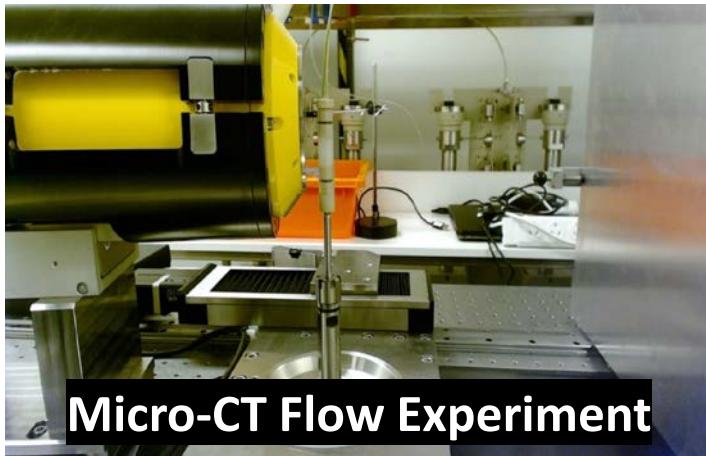
# Application: Validation of Pore Scale Simulation

Can we obtain imbibition relative permeability from a quasi-static approach ?

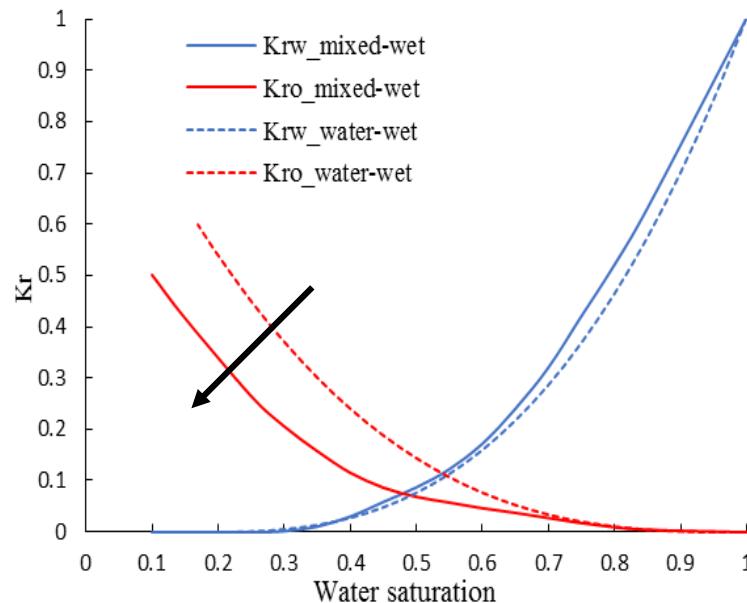


# Influence of Wettability on Topology

Dynamic Connectivity → [PNAS, Reynolds et al. 2017]



Micro-CT Flow Experiment



## Power Associated with Flow

$$\mathcal{P}_i = \frac{dW}{dt} = -\mathbf{q}_i \cdot \nabla p \approx \frac{\mu_i q_i^2}{K_i}$$

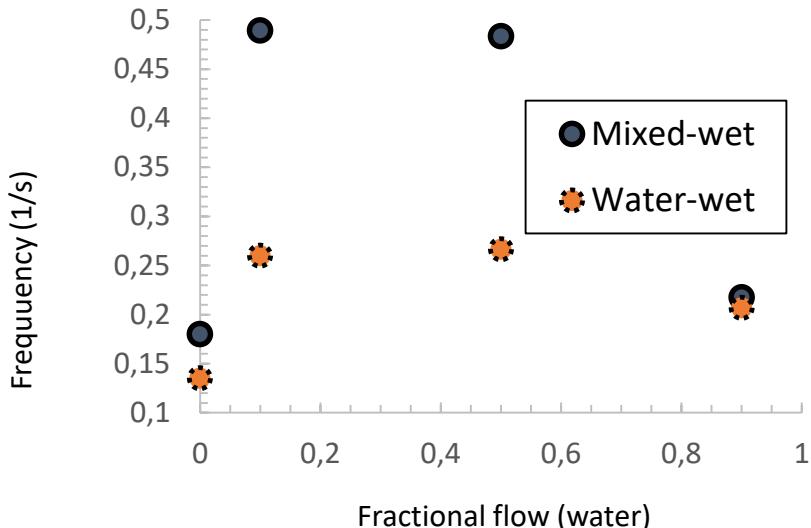
## Surface Energy of WP/NWP Interface

$$(\text{Surface energy}) E = \frac{\sigma}{l}$$

## Tested Two Different Wettabilities

Water Wet:  $l = 0.72$   
Mixed-Wet = -0.11

## Propensity of Interface Creation



## Generation of Unknown Phase

